

FLOOD ROUTING THROUGH STORM DRAINS  
Part IV  
NUMERICAL COMPUTER METHODS OF SOLUTION

By  
V. YEVJEVICH and A. H. BARNES

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## ABSTRACT

This fourth part of a four-part series of hydrology papers on flood routing through storm drains presents computer-oriented numerical methods for solving the two quasi-linear hyperbolic partial differential equations known as the De Saint-Venant equations of gradually varied free-surface unsteady flow. Formulation and description of various finite-difference schemes based on explicit methods include the "unstable", diffusing, upstream differencing, leap frog, and Lax-Wendroff schemes. Stability and convergence are examined for these various schemes of the explicit method. Using various criteria of comparison, the specified intervals scheme of the method of characteristics, the Lax-Wendroff scheme, and the diffusing scheme are compared. Of the above explicit schemes in using the finite-difference ratios in the two partial differential equations, it is found that the Lax-Wendroff scheme with the second-order interpolation for dependent variables is the most accurate stable scheme. The specified intervals scheme of the method of characteristics, using either the first-order or second-order interpolations for the dependent variables, is also discussed. It is concluded that this scheme, based on the method of characteristics and using the second-order interpolations, is the most accurate numerical integration scheme of all those studied. Flow charts, computer programs, variable conversion tables, and sample inputs and outputs, for the three numerical computer schemes, the diffusing scheme, the Lax-Wendroff scheme, and the specified intervals scheme of the method of characteristics, used in the solution of the De Saint-Venant equations, are given in appendices 1 through 3.

# FLOOD ROUTING THROUGH STORM DRAINS

## Part IV

### NUMERICAL COMPUTER METHODS OF SOLUTION

by

V. Yevjevich\* and A. H. Barnes\*\*

#### Chapter 1

#### INTRODUCTION

##### 1.1 General Classification of Partial Differential Equations

Partial differential equations of physical processes fall within one of three forms, depending on the character of the coefficients of the partial derivatives. The equations expressing the one-dimensional gradually varied free-surface unsteady flow result in what is termed the hyperbolic form of partial differential equations. These equations are characterized by the initial conditions of the dependent variables being known, given, or independently evaluated at all distance positions for the time selected as zero, the boundary conditions being independently established at two distance locations, and the process being continued indefinitely in time within the established boundary conditions. As time increases, the effect of the initial conditions becomes less influential as the boundary conditions dominate the process.

The hyperbolic partial differential equations contrast the elliptic differential equations in which the process is not time dependent. In this case the initial conditions are the boundary conditions and are independent of time. A typical process described by this form is a two-dimensional temperature distribution in a thin plate with prescribed boundary conditions along the edges.

The third type of partial differential equations are parabolic equations, with the solution requirements being similar to the hyperbolic form. The simplest parabolic equation is the one-dimension heat-flow equation.

In subsequent text only the hyperbolic partial differential equation for gradually varied free-surface unsteady flow are discussed.

##### 1.2 Continuity and Momentum Equations of Unsteady Flow

The two basic quasi-linear hyperbolic partial differential equations of gradually varied free-surface unsteady flow are derived in Chapter 3, Part I, Hydrology Paper No. 43, as Eqs. 3.23 and 3.19, and are reproduced here in their final dimensionless forms. The continuity equation is

$$\frac{A}{VB} \frac{\partial V}{\partial x} + \frac{\partial y}{\partial x} + \frac{1}{V} \frac{\partial y}{\partial t} = \frac{q}{VB} \quad (1.1)$$

and the momentum equation is

$$\frac{\alpha V}{g} \frac{\partial V}{\partial x} + \frac{\beta}{g} \frac{\partial V}{\partial t} + \frac{\partial y}{\partial x} = (S_o - S_f) - \beta \frac{Vq}{Ag} \quad (1.2)$$

in which

A = the cross-section area,  
V = the mean cross-section velocity as a dependent variable,  
y = the water depth in the conduit as a dependent variable,  
x = the length along the conduit as an independent variable,  
t = the time as an independent variable,  
B = the water surface width,  
 $\alpha$  = the energy velocity distribution coefficient,  
 $\beta$  = the momentum velocity distribution coefficient,  
g = the gravitational acceleration,  
 $S_o$  = the slope of the conduit invert,  
 $S_f$  = the energy gradient, and  
q = the distributed lateral inflow (or outflow) as discharge per unit length of the conduit.

The energy gradient, measuring the energy head loss along the conduit, is expressed in this study by the Darcy-Weisbach equation in the form

$$S_f = \frac{fV^2}{8gR} \quad (1.3)$$

in which f is the Darcy-Weisbach friction factor, R is the hydraulic radius of a partially full conduit, with  $R = A/P$ , and P is the wetted perimeter.

The friction factor (f) is expressed as a function of Reynolds number,  $R_e = VR/\nu$ , with  $\nu$  the kinematic viscosity of the water.

Equations 1.1 and 1.2 generally give the closest approximations of the actual flood movement through channels and conduits, if the basic conditions for applying the two equations are approximately satisfied. The most important condition is that of gradual variability of the flood hydrograph; this condition is nearly always fulfilled for storm floods entering into and moving along storm drains.

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### 1.3 Methods of Solving Equations of Unsteady Flow

All methods available in literature for solving Eqs. 1.1 and 1.2 may be grouped into analytical, graphical, and numerical procedures. The numerical procedures depend on the computational devices available.

Analytical solutions. The partial differential equations 1.1 and 1.2 have a friction slope,  $S_f$ , proportional to the square of the velocity or to the square of the discharge. Because their coefficients are functions of dependent variables ( $V$ ,  $y$ ), they are non-linear differential equations of the hyperbolic type. Because of the inherent mathematical difficulties of these non-linear and non-homogeneous equations, there is no way to carry out the analytical integration in closed form, unless many simplifications are introduced.

The classical approach, first performed by De Saint-Venant, neglects friction resistance and assumes the channel to be horizontal with wide rectangular cross sections. These assumptions deviate so much from the reality of flood-wave movement in channels and conduits that the wave characteristics resulting from analytical integration are generally not comparable with true wave characteristics. This classical approach by means of analytical integration is an extreme; it may be considered to be a rough approximation, and, in accuracy, can be compared with some of the very simple integration procedures of flood routing that are based on the water storage ordinary differential equation.

The use of analytical integration makes it necessary to approximate and simplify both the initial conditions and the boundary conditions by analytical expressions, which are used in Eqs. 1.1 and 1.2. The inflow hydrograph as the boundary condition, and the wave profile along the conduit, as the initial condition, must be mathematically approximated by considering them to be either symmetrical or asymmetrical waves, with functions of bell-shaped curves (gamma-functions, and others). The channel conditions may be represented by the cross section area or width as functions of water depth and distance along the conduit, with a roughness coefficient usually a constant, and the bottom slope being either a constant or a function of distance. The lateral inflow and outflow are taken as constant or are approximated by simple functions of channel and lateral flow characteristics, and of time.

The great diversity in shape and roughness of natural channels, free-surface flow conditions and the complexity of the pattern of the lateral inflows and outflows tend to complicate the analytical expressions that approximate these conditions to the extent that the analytical integration of the two partial differential equations becomes impossible. In summary, the two partial differential equations for unsteady flow can be integrated analytically, with expressions for wave evolution, by rather restrictive and very simplifying conditions, which generally are not acceptable for the solution of current practical problems.

For some discussions and abstracted references about the analytical solutions of simplified conditions for flood routing through conduits and channels, as well as of graphical and numerical solutions, see the "Bibliography and Discussion of Flood-Routing Methods

and Unsteady Flow in Channels" [1]\*, and the general reference list in Appendix 2 of Hydrology Paper No. 43 (Part I of this series of four papers).

Graphical solutions. The graphical solutions of equations for free-surface unsteady flow may be characterized by the following procedure. The celerity of the disturbance in the distance-time reference plane,  $(x,t)$  - plane, is computed from the simplified wave relation

$$\frac{dx}{dt} = V \pm \sqrt{gy_*}, \quad (1.4)$$

in which  $V$  is the mean velocity of flow,  $y_*$  is the hydraulic depth  $(A/B)$  in any cross-sectional shape, and  $g$  is the gravitational acceleration.

The term  $C = \sqrt{gy_*}$  is usually referred to as the celerity of a small disturbance moving in a quiescent water of a channel. The terms  $V + \sqrt{gy_*}$  and  $V - \sqrt{gy_*}$  are called either the wave velocity [2, p. 540], or the celerity of a small disturbance in the moving fluid [1, p. 10]. This latter term will be used in this paper when Eq. 1.4 is discussed or used. If the first derivative,  $dt/dx$ , in the  $(x,t)$  - plane is used as the measure of the celerities of disturbances in the moving water, then the inverse of Eq. 1.4 should be used as

$$\frac{dt}{dx} = \frac{1}{V \pm \sqrt{gy_*}} \quad (1.5)$$

In case of the circular conduit in which flood waves move with gradually varied free-surface flow,  $y_*$  should be replaced by  $y_* = f(y)$ , a function of water depth.

In the discussion to follow the two directions of Eq. 1.5 will be referred to as the characteristic directions, which are first derivatives of characteristic curves, defined in Chapter 3, Part I, Hydrology Paper No. 43. Along the characteristic curves, the wave phenomenon may be expressed by the two ordinary differential equations with two dependent variables as unknowns. Thus, starting from the known values of the dependent variables ( $V$  and  $y$ ) at two locations in time ( $t$ ) and position ( $x$ ), the direction of the characteristics may be graphically plotted. From these plots, the location of the intersections in time and position can be determined. With the known time ( $t$ ) and position ( $x$ ) a finite difference solution to the two ordinary differential equations gives the corresponding dependent variables ( $V$  and  $y$ ). Repeating the procedure, the integration proceeds along the time scale for the given length of channel or conduit.

This procedure has been used extensively by Parmakian in his book on waterhammer analysis [3]. Akers and Harrison presented a similar analysis for free-surface unsteady flow in a circular channel in their paper on attenuation of flood waves in partially full pipes, [4].

The limitations of graphical procedures are immediately evident when one considers the effect of

\*[ ] Reference numbers refer to the list of references at the end of this paper.



various parameters, initial and boundary conditions, in a given problem. Thus the graphical solution has limited application at present because of the labor involved, except perhaps for the visualization of the digital computer schemes and the results to be presented.

Numerical solutions. Various numerical procedures have been used in the past. The excessive number of calculations in order to progress the solution in time, however, has limited the application of these solutions.

The two partial differential equations, 1.1 and 1.2, are usually approximated by the two finite-differences equations, replacing the increments  $(dx, dt, dV, dy)$  by the finite differences  $(\Delta x, \Delta t, \Delta V, \Delta y)$ . At the same time the partial derivatives are replaced by ratios of finite differences:  $\partial V/\partial x$  by  $\Delta V/\Delta x$ ,  $\partial V/\partial t$  by  $\Delta V/\Delta t$ ,  $\partial y/\partial x$  by  $\Delta y/\Delta x$ , and  $\partial y/\partial t$  by  $\Delta y/\Delta t$ . With  $\Delta x$  and  $\Delta t$  given,  $\Delta V$  and  $\Delta y$  are changes of dependent variables which occur for these finite differences.

The basic characteristics of the above finite-difference approximations are: (1) the accuracy depends on the size and relation of finite differences  $\Delta t$  and  $\Delta x$ ; (2) the smaller the  $\Delta x$ , the more involved the computation work, but also the greater the accuracy may be, and (3) the values of dependent variables computed for the end of a  $\Delta t$  become the initial values for the next  $\Delta t$ .

With the development of electronic computers, which provide fast and relatively inexpensive computations, the past drawbacks in economy of performing the operations of the finite-differences method of integration are largely eliminated. The method is highly favored inasmuch as it is the most accurate of all practical methods of flood routing in channels and conduits. The advent of new numerical schemes helped this progress in the use of numerical methods of solution by digital computers.

The results of integration are given for two dependent variables as functions  $V = F_1(x, t)$  and  $y = F_2(x, t)$ . These two functions represent surfaces in the space  $(V, x, t)$  and  $(y, x, t)$ . If there is any discontinuity in the four partial derivatives of Eqs. 1.1 and 1.2, these discontinuities propagate along the channel, and the projection of the position of discontinuities at surfaces  $F_1$  and  $F_2$  in the  $(x, t)$ -plane produces lines that are called "characteristics", or "characteristic lines". These lines are usually curves, but in application may be replaced by straight lines along the finite differences  $\Delta x$  and  $\Delta t$ .

The simplified characteristic lines are usually given in the form

$$dx = (V \pm \sqrt{gy_*}) dt, \quad (1.6)$$

and

$$d(V \pm 2\sqrt{gy_*}) = g(S_0 - S_f)dt, \quad (1.7)$$

which are equivalent to Eqs. 1.1 and 1.2. The hydraulic depth  $(y_*)$  should be expressed as a function of  $y$  for the free-surface flow in circular conduits.

Equations 1.6 and 1.7 are usually numerically integrated by replacing  $dx$  and  $dt$  with  $\Delta x$  and  $\Delta t$ , and  $d(V \pm 2\sqrt{gy_*})$  with  $\Delta(V \pm 2\sqrt{gy_*})$ . Several numerical procedures have been developed for these approximations in the finite-differences form.

Certain features of the method of numerical integration by characteristics are important for applicability in practical cases in flood routing by finite differences: (1) the long wave is assumed to be composed of many elementary waves in the form of small surges so that for the time  $\Delta t$  and the reach  $\Delta x$ , the velocity change,  $\Delta V$ , and height change,  $\Delta y$ , are considered as discontinuities traveling with celerities  $V \pm \sqrt{gy_*}$  (providing only a rough approximation in the case of long flood waves, where the friction forces are not negligible); (2) the straight-line characteristics are used as approximations instead of curve-line characteristics for  $\Delta x$  and  $\Delta t$ , and (3) some complexity of procedure when friction factors, channel slope, sudden changes of cross section, bifurcations, junctions, and similar changes, are to be taken into consideration.

With the advent of computers and new numerical schemes, numerical integration by finite differences of Eqs. 1.6 and 1.7 has become economical. The general applicability of various electronic computers (analog, hybrid, digital) to the numerical integration either of Eqs. 1.1 and 1.2, or of Eqs. 1.6 and 1.7, is discussed in the next subchapter.

Concluding remarks. All three methods -- analytical, graphical, and numerical -- by finite differences applied either to partial differential equations or to characteristic differential equations, when applicable, give sufficiently accurate results if the methods are extended to their limits of accuracy. These methods can be successfully applied to the analysis of particular waves that have been observed. The practical prediction of wave movement, however, requires a considerable amount of work, especially when the network of drains is complex.

The mathematical difficulties of analytical integration of the two partial differential equations, the need for a large amount of data for the graphical methods, the accompanying drawbacks of time-consuming procedures and the cost in applying the approximate methods of numerical integration have provided incentive for developing simpler, but generally less accurate, flood-routing methods [1]. Since the objective of this study is to produce research results that lead to practical methods in using complete Eqs. 1.1 and 1.2, or Eqs. 1.6 and 1.7, in routing flood hydrographs through storm drains, the only acceptable integration methods from both economic and accuracy standpoints are numerical methods by finite differences and the use of electronic computers. This paper is, therefore, concerned only with these latter methods.

#### 1.4 Computer Oriented Numerical Solutions

The obvious conclusion to the dilemma of excessive repetitive calculations and the limit of manual computations is the use of electronic computers. Three possibilities exist for the solution of the problem equations.

One type of computer is the analog computer in which the mathematical functions are simulated by suitable amplifiers, potentiometers or other electronic elements. The combination of these elements simulate the mathematical equations of the physical phenomenon.

This technique is particularly desirable for a physical system with fixed parameters and repetitive operations. This analog system permits an evaluation of the effect of variations in boundary conditions. A disadvantage of the analog solution would be the problems of generating the geometric and hydraulic parameters at each stage in the computations.

The hybrid electronic computer permits continuous evaluation of the differential equations by analog and evaluates the required parameters by digital computation. Thus, a continuous solution can be obtained with the geometric and hydraulic parameters evaluated by direct computation. The availability of such computers is still limited, but hybrid computers may become the best computational device for unsteady flow. The programming is specialized and not readily usable by most programmers. For these reasons the more conventional digital computer has been generally used and will be discussed exclusively in this paper.

The digital computer presents the advantage of rapid arithmetical operations and a relatively simple and versatile programming capability. The basic limitation is that integration cannot be expressed as a continuous function as is done in the analog computer. This requires that any integration of an equation or a set of equations be represented by a series of discrete elements. The approximation to the correct integration would be expected to improve as the size of the discrete elements decreased and their number increased. This is an acceptable assumption for many integration processes. However, it cannot be assumed that it is correct for all cases. This is due to the effect of round-off and truncation errors within the computer. For this study it has been assumed that the functions to be integrated are "well

behaved" and may be reasonably integrated by the assumption of discrete increments of the variables of integration.

There are a large variety of numerical integration procedures available for the solution of the St-Venant partial differential equations of gradually varied free-surface unsteady flow. One method of categorization of these basic procedures is to consider solutions depending on the two partial differential equations of 1.1 and 1.2 of the phenomenon; in the other method solutions depend on the ordinary differential equation forms, Eqs. 1.6 and 1.7, of the same equations. How the forms of the ordinary differential equations are derived from the partial differential equations is shown in Chapter 3 of Part I, Hydrology Paper No. 43.

#### 1.5 Objectives of Studies Presented in this Paper

The objectives of this paper are to present only the results of studies concerning the numerical solutions by various finite-differences schemes, either for the case of the two partial differential equations, 1.1 and 1.2, or for the case of the four characteristics equations, 1.6 and 1.7. Chapter 2 analyzes the applicability of various finite-difference schemes in the numerical solution of the two partial differential equations. Chapter 3 analyzes the various finite-difference schemes in the numerical solution of the four characteristic equations. The applicability of various schemes is discussed at the end of each of these two chapters. Chapter 4 is a comparison of the best finite-difference schemes in the case of numerical solution of partial differential equations and numerical solution of characteristic equations. Chapter 5 presents the conclusions and recommendations for further research.

## INTEGRATION OF PARTIAL DIFFERENTIAL

## EQUATIONS BY FINITE DIFFERENCES

## 2.1 Finite-Difference Methods

The finite-difference methods of numerical integration to be discussed refer to the partial differential equations of gradually varied free-surface unsteady flow. Because these equations do not permit a closed analytical solution, approximate numerical methods of integration must be employed. Since all numerical integration methods are fundamentally finite-difference procedures some distinctions between various methods or schemes are appropriate.

For this presentation, the term "finite-difference method" will refer to the approximation to the partial derivatives as the ratios of differences of finite values of the dependent variables at fixed uniform intervals. The ratios of finite differences will approach the partial derivatives as the intervals or differences become smaller. The basic definition of a partial derivative in  $x$  of a two-variable function,  $f(x, y)$ , is

$$\frac{\partial f(x, y)}{\partial x} = \lim_{\Delta x \rightarrow 0} \left[ \frac{f(x + \Delta x, y) - f(x, y)}{\Delta x} \right] \quad (2.1)$$

Using the right side of this equation, the partial derivative may be approximated as nearly accurate as desired by selecting a small difference  $\Delta x$ .

For solving De Saint-Venant equations 1.1 and 1.2 difference approximations are made as follows. Since there are two independent variables and two dependent variables, designation of the time-distance locations of the variables will be based on the subscripts and superscripts of the variables. The subscript will refer to the distance (space) location, and the superscript to the time location as shown in Fig. 2.1.

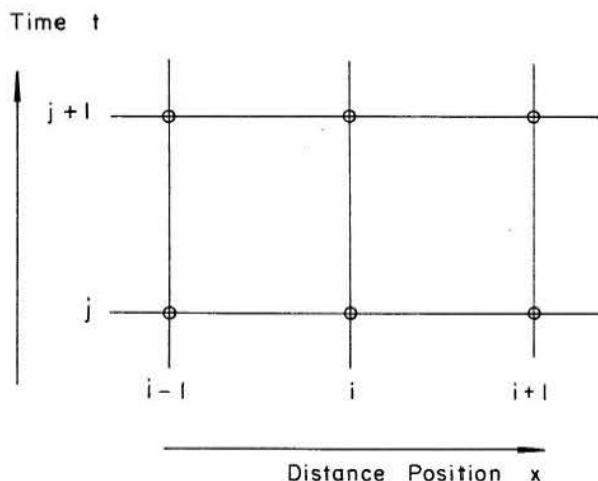


Fig. 2.1. Definition graph for the finite-difference scheme.

Thus, the depth at distance location  $i$  and at time location  $j$  is designated as  $y_i^j$ . The four partial derivatives of Eqs. 1.1 and 1.2 may be approximated by

$$\frac{\partial V}{\partial x} \approx \frac{V_{i+1}^j - V_i^j}{x_{i+1}^j - x_i^j} \quad (2.2)$$

$$\frac{\partial V}{\partial t} \approx \frac{V_i^{j+1} - V_i^j}{t_i^{j+1} - t_i^j} \quad (2.3)$$

$$\frac{\partial y}{\partial x} \approx \frac{y_{i+1}^j - y_i^j}{x_{i+1}^j - x_i^j} \quad (2.4)$$

and

$$\frac{\partial y}{\partial t} \approx \frac{y_i^{j+1} - y_i^j}{t_i^{j+1} - t_i^j} \quad (2.5)$$

The unknown quantities in these expressions are generally the values at the incremental time locations,  $j+1$ . Thus  $V_i^{j+1}$  and  $y_i^{j+1}$  are the unknown values.

With the two equations of unsteady flow, these two unknowns may be solved for simultaneously. This procedure is referred to as an explicit scheme in that the conditions at a later time,  $j+1$ , are determined directly from the conditions at the preceding time,  $j$ . Other explicit schemes are presented in the next subchapter.

Another manner of expressing the partial derivatives with respect to the distance position is

$$\frac{\partial V}{\partial x} \approx \frac{V_{i+1}^{j+1} - V_i^{j+1}}{x_{i+1}^{j+1} - x_i^{j+1}} \quad (2.6)$$

and

$$\frac{\partial y}{\partial x} \approx \frac{y_{i+1}^{j+1} - y_i^{j+1}}{x_{i+1}^{j+1} - x_i^{j+1}} \quad (2.7)$$

The partial derivatives in the case of Eqs. 2.6 and 2.7 are described in terms of the independent variable  $x$  along the incremental time locations. Therefore, there are four unknowns of  $V$  and  $y$ , at two distance locations at a given incremental time location. The two equations of unsteady flow at a given point in time and distance are insufficient for

the solution. However, if a system of simultaneous equations are developed for each point, there will be as many equations as the total number of unknowns. A simultaneous solution of this set then results in the desired solution. This scheme is referred to as the implicit solution since all solutions are directly interrelated. No attempt was made to use this method, however, because of the limits in solving equations for the dependent variables at an unlimited number of distance locations.

A physical and, consequently, mathematical limitation to either an explicit or implicit scheme is imposed by the direction a disturbance travels in the time-distance reference plane. The directions of a disturbance are commonly referred to as the characteristic directions and are defined by Eq. 1.5. The two expressions for  $dt/dx$  of Eq. 1.5 represent the two directions the disturbances propagate along.

If one considers these directions as emanating from a single given point in the time-distance plane, where a disturbance occurred, the region  $x$  and  $t$  between these two directions is affected by the disturbance. This region is the "region of influence". If one considers the disturbances as having occurred at two different locations in the time-distance plane, two of the four directions will intersect. The region bounded by this intersection is the "domain of dependence." The dependent variables in this region are functions of all their previous values within this region. As a corollary, the values of dependent variables outside this region do not affect the values of  $V$  and  $y$  inside this region.

Thus, the directions of the disturbance or characteristic directions in the  $(x, t)$  - plane divide the time-distance plane into a region wherein solutions from given conditions are possible, and a region in which solutions are theoretically impossible. It is necessary to consider this in any finite-difference method of integrating the two partial differential equations. The general criterion to be applied is that

$$\frac{dt}{dx} \approx \frac{1}{V \pm \sqrt{g A/B}} \quad (2.8)$$

in which  $V$  and  $A/B$  are the average values for the specified finite differences,  $\Delta x$  and  $\Delta t$ . The criteria of Eq. 2.8 is valid for all values of the dependent variables in the solution. The nearer the two points in the  $(x, t)$  - plane are, the more nearly the numerical solution will approach the true solution.

## 2.2 Various Finite-Difference Schemes

Equations 2.2 through 2.5 present the simplest approximation by the finite-difference expressions to the partial derivatives. A wide variety of schemes, usually more sophisticated than Eqs. 2.2 through 2.5, have been developed by various authors to provide better accuracy and to maintain the stability of the solutions with minimum computational work.

Richtmeyer [5] presented six schemes with their corresponding truncation errors. These schemes are presented in Table 2.1. This table displays the computational template of the  $(x, t)$  - plane, the approximation to the partial derivatives, and the order of the truncation error  $O(\Delta)$ , due to the approximation where  $\Delta$  is the symbol of increment, either  $\Delta x$  or  $\Delta t$ .

Substituting these approximations into the basic equations results in a pair of equations with two unknowns, velocity and depth, at the end of the time interval.

The "unstable scheme" is inherently unstable. It is presented to demonstrate the simplest scheme, and to permit comparison of stable schemes with this basic scheme.

The diffusing scheme is the simplest stable scheme. It offers two approaches for computation. One approach consists of the staggered scheme as presented in Table 2.1. It uses known values of  $V$  and  $y$  at the  $i-1$  and the  $i+1$  distance positions at time  $t$  to compute the dependent variables at the distant position  $i$ , at time  $t + \Delta t$ . This approach determines values at all locations defined by  $i+j$  equal an even number. The other approach is to advance one  $\Delta x$  and thus compute the dependent variables at each intersection. This approximately doubles the computational time but produces results at one-half the intervals of the first method.

In order for the diffusing scheme to be stable, it is necessary that

$$\frac{\Delta t}{\Delta x} \leq \left| \frac{1}{V \pm \sqrt{g A/B}} \right|$$

be a condition throughout the computation. As the flow progresses into the super-critical range, this condition is less likely to be fulfilled unless an arbitrary reduction in  $\Delta t$  is made. An additional limitation of this scheme is the assumed linearity of the dependent variables within the interval from  $i-1$  to  $i+1$ .

The upstream differencing scheme is similar to the diffusing scheme. The computer programming, however, is somewhat more involved because of the necessity of deciding which representation of the distance derivative to use for each computation. For this reason this scheme was not investigated in this study.

The leap-frog scheme is an improvement over the diffusing scheme in that the time derivative is estimated from the computed values of the dependent variables at the  $t - \Delta t$  time position. The limitation of this procedure is similar to that of the diffusing scheme. An additional limitation is the required computer storage of computed values at three successive times as compared to two successive times for the other schemes.

The previously described schemes all depend on an assumption of linearity between the time-distance junctions for the description of the partial derivatives at the pivot point  $(i, j)$ . An improvement to this assumption is to recognize the rate-of-change of the derivative as defined by the known values of the dependent variables at three points. The Lax-Wendroff method provides this recognition. The procedure is described in detail in a following subchapter. The consistent reproduction of initial conditions for a constant discharge, regardless of the curvature of the water surface, is the benefit derived from this method.

The implicit scheme requires the solution of a system of simultaneous equations equal in number to the number of distance intervals plus one. Two of

Table 2.1 Various finite-difference schemes

	UNSTABLE	DIFFUSING	UPSTREAM DIFFERENCING	LEAP FROG	LAX WENDROFF	IMPLICIT
COMPUTATIONAL TEMPLATE						
PARTIAL DERIVATIVE APPROXIMATION	$\frac{\partial U}{\partial x} \approx \frac{U_{i+1}^j - U_{i-1}^j}{2\Delta x}$	$\frac{\partial U}{\partial x} \approx \frac{U_{i+1}^j - U_{i-1}^j}{2\Delta x}$	$\frac{\partial U}{\partial x} \approx \frac{U_{i+1}^j - U_i^j}{\Delta x}$ or $\frac{U_i^j - U_{i-1}^j}{\Delta x}$	$\frac{\partial U}{\partial x} \approx \frac{U_{i+1}^j - U_{i-1}^j}{2\Delta x}$	Depends on form of partial differential equation.	$\frac{\partial U}{\partial x} \approx \frac{U_{i+1}^{j+1} - U_{i-1}^{j+1} + U_{i+1}^j - U_{i-1}^j}{4\Delta x}$
TRUNCATION ERROR	$O[\Delta^2]$	$O[\Delta^2]$	$O[\Delta^2]$	$O[\Delta^3]$	$O[\Delta^3]$	$O[\Delta^3]$

these equations involve the boundary conditions. This system was not used because of the number of equations that needed to be solved simultaneously, for an arbitrarily long conduit.

All but one of the above schemes are explicit. Two of the schemes, the diffusing scheme and the Lax-Wendroff scheme, are used in this study to solve the De Saint Venant equations. These solutions provide good accuracy and require only reasonable computer time. The diffusing and Lax-Wendroff schemes are summarized in the following two subchapters.

### 2.3 Diffusing Scheme

The diffusing scheme evolves from the following approximation to the partial derivatives with respect to time. The schemes in Table 2.1 is the definition graph for the location of significant variables. It is assumed that the dependent variables are known for all positions at time  $j$ . The dependent variable will be designated as  $U$  in this development, and it may refer either to the  $V$  or  $y$  dependent variables of the two partial differential equations. The objective is to represent the partial derivatives as functions of the unknown dependent variable  $U$  at distance location  $i$  and time location  $j+1$ . The partial derivative of  $U$  with respect to  $t$  is approximated by

$$\left(\frac{\partial U}{\partial t}\right)_i \approx \left(\frac{\Delta U}{\Delta t}\right)_i, \quad (2.9)$$

in which

$$\Delta U_i = U_i^{j+1} - U_i^j. \quad (2.10)$$

Expressing  $U_i^j$  as an average

$$U_i^j = \frac{U_{i+1}^j + U_{i-1}^j}{2}, \quad (2.11)$$

then

$$\Delta U_i = U_i^{j+1} - \frac{U_{i+1}^j + U_{i-1}^j}{2}, \quad (2.12)$$

and finally the finite difference approximation to this partial derivative is

$$\begin{aligned} \left(\frac{\Delta U}{\Delta t}\right)_i &= \frac{U_i^{j+1} - \frac{U_{i+1}^j + U_{i-1}^j}{2}}{\Delta t} \\ &= \frac{2U_i^{j+1} - U_{i+1}^j - U_{i-1}^j}{2\Delta t}. \end{aligned} \quad (2.13)$$

Similarly, the partial derivative with respect to the distance  $x$  is approximated by

$$\left(\frac{\partial U}{\partial x}\right)_i \approx \left(\frac{\Delta U}{\Delta x}\right)_i, \quad (2.14)$$

in which

$$\left(\frac{\Delta U}{\Delta x}\right)_i = \frac{1}{2} \left[ \frac{U_{i+1}^j - U_i^j}{\Delta x} + \frac{U_i^j - U_{i-1}^j}{\Delta x} \right], \quad (2.15)$$

so that

$$\left(\frac{\Delta U}{\Delta x}\right)_i = \frac{1}{2\Delta x} (U_{i+1}^j - U_{i-1}^j). \quad (2.16)$$

It is to be noted that both partial derivatives are approximated for the location  $i, j$ .

### 2.4 Lax-Wendroff Scheme

The Lax-Wendroff finite difference scheme was investigated to eliminate some of the deficiencies of the diffusing scheme. The summary of the scheme is as follows. It is assumed that all functions are continuous and contain as many continuous derivatives as required. It is also assumed that products of first-order partial derivatives, and any derivative of  $S_f$  in  $x$  and  $t$  are negligible quantities.

The expressions  $\frac{\partial A}{\partial t} = B \frac{\partial y}{\partial t}$  and  $\frac{\partial A}{\partial x} = B \frac{\partial y}{\partial x}$  relate  $A, B,$  and  $y$ . Therefore, the equation of continuity reduces to

$$\frac{\partial y}{\partial t} = -\frac{A}{B} \frac{\partial V}{\partial x} - V \frac{\partial y}{\partial x}. \quad (2.17)$$

The intended application of the Taylor series requires the use of second-order partial derivatives. Thus,

$$\frac{\partial^2 y}{\partial t^2} = -\frac{A}{B} \frac{\partial^2 V}{\partial x \partial t} - V \frac{\partial^2 y}{\partial x \partial t}, \quad (2.18)$$

and

$$\frac{\partial^2 y}{\partial x \partial t} = -\frac{A}{B} \frac{\partial^2 V}{\partial x^2} - V \frac{\partial^2 y}{\partial x^2}. \quad (2.19)$$

The momentum equation, 1.2, is rewritten here in the form

$$\frac{\partial V}{\partial t} = -\frac{\alpha}{\beta} V \frac{\partial V}{\partial x} - \frac{g}{\beta} \frac{\partial y}{\partial x} - \frac{g}{\beta} (S_f - S_0), \quad (2.20)$$

which gives then

$$\frac{\partial^2 V}{\partial x \partial t} = -\frac{\alpha}{\beta} V \frac{\partial^2 V}{\partial x^2} - \frac{g}{\beta} \frac{\partial^2 y}{\partial x^2}. \quad (2.21)$$

Hence, Eq. 2.18 becomes

$$\frac{\partial^2 y}{\partial t^2} = \frac{A}{B} \frac{1}{\beta} \left( \alpha V \frac{\partial^2 V}{\partial x^2} + g \frac{\partial^2 y}{\partial x^2} \right) + \frac{VA}{B} \frac{\partial^2 V}{\partial x^2} + V^2 \frac{\partial^2 y}{\partial x^2}$$

or

$$\frac{\partial^2 y}{\partial t^2} = \left(\frac{\alpha}{\beta} + 1\right) \frac{AV}{B} \frac{\partial^2 V}{\partial x^2} + \left(\frac{g}{\beta} \frac{A}{B} + V^2\right) \frac{\partial^2 y}{\partial x^2} \quad (2.22)$$

Equation 2.20 then gives

$$\frac{\partial^2 V}{\partial t^2} = -\frac{\alpha}{\beta} V \frac{\partial^2 V}{\partial x \partial t} - \frac{g}{\beta} \frac{\partial^2 y}{\partial x \partial t} \quad (2.23)$$

Substituting Eqs. 2.19 and 2.21 into Eq. 2.23 yields

$$\frac{\partial V^2}{\partial t^2} = -\frac{\alpha}{\beta} V \left(-\frac{\alpha}{\beta} V \frac{\partial^2 V}{\partial x^2} - \frac{g}{\beta} \frac{\partial^2 y}{\partial x^2}\right) - \frac{g}{\beta} \left(-\frac{A}{B} \frac{\partial^2 V}{\partial x^2} - V \frac{\partial^2 y}{\partial x^2}\right)$$

or

$$\frac{\partial^2 V}{\partial t^2} = \left[\left(\frac{\alpha}{\beta}\right)^2 V^2 + \frac{g}{\beta} \frac{A}{B}\right] \frac{\partial^2 V}{\partial x^2} + \left(\frac{\alpha}{\beta} + 1\right) \frac{g}{\beta} V \frac{\partial^2 y}{\partial x^2} \quad (2.24)$$

Putting U as the symbol for any dependent variable V or y, then for any U(x, t) and a fixed x, a Taylor series expansion gives

$$U(t+\Delta t) = U(t) + \Delta t \frac{\partial U}{\partial t} + \frac{(\Delta t)^2}{2} \frac{\partial^2 U}{\partial t^2} + 0[(\Delta t)^3] \quad (2.25)$$

in which both  $\partial U/\partial t$  and  $\partial^2 U/\partial t^2$  are functions of t. Similarly, for a fixed t,

$$U(x+\Delta x) = U(x) + \Delta x \frac{\partial U}{\partial x} + \frac{(\Delta x)^2}{2} \frac{\partial^2 U}{\partial x^2} + 0[(\Delta x)^3] \quad (2.26)$$

and

$$U(x-\Delta x) = U(x) - \Delta x \frac{\partial U}{\partial x} + \frac{(\Delta x)^2}{2} \frac{\partial^2 U}{\partial x^2} - 0[(\Delta x)^3] \quad (2.27)$$

Subtracting Eq. 2.27 from Eq. 2.26 yields

$$\frac{\partial U}{\partial x} = \frac{U(x+\Delta x) - U(x-\Delta x)}{2\Delta x} + 0[(\Delta x)^3] \quad (2.28)$$

Adding Eq. 2.27 and Eq. 2.26 yields the approximation of the second-order partial derivative of U with respect to x

$$\frac{\partial^2 U}{\partial x^2} = \frac{U(x+\Delta x) - 2U(x) + U(x-\Delta x)}{(\Delta x)^2} + 0[(\Delta x)^4] \quad (2.29)$$

Substituting V and y for U, respectively, and using Eqs. 2.17, 2.20, 2.22, and 2.24 for the appropriate partial derivatives with respect to t in Eq. 2.25 produces

$$V(t+\Delta t) = V(t) - \frac{\Delta t}{\beta} \left[\alpha V \frac{\partial V}{\partial x} + g \frac{\partial y}{\partial x} + g(S_f - S_o)\right]$$

$$+ \frac{(\Delta t)^2}{2} \left[ \left( \frac{\alpha^2 V^2}{\beta^2} + \frac{g}{\beta} \frac{A}{B} \right) \frac{\partial^2 V}{\partial x^2} + \left(\frac{\alpha}{\beta} + 1\right) \frac{g}{\beta} V \frac{\partial^2 y}{\partial x^2} \right] + 0[(\Delta t)^3] \quad (2.30)$$

and

$$y(t+\Delta t) = y(t) - \Delta t \left( \frac{A}{B} \frac{\partial V}{\partial x} + V \frac{\partial y}{\partial x} \right) + \frac{(\Delta t)^2}{2} \left[ \left(\frac{\alpha}{\beta} + 1\right) \frac{AV}{B} \frac{\partial^2 V}{\partial x^2} + \left(\frac{A}{B} \frac{g}{\beta} + V^2\right) \frac{\partial^2 y}{\partial x^2} \right] + 0[(\Delta t)^3] \quad (2.31)$$

Let j index the t intervals and i index the x intervals. Referring to Eqs. 2.28 and 2.29, the first and second partial derivatives with respect to x are approximated by

$$\frac{\partial U}{\partial x} = \frac{U_{i+1}^j - U_{i-1}^j}{2\Delta x} \quad (2.32)$$

and

$$\frac{\partial^2 U}{\partial x^2} = \frac{U_{i+1}^j - 2U_i^j + U_{i-1}^j}{(\Delta x)^2} \quad (2.33)$$

Thus, recurrence relations for finding approximate solutions to V and y in Eqs. 2.30 and 2.31 are

$$y_i^{j+1} = y_i^j - \frac{\Delta t}{2\Delta x} \left[ \left(\frac{A}{B}\right)_i^j (V_{i+1}^j - V_{i-1}^j) + V_i^j (y_{i+1}^j - y_{i-1}^j) \right] + \frac{1}{2} \left(\frac{\Delta t}{\Delta x}\right)^2 \left\{ \left(\frac{\alpha}{\beta} + 1\right) \left(\frac{A}{B}\right)_i^j V_i^j (V_{i+1}^j - 2V_i^j + V_{i-1}^j) + \left[\frac{g}{\beta} \left(\frac{A}{B}\right)_i^j + (V_i^j)^2\right] (y_{i+1}^j - 2y_i^j + y_{i-1}^j) \right\} \quad (2.34)$$

and

$$V_i^{j+1} = V_i^j - \frac{\Delta t}{2\beta\Delta x} \left[ \alpha V_i^j (V_{i+1}^j - V_{i-1}^j) + g(y_{i+1}^j - y_{i-1}^j) + 2g\Delta x(S_f - S_o) \right] + \frac{1}{2} \left(\frac{\Delta t}{\Delta x}\right)^2 \left\{ \left[\left(\frac{\alpha}{\beta}\right)^2 (V_i^j)^2 + \frac{g}{\beta} \left(\frac{A}{B}\right)_i^j\right] (V_{i+1}^j - 2V_i^j + V_{i-1}^j) + \left(\frac{\alpha}{\beta} + 1\right) \frac{g}{\beta} V_i^j (y_{i+1}^j - 2y_i^j + y_{i-1}^j) \right\} \quad (2.35)$$

For those cases in which the products of the first order partial derivatives and the derivatives of  $S_f$  cannot be disregarded, difference equations analogous to Eqs. 2.34 and 2.35 may be derived by appropriate substitutions of relations from Table 2.2 into Eqs. 2.25, 2.26, and 2.27.

TABLE 2.2  
Substitutions

The substitutions in the following equations are:

$$M = \frac{(1 - \frac{2y}{D})}{\sqrt{\frac{y}{D}(1 - \frac{y}{D})}}, \text{ with } D \text{ the conduit diameter;}$$

$$N = \frac{1}{D} \left\{ \frac{B}{\cos^{-1}(1 - \frac{2y}{D})} - \frac{A}{D \sqrt{\frac{y}{D}(1 - \frac{y}{D})} [\cos^{-1}(1 - \frac{2y}{D})]^2} \right\};$$

$$\frac{\partial B}{\partial x} = M \frac{\partial y}{\partial x}, \quad \frac{\partial B}{\partial t} = M \frac{\partial y}{\partial t}, \quad \frac{\partial R}{\partial x} = N \frac{\partial y}{\partial x}, \quad \text{and} \quad \frac{\partial R}{\partial t} = N \frac{\partial y}{\partial t};$$

$$\frac{\partial^2 y}{\partial x \partial t} = \frac{\partial V}{\partial x} (-2 \frac{\partial y}{\partial x} + \frac{A}{B^2} \frac{\partial B}{\partial x}) - \frac{A}{B} \frac{\partial^2 V}{\partial x^2} - V \frac{\partial^2 y}{\partial x^2};$$

$$\frac{\partial^2 y}{\partial t^2} = - \frac{\partial V}{\partial x} \frac{\partial y}{\partial t} - \frac{A}{B^2} \frac{\partial B}{\partial t} - \frac{A}{B} \frac{\partial^2 V}{\partial x \partial t} - V \frac{\partial^2 y}{\partial t^2} - \frac{\partial V}{\partial t} \frac{\partial y}{\partial t};$$

$$\frac{\partial^2 V}{\partial x \partial t} = - \frac{\alpha}{\beta} \frac{\partial V}{\partial x} \frac{\partial V}{\partial x} + V \frac{\partial^2 V}{\partial x^2} - \frac{g}{\beta} \frac{\partial^2 y}{\partial x^2} - \frac{\alpha}{\beta} \frac{f}{8} \left( \frac{2RV \frac{\partial V}{\partial x} - V^2 \frac{\partial R}{\partial x}}{R^2} \right)$$

and

$$\frac{\partial^2 V}{\partial t^2} = - \frac{\alpha}{\beta} \left( \frac{\partial V}{\partial x} \frac{\partial V}{\partial t} + V \frac{\partial^2 V}{\partial x \partial t} \right) - \frac{g}{\beta} \frac{\partial^2 y}{\partial x \partial t} - \frac{\alpha}{\beta} \frac{f}{8} \left( \frac{2RV \frac{\partial V}{\partial t} - V^2 \frac{\partial R}{\partial t}}{R^2} \right).$$

## 2.5 Comparison of Solutions by the Two Schemes

Comparing the solutions of both water depth and water velocity at various times and distances would be redundant. Since the analytical and physical waves will be compared by their water depths at a given position, solutions of  $y$  alone are considered. In this analysis, comparison is made for the theoretical dimensions of the experimental conduit, approximately 3 feet in diameter and 822 feet long. In the subsequent plots of these solutions of  $y$  let  $A_w$  be the solutions with all the derivative terms, and  $A_{w0}$  be the solutions without the terms consisting of the product of the first order derivatives and the derivatives of the energy slope, and  $D$  the solutions based on the diffusing scheme.

An important criterion of any numerical solution is the ability to repeat the values of  $y$  given at the initial conditions as best as possible over a period of time under a constant discharge. Under this steady flow, a critical  $x$  position is that which is near the downstream end of the pipe. Figure 2.2 shows the plots of  $y$  versus  $t$  at  $x = 796.7$  ft using the Lax-Wendroff Scheme developed in the previous subchapter, and the method based on the

diffusing scheme. In these two methods the total number  $n$  of  $x$  intervals used was 160, or  $\Delta x = L/n = 822/160$ . It is to be noted that after 175 seconds the maximum drops are about 0.01 and 0.07 ft for  $A_w$  and  $D_i$  schemes, respectively.

Another important criterion in a numerical solution is stability. Paraphrasing material from the Journal of Mathematics and Physics [6] stability is related to the difference between the exact solution of the difference equations and the numerical solution of these equations. This difference may be called the round-off error. In the Journal stability is defined in terms of the growth of round-off errors. That is, strong stability exists if the over-all error due to round-off errors does not grow, and weak stability exists if single round-off errors do not grow. Strong and weak instability occurs if neither of the above is true. Also stated is the assumption that weak stability implies strong stability. Thus, stability is a measure of error propagation.

The first series of tests studying the measure of error propagation was that of strong stability under a constant discharge or steady flow. That is, for both the Lax-Wendroff method and the method based on the diffusing scheme, an error of 0.001 feet was added to the initial condition at each  $x$  partition point. Simultaneously, these schemes were run over a period of time using the correct initial conditions, and these same conditions, plus the induced error were used as the starting lines. In both cases the induced error did not grow but approached zero with the developed scheme tending to zero at a faster rate.

Some effects were observed in the second series of tests with reference to weak stability, as the induced error was added only to the middle partition point. Using 81 partition points and observing the solutions of  $y$  at  $x = 4n - 3$  and  $t = 2n - 1$ , it was found that the developed solution took 225.3 seconds to zero out to five decimal places, and the diffusing scheme took 520.9 seconds.

Of more importance in the matter of stability is the third series of tests studied. This time the constant discharge input hydrograph was replaced by a varying hypothetical input hydrograph. An error of 0.001 feet was added to the initial conditions at the 81st point of a total of 160 partition points in both the Lax-Wendroff scheme and the diffusing scheme. The solutions of  $y$  for the same  $t$  and  $x$  partition points were the same as those observed for the second series of tests. After 180.9 seconds the error at point  $i = 5$  was 0.00001, and the error at the other points has zeroed out to 5 decimal places using the Lax-Wendroff scheme. The diffusing scheme solutions did not show an induced error growth either; this time the error did not stop at zero but became negative.

Thus, these series of tests indicate that both the diffusing scheme and the Lax-Wendroff scheme are stable with the latter showing the greater stability.

The next consideration regarding comparisons of solutions using the hypothetical flood input hydrograph, is that of the effect of interval size. In both the Lax-Wendroff scheme and the diffusing scheme  $\Delta t = \Delta x/4z$ , where  $z$  is the initial discharge ( $Q$ ) divided by the initial area ( $A$ ). This is done to insure that  $\Delta t$  will be small enough to fall within the domain of dependence. Figure 2.3 shows the plots



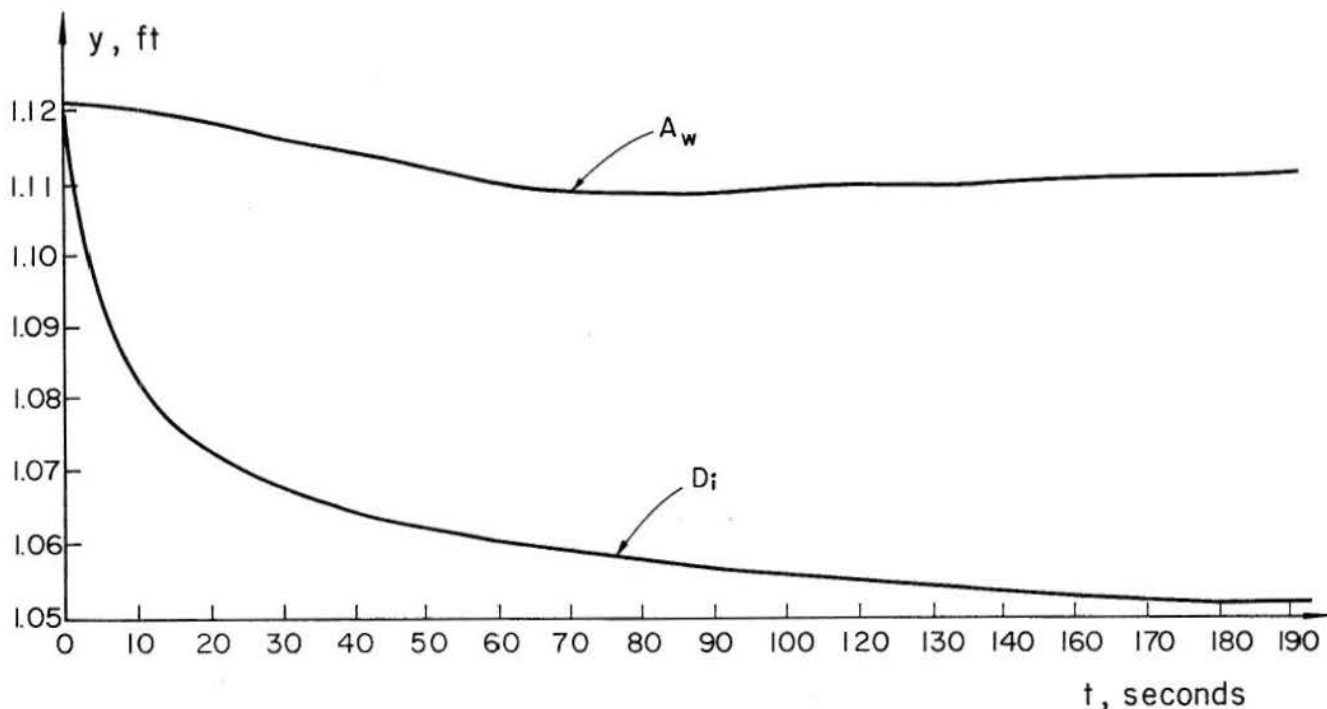


Fig. 2.2. Comparison of Lax-Wendroff scheme ( $A_w$ ) and the diffusing scheme ( $D_i$ ) in reproducing the steady initial conditions along the conduit, at the distance  $x = 796.7$  ft.

of  $y$  in feet at  $x = 735.8$  ft versus the number  $n$  of  $\Delta x$  intervals used ( $n = 80$ ,  $n = 160$ , and  $n = 320$ ) for both schemes and for three different times. The entire length of 822 ft of the conduit was divided by  $n$  to obtain the corresponding  $\Delta x$ . From top to bottom in Fig. 2.3, the given times  $t$  represent  $y$  rising (upper graph),  $y$  near maximum (central graph), and  $y$  falling (lower graph). The effects of the size of the  $\Delta x$  intervals are noticeable, and, thus, the corresponding size of  $\Delta t$  intervals are also noticeable, when comparing the diffusing scheme to the Lax-Wendroff scheme. Since the error in the Taylor series expansion is on the order of  $(\Delta t)^3$ , in which  $\Delta t$  is a function of  $\Delta x$ , the difference in  $y$  due to different  $\Delta x$  sizes is not as profound in the Lax-Wendroff scheme solutions as in the diffusing scheme. Figure 2.3 also shows the underestimation by the diffusing scheme similarly shown before in Fig. 2.2 in the study of ability of this scheme to repeat the initial condition under a constant input discharge.

The last consideration in this comparison of solutions involves the Lax-Wendroff scheme but with the assumption ( $A_{w0}$ ), or without this assumption ( $A_w$ ), that all products of first-order partial derivatives and any derivative of  $S_f$  are negligible.

Using the same hypothetical input hydrograph, Figs. 2.4 and 2.5 show plots of the depth  $y$  versus time  $t$  at positions  $x = 409.1$  ft, and  $x = 797.8$  ft, respectively. These figures give the comparisons of results for the developed Lax-Wendroff scheme ( $A_w$ ) and the simplified scheme with the above assumption ( $A_{w0}$ ). The difference occurs in the computed hydrographs when the first-order partial derivatives are

such that the assumption becomes less valid. That is for example,  $\partial y / \partial t$  is negligible only until the computed water wave reaches a particular  $x$  position and causes an increase in  $y$ .

## 2.6 Concluding Remarks

Among the finite-difference schemes, the Lax-Wendroff scheme is considered as superior not only to the diffusing scheme but to all others investigated for the purpose of flood routing through storm drains under the conditions of application of Eqs. 1.1 and 1.2. Taking into account all six schemes, either discussed briefly or analyzed, it is concluded that the Lax-Wendroff scheme is an optimal scheme between the accuracy in the results produced and the computer time necessary for the corresponding numerical solutions. It is, therefore, considered as the feasible numerical computational scheme whenever a gradually varied free-surface unsteady flow is computed directly by numerically integrating the two partial differential equations stated in Chapter 1 as Eqs. 1.1 and 1.2.

For benefit to other investigators and users, the computational procedures and programs are reproduced here in the two appendices.

Appendix 1 gives the computation details of the diffusing scheme and Appendix 2 gives the computation details of the Lax-Wendroff scheme. Each appendix contains the following items, (1) Flow chart; (2) Computer program, (3) Definition of variables; this gives the conversion table between the mathematical symbols used in this paper and the symbols used in Fortran language for a CDC 6600 or CDC 6400 digital computer; and (4) Sample input and output.

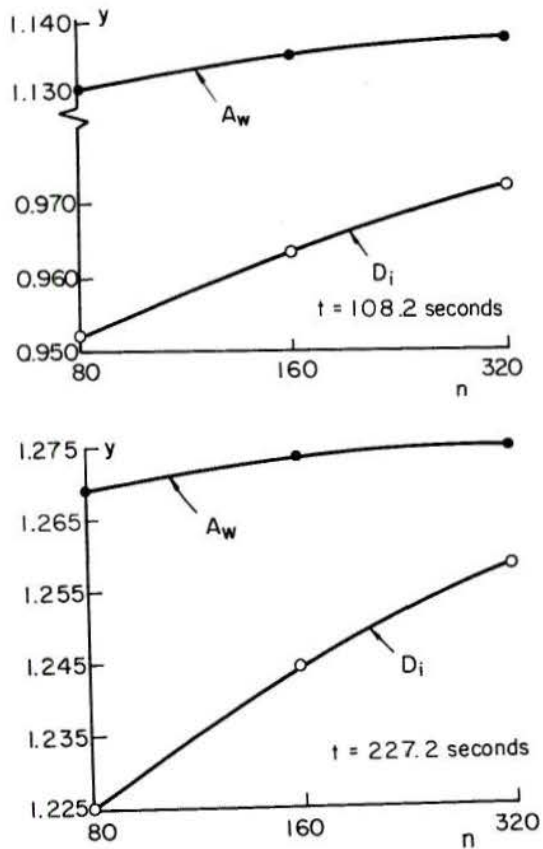


Fig. 2.3. Study of effects of the size of  $\Delta x$  and  $\Delta t$  intervals (measured by  $n$ , the number of  $\Delta x$  intervals over the length  $L = 822$  ft), on the predicted depth  $y$  at  $x = 735.8$  ft.

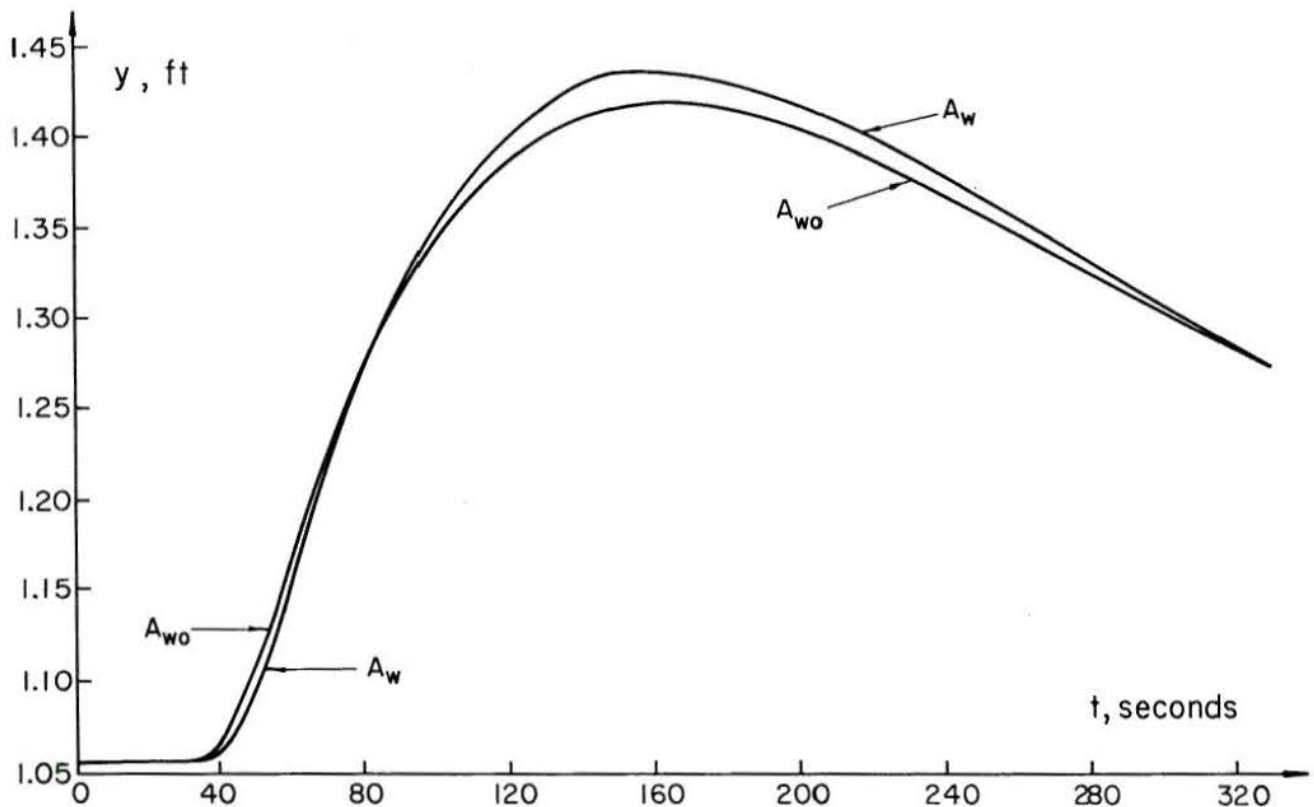


Fig. 2.4. Comparison of the hydrographs at the position  $x = 409.1$  computed with the Lax-Wendroff scheme without the assumption ( $A_w$ ) and with the assumption ( $A_{wo}$ ) of products of partial derivatives or the derivatives of  $S_f$  in  $x$  and  $t$  being negligible.

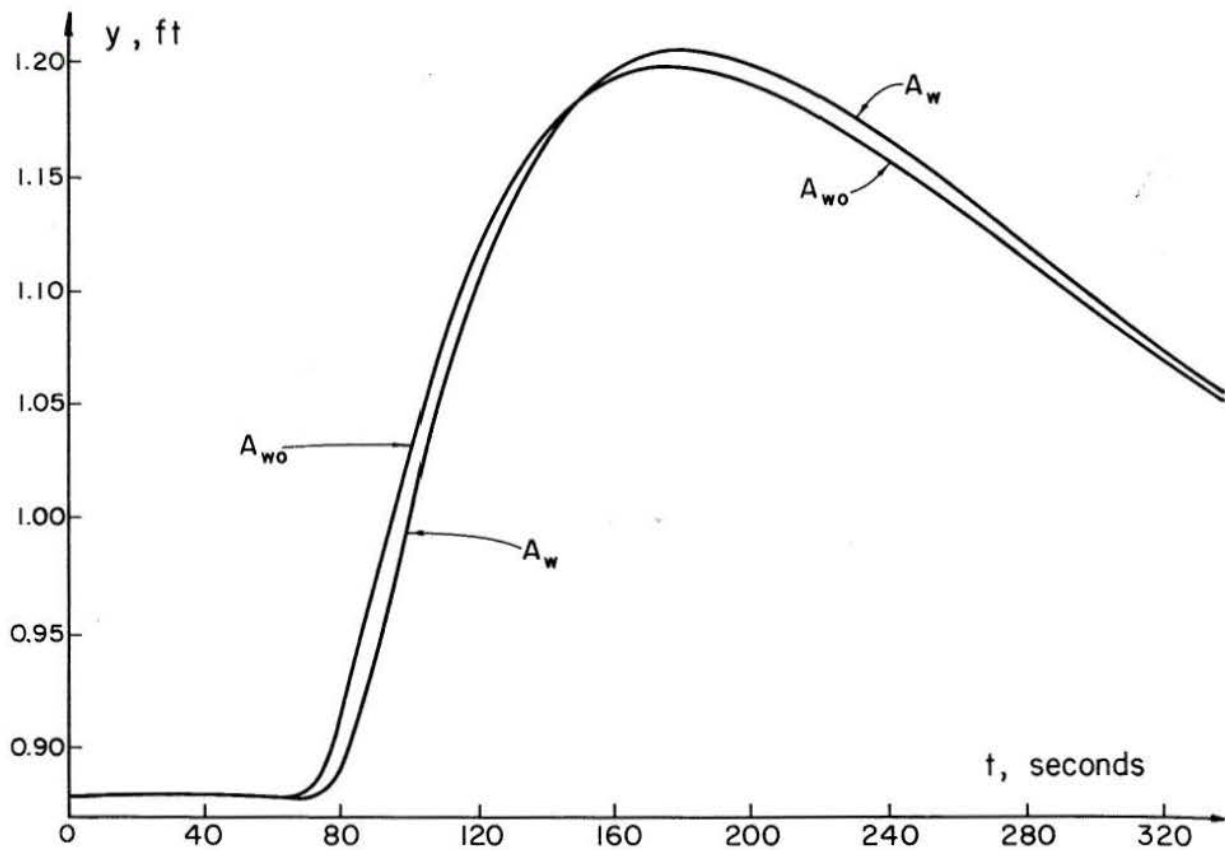


Fig. 2.5. The same comparison as in Fig. 2.4, except at the position  $x = 797.8$  ft.

## INTEGRATION OF CHARACTERISTIC DIFFERENTIAL EQUATIONS BY FINITE DIFFERENCES

## 3.1 Statement of Characteristic Equations

The two partial differential equations of gradually varied free-surface unsteady flow, Eqs. 1.1 and 1.2, when transformed give the four ordinary characteristic differential equations. Their development is shown in Chapter 3, Part I, Hydrology Paper No. 43. The equations with  $\alpha = \beta = 1$ , and  $q = 0$  (Eqs. 3.50 to 3.53 of Part I), are the starting equations and are given here as:

$$\xi_+ = \left( \frac{dt}{dx} \right)_+ = \frac{1}{V + \sqrt{gA/B}}, \quad (3.1)$$

$$\xi_- = \left( \frac{dt}{dx} \right)_- = \frac{1}{V - \sqrt{gA/B}}, \quad (3.2)$$

$$\left\{ \left( \frac{A}{VB} - \frac{V}{g} \right) \xi_+ + \frac{1}{g} \right\} \frac{dy}{dx} + \frac{A}{gVB} \frac{dV}{dx} + \frac{A}{VB} (S_o - S_f) \xi_+ = 0, \quad (3.3)$$

and

$$\left\{ \left( \frac{A}{VB} - \frac{V}{g} \right) \xi_- + \frac{1}{g} \right\} \frac{dy}{dx} + \frac{A}{gVB} \frac{dV}{dx} + \frac{A}{VB} (S_o - S_f) \xi_- = 0. \quad (3.4)$$

These four dependent equations form the basis for numerical solutions in the method of characteristics. There are a variety of procedures that may be used and these procedures may be broadly divided into two categories, the grid system and the specified intervals system.

## 3.2 Various Schemes

The first category uses the grid system generated by the intersecting characteristics curves in the time-distance plane. In this case, solutions to the problem are made at the intersections. These intersections occur at the nonuniform spacings in both  $x$  and  $t$  directions, thus, interpolations are required in order to develop time or distance relations. These relations are commonly referred to as the Lagrangian description for the distance relations at an instant of time, and the Eulerian description for the time relations at a fixed position. This method of using grids of characteristics is based on establishing the initial characteristic curves from the initial conditions. The receding characteristic curves emanate from it. In Fig. 3.1 the initial characteristic curve  $\xi_0$ , first determined from the inflow hydrograph and the initial steady conditions, is drawn from  $x = 0$  and  $t = 0$ . By introducing the values of the dependent variables  $V$  and  $y$  along the initial characteristic curve  $\xi_0$ , at the appropriate points in the computational scheme, the values of  $V$  and  $y$  as functions of the independent variables  $x$  and  $t$  are obtained at successive points. For example, the values of the depths and velocities at points  $Q_1$ ,  $Q_2$  and  $Q_3$  in Fig. 3.1 are obtained from the values of

depths, velocities, and coordinates  $(x, t)$  of the points  $Q_0$ ,  $P_1$ ,  $P_2$  and  $P_3$ , respectively. In the same manner, all values of the dependent variables  $V$  and  $y$  as functions of the independent variables  $x$  and  $t$  can be computed.

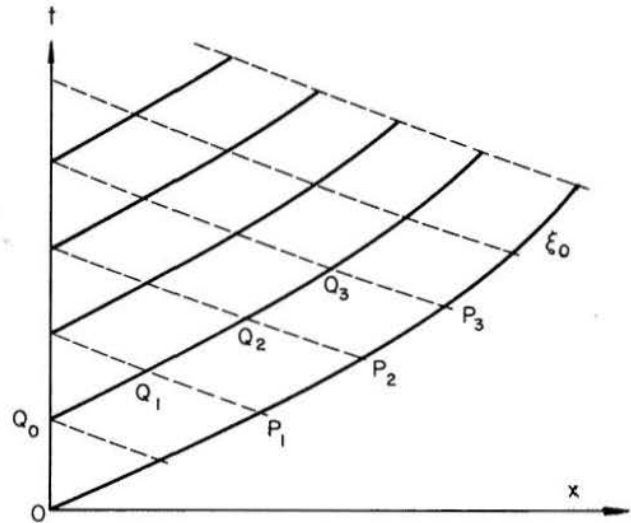


Fig. 3.1. Network of characteristics in the method of grid system for the solution of unsteady flow equations.

It is evident from the preceding brief description that the values in the solution at each intersection of characteristics must be retained in the computer for the later interpolation for fixed times and positions. No attempt was made in this study to use the method of characteristic curves. The principal reason was the need for excessive computer storage of solutions at each intersection.

The second category is the specified intervals system for independent variables. In this approach, the dependent variables  $V$  and  $y$  are known functions of the independent variables  $x$  and  $t$  either as initial conditions of  $t = 0$  or as the results of previous time computations. For example, it is assumed that  $V$  and  $y$  are known along distance  $x$  at time  $t$ . Figure 3.2 represents the rectangular grid in the  $(x, t)$ -plane with intervals  $\Delta x$  and  $\Delta t$  in  $x$  and  $t$  coordinates, respectively. In this case,  $V$  and  $y$  at points  $M_j, A_j, B_j, \dots, N_j$  are known. The values of  $V$  and  $y$  at time  $t + \Delta t$ , and particularly at points  $M_{j+1}, A_{j+1}, B_{j+1}, \dots, N_{j+1}$ , can then be computed from equations 3.1 through 3.4 and from the boundary conditions. In this manner,  $V$  and  $y$  at time  $t + \Delta t$  at various points along distance  $x$  can also be computed. This process can be continued as far as desired or meaningful. This method was selected and used in this study because the values of  $x$  and  $t$  at points  $M_{j+1}, A_{j+1}, B_{j+1}, \dots, N_{j+1}$  are exactly known, and only the values of  $V$  and  $y$  at these points must be determined.

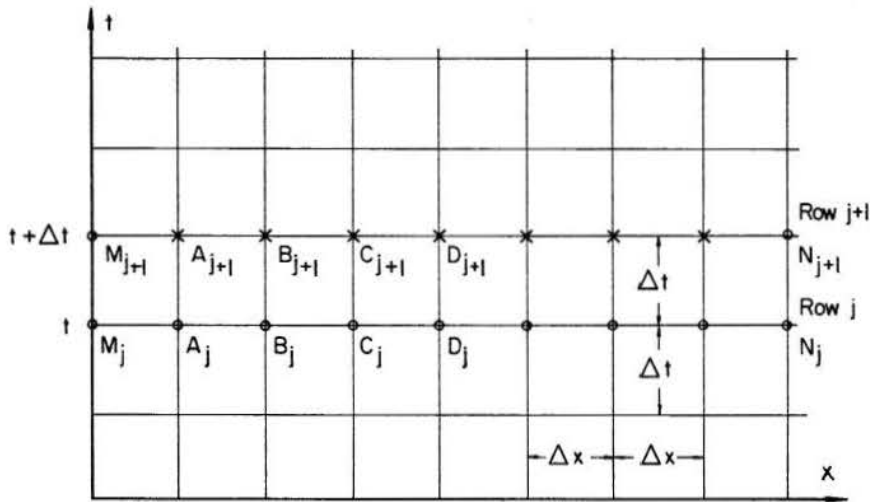


Fig. 3.2. Network of specified intervals for the solution of characteristic equations.

This method has the advantage that it gives results directly and in the form most needed and useable, such as the hydrograph at each position along the channel and also the water surface profile at any given time. From the view of computer programming, arrangement of the steps of computation for the methods of the second category appears to offer advantages over the methods of the first category. Since the values of the dependent variables at time  $t$  in the second category are known at predetermined points, the only information needed to be stored in the computer is the values of the dependent variables at time  $t + \Delta t$ . Therefore, this category needs computer storage of only two time lines as indicated in Fig. 3.2 and designated by  $j$  and  $j+1$  rows, respectively. Values of the dependent variables  $V$  and  $y$  of row  $j$  are known and stored while the values of  $V$  and  $y$  of row  $j+1$  are being computed for the next time interval. After completion of this time interval, the values of  $V$  and  $y$  of row  $j+1$  are stored for computation at the next time interval; the values of  $V$  and  $y$  of row  $j$  are then printed out and the storage space is replaced by the values of row  $j+1$ .

### 3.3 Numerical Solution by the Specified Intervals System

This section discusses the numerical solution of the equations of free-surface unsteady flow by the method of characteristics with the specified time interval,  $\Delta t$ , and the specified distance interval  $\Delta x$ . In this method,  $V$  and  $y$  at point  $P$  on the  $(x, t)$ -plane of Fig. 3.3 are to be computed from the initial conditions or from previous values of  $V$  and  $y$  at points  $A$ ,  $B$ , and  $C$  using two assumptions:

(a)  $\Delta t$  is sufficiently small so that the parts of the characteristics between  $P$  and  $R$  and between  $P$  and  $S$  may be considered as straight lines, and

(b) The slope of the straight line  $PR$  at point  $P$  is the positive characteristic direction of the position  $C$ ,  $(\xi_+)_C$ , and the slope of the straight line  $PS$  at point  $P$  is the negative characteristic direction of the position  $C$ ,  $(\xi_-)$ .

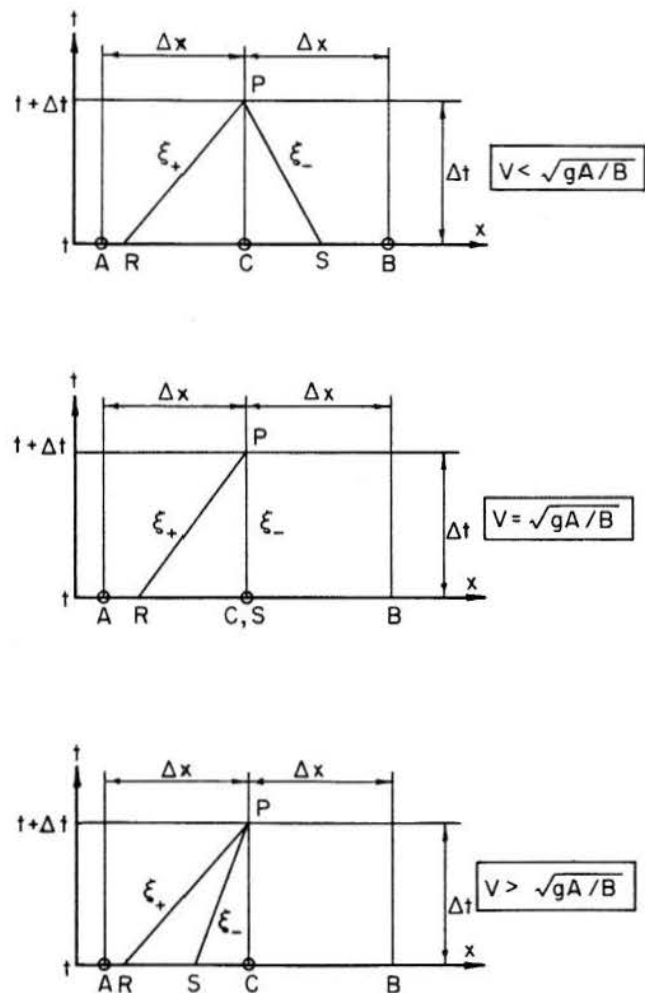


Fig. 3.3. Rectangular grid for the solution by the system of specified intervals,  $\Delta t$  and  $\Delta x$ : subcritical flow (upper graph), critical flow (center graph), and supercritical flow (lower graph).

Since  $x_p$  and  $t_p$  are known, the velocity at point P,  $V_p$ , and the depth at point P,  $y_p$ , are then computed. The computations proceed as follows.

(1) The coordinates of R and S are determined from the relations of  $(\xi_+)_C$ ,  $(\xi_-)_C$ , and the geometry of the grid by

$$t_p - t_R = (\xi_+)_C (x_p - x_R), \quad (3.5)$$

and

$$t_p - t_S = (\xi_-)_C (x_p - x_S), \quad (3.6)$$

in which  $(\xi_+)_C$  and  $(\xi_-)_C$  are computed from Eqs. 3.1 and 3.2, respectively, at point C.

(2) The values of  $V_R$ ,  $V_S$ ,  $y_R$ , and  $y_S$  are determined by interpolation from the Taylor expansion, with  $h$  the symbol of either  $\Delta x$  or  $\Delta h$ , as

$$f(x+h) = f(x) + hf'(x) + \frac{h^2}{2!} f''(x) + \dots + 0(h^n), \quad (3.7)$$

and

$$f(x-h) = f(x) - hf'(x) + \frac{h^2}{2!} f''(x) + \dots + 0(h^n), \quad (3.8)$$

For a first order interpolation, the second and higher derivatives are neglected. The first derivative of Eq. 3.7 becomes, in finite difference form,

$$f'(x) = \frac{f(x + \Delta x) - f(x)}{\Delta x},$$

and that of Eq. 3.8 becomes, in finite-difference form,

$$f'(x) = \frac{f(x) - f(x - \Delta x)}{\Delta x}.$$

The value of the function ( $U = V$  or  $y$ ) at points R and S are then, from Eq. 3.8 and Eq. 3.7, respectively,

$$U_R = U_C - \frac{U_C - U_A}{\Delta x} (x_C - x_R) \quad (3.9)$$

$$U_S = U_C + \frac{U_C - U_B}{\Delta x} (x_C - x_S) \quad (3.10)$$

For the second order interpolation, the third and higher derivatives of Eq. 3.7 and Eq. 3.8 are neglected, the first and second derivatives in these two equations become, in finite-difference form,

$$f'(x) = \frac{f(x+\Delta x) - f(x-\Delta x)}{2\Delta x}$$

and

$$f''(x) = \frac{f(x+\Delta x) - 2f(x) + f(x-\Delta x)}{(\Delta x)^2}$$

The value of the function ( $U = V$  or  $y$ ) at points R and S are then

$$U_R = U_C - \frac{U_B - U_A}{2\Delta x} (x_C - x_R) + \frac{U_B - 2U_C + U_A}{2(\Delta x)^2} (x_C - x_R)^2, \quad (3.11)$$

$$U_S = U_C - \frac{U_B - U_A}{2\Delta x} (x_C - x_S) + \frac{U_B - 2U_C + U_A}{2(\Delta x)^2} (x_C - x_S)^2, \quad (3.12)$$

from which  $V_R$ ,  $V_S$ ,  $y_R$  and  $y_S$  may be computed knowing the  $V$  and  $y$  at points A, C, and B.

(3) Then  $V_p$  and  $y_p$  are obtained by solving simultaneously the finite-difference forms of Eqs. 3.3 and 3.4, or by

$$(F_+)_C (y_p - y_R) + (G_+)_C (V_p - V_R) + (S_+)_C (x_p - x_R) = 0 \quad (3.13)$$

and

$$(F_-)_C (y_p - y_S) + (G_-)_C (V_p - V_S) + (S_-)_C (x_p - x_S) = 0 \quad (3.14)$$

in which the above values of  $F$ ,  $G$ , and  $S$  at point C are defined as

$$(F_+)_C = (A_1 C_2 - A_2 C_1)_C (\xi_+)_C - (B_1 C_2 - B_2 C_1)_C;$$

$$(G_+)_C = (A_1 B_2 - A_2 B_1)_C;$$

$$(S_+)_C = (A_1 E_2 - A_2 E_1)_C (\xi_+)_C - (B_1 A_2 - B_2 A_1)_C;$$

$$(F_-)_C = (A_1 C_2 - A_2 C_1)_C (\xi_-)_C - (B_1 C_2 - B_2 C_1)_C;$$

$$(G_-)_C = (A_1 B_2 - A_2 B_1)_C, \quad \text{and}$$

$$(S_-)_C = (A_1 E_2 - A_2 E_1)_C (\xi_-)_C - (B_1 E_2 - B_2 E_1)_C,$$

in which the above coefficients of the two general partial differential equations (Eqs. 3.24 and 3.25, Part I, Hydrology Paper No. 43) are:  $A_1 = A/VB$ ,  $A_2 = V/g$ ,  $B_1 = 0$ ,  $B_2 = 1/g$ ,  $C_1 = C_2 = 1$ ,  $D_1 = 1/V$ ,  $D_2 = 0$ ,  $E_1 = 0$ , and  $E_2 = S_f - S_o$ . Solving equations 3.13 and 3.14 simultaneously,

$$y_p = \frac{\begin{vmatrix} (T_+)_C & (G_+)_C \\ (T_-)_C & (G_-)_C \end{vmatrix}}{\begin{vmatrix} (F_+)_C & (G_+)_C \\ (F_-)_C & (G_-)_C \end{vmatrix}} \quad (3.15)$$

and

$$V_p = \frac{\begin{vmatrix} (F_+)_C & (T_+)_C \\ (F_-)_C & (T_-)_C \end{vmatrix}}{\begin{vmatrix} (F_+)_C & (G_+)_C \\ (F_-)_C & (G_-)_C \end{vmatrix}} \quad (3.16)$$

in which

$$(T_+)_C = (F_+)_C y_R + (G_+)_C V_R - (S_+)_C (x_p - x_R), \quad (3.17)$$

and

$$(T_-)_C = (F_-)_C y_S + (G_-)_C V_S - (S_-)_C (x_p - x_S). \quad (3.18)$$

By these computations, velocities and depths at time  $t + \Delta t$  are obtained for all points along the channel, except for the two boundary points. The values for the boundary points are provided by previous computations of the known boundary conditions.

The procedure in the solution requires first the determination of the intervals within which the points R and S lie. A linear interpolation is then performed within the appropriate interval for the dependent variables at time  $t$ . This linear interpolation has the same effect as the linear interpolation in the diffusing finite-difference scheme, namely a systematic positive or negative shift in the computed values  $V$  and  $y$ .

In an attempt to eliminate this deficiency, a second-order interpolation was developed. Referring again to Fig. 3.3 (upper graph), a second-degree polynomial of the form

$$U = a + bx + cx^2 \quad (3.19)$$

is assumed to fit the function of  $V$  and  $y$  through points A, C, and B. This is the same interpolation as in Eqs. 3.9 and 3.10, except in a different way of implementing it. If the function is centered on the location of C, then the constants are

$$a = U_C, \quad b = \frac{U_B - U_A}{2\Delta x}, \quad \text{and} \quad c = \frac{U_B - 2U_C + U_A}{2\Delta x^2}. \quad (3.20)$$

Thus, the value of the function of the location of R is

$$U_R = U_C - \frac{1}{2}(UP)(U_B - U_A) + \frac{1}{2}(UP)^2(U_B - 2U_C + U_A) \quad (3.21)$$

in which

$$UP = - \frac{\Delta t}{\Delta x} \left/ \left( \frac{dt}{dx} \right)_+ \right. \quad (3.22)$$

The ratio of  $\Delta t$  to  $\Delta x$  is the selected grid mesh ratio and  $(dt/dx)_+$  is the direction of the positive characteristic estimated from the conditions at location C.

Similarly, the value of the function at location S is

$$U_S = U_C - \frac{1}{2}(UN)(U_B - U_A) + \frac{1}{2}(UN)^2(U_B - 2U_C + U_A) \quad (3.23)$$

in which

$$UN = - \frac{\Delta t}{\Delta x} \left/ \left( \frac{dt}{dx} \right)_- \right. \quad (3.24)$$

This interpolation scheme offers two advantages. First, the curvature of the function at a given time is approximated. Second, it is not necessary to compute within which interval the intersection of the characteristic and the  $x$ -axis falls. The assumptions

in this scheme are that the functions of velocity and depth are continuous and may be approximated by a parabolic relation within the interval. Any other similar non-linear interpolation scheme may be designed if it suits the general types of the  $V(x)$  and  $y(x)$  functions for various values of  $t$ .

### 3.4 Initial Conditions

The necessary initial conditions for the unsteady free-surface flow are that all velocities and depths of water along the channel must be known at a given time. In this study, it was assumed that at the initial time the discharge was constant throughout the reach. Thus, the problem can be treated as a steady non-uniform flow. Velocities and depths along the channel were then determined by computations of conventional backwater or drawdown surface profiles, depending on the downstream control conditions. This procedure uses the standard step method [2, p. 265].

### 3.5 Boundary Conditions

The two governing partial differential equations for unsteady flow require two independent boundary conditions relating velocity and depth at certain locations along the channel. One of these conditions is the discharge-time relation existing at the inlet end to the section of channel under study. This relation can be either expressed in a mathematical form, or given as discrete points of discharge at selected intervals of time.

The other boundary condition imposed on the problem is that of a discharge-versus-depth relation at the downstream end, characterized either by a control structure or by the critical depth at a free outfall. This is the boundary condition that must exist for subcritical flow of the base discharge.

If the base discharge is in the supercritical range or on a supercritical slope the boundary condition must be expressed at the inlet end. This function takes the form of a discharge-versus-depth relation. This condition, in combination with the condition of a discharge-versus-time relation, is somewhat difficult to visualize physically; however, it is a necessary condition because the characteristic directions both have a positive slope and thus there is no influence of the downstream conditions on the upstream conditions.

The following discussion presents a detailed analysis of these boundary conditions. Arbitrary inflow hydrographs were investigated to test and verify the computer program and also to provide results for evaluating the significance of variations in the hydraulic parameters.

Upstream boundary conditions - The boundary condition at the upstream inlet is given by an inflow hydrograph,  $Q(t)$ , with no limitation on the shape of the hydrograph. A hypothetical hydrograph, having a Pearson Type III distribution with four parameters, was selected for evaluating the effect of variations in the parameter and is shown by Fig. 3.4. Thus, the inflow  $Q$  at time  $t$  designated by  $Q(t)$  may be described by

$$Q(t) = Q_b + Q_o e^{-\frac{(t-t_p)}{(t_g-t_p)}} \frac{t}{(t/t_p)^{t/(t_g-t_p)}}, \quad (3.25)$$

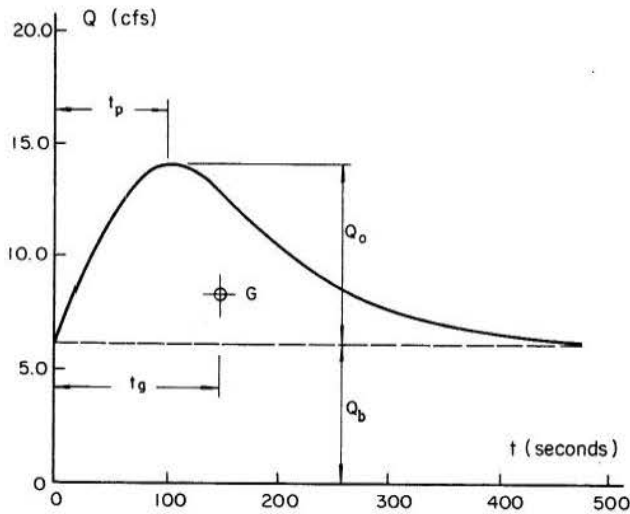


Fig. 3.4. Hypothetical inflow hydrograph of the Pearson Type III function, Eq. 3.25, with the selected parameters:  $Q_b = 6.21$  cfs,  $Q_o = 8.00$  cfs,  $t_p = 100.00$  sec, and  $t_g = 150.0$  sec.

in which  $Q_b$  is the constant base flow,  $Q_o$  is the peak flow,  $t_p$  is the time from the beginning of storm runoff to peak discharge and  $t_g$  is the time from the beginning of the storm runoff to the center of mass of storm runoff,  $G$ . One hydrograph with arbitrary values of  $Q_b$ ,  $Q_o$ ,  $t_p$ , and  $t_g$  were used in this study. The shape and these arbitrary values of parameters are shown in Fig. 3.4.

The depth and the velocity at the upstream boundary point  $P$  in Fig. 3.5, which is at  $x = 0$  and at the time  $t + \Delta t$ , can be computed from initial conditions at  $C$  and  $B$ , with the boundary conditions given by the inflow hydrograph

$$AV = Q(t) \quad (3.26)$$

in which  $A$  is the cross-sectional area and  $V$  is the velocity at  $P$ .

Using the previously discussed assumptions and procedure of computing velocities and depths at other points along the channel the negative characteristic direction at point  $C$  is also given by the initial conditions. The relation between the depth  $y_p$  and velocity  $V_p$  at point  $P$  can be determined from Eq. 3.4. Substituting the boundary condition of Eq. 3.26 into Eq. 3.14 gives

$$y_p = y_s - \frac{(G_-)_C \left\{ \frac{Q(t)}{A} - V_s \right\} + (S_-)_C (x_p - x_s)}{(F_-)_C} \quad (3.27)$$

in which  $A$  is the cross-sectional area at  $P$  and  $A$  is a function of  $y_p$ .

Solving for  $y_p$  from Eq. 3.27 and substituting  $y_p$  into Eq. 3.26 makes it possible to determine  $V_p$ . Since Eq. 3.27 is not linear in  $y_p$ , a Newton-Raphson iteration was used for its solution.

Downstream boundary conditions - The boundary conditions at the downstream outlet may generally be

given by a stage-discharge relation. In this portion of the study only a free outfall at the end of conduit was assumed. Therefore, a critical flow at the downstream end exists

$$\frac{V}{\sqrt{g \frac{A}{B}}} = 1 \quad (3.28)$$

where  $A$  is the cross-sectional area and  $B$  is the top width of the downstream boundary.

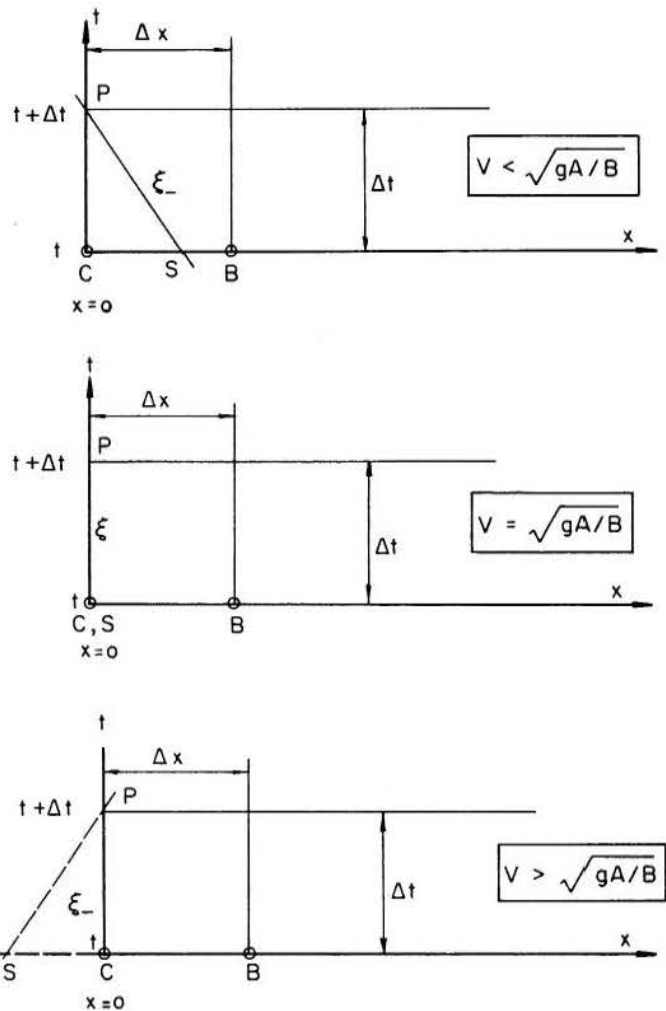


Fig. 3.5. Upstream boundary conditions: subcritical flow (upper graph), critical flow (central graph), and supercritical flow (lower graph).

Figure 3.6 shows the downstream boundary where the critical depth occurs. For the free outfall, it was assumed that critical depth occurred at a distance of 4.5 times the critical depth from the end. This assumption was also applied to the unsteady case, with critical depth computed from the base discharge,  $Q_b$ . Therefore, the total distance  $x_L$  from the inlet to the downstream boundary is determined by



$$x_L = x_F - 4.5 y_c \quad (3.29)$$

in which  $x_F$  is the total length of the channel and  $y_c$  is the critical depth for discharge  $Q_b$ .

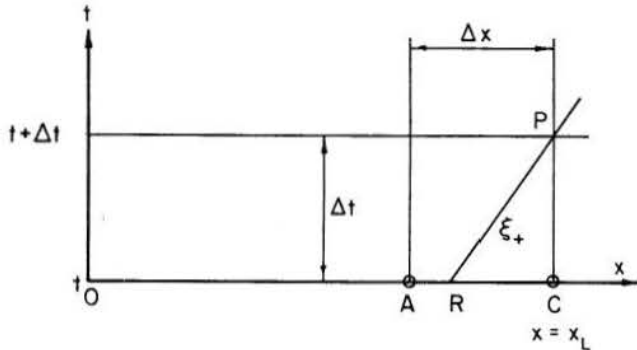


Fig. 3.6. Downstream boundary conditions for the subcritical flow, with  $x_L$  the computational conduit length.

The depth and velocity at the downstream boundary point P at time  $t + \Delta t$  can be computed from the initial conditions at A and C, and from the boundary conditions given by Eq. 3.28.

Using the same assumptions and computational procedures, the initial conditions also give the relation between the depth  $y_p$  and the velocity  $V_p$  by applying Eq. 3.3. Substituting the boundary conditions of Eq. 3.28 into Eq. 3.13 results in

$$y_p = y_R - \frac{(G_+)_C(\sqrt{gA/B} - V_R) + (S_+)_C(x_p - x_R)}{(F_+)_C} \quad (3.30)$$

in which A is the cross-sectional area and B is the top width at P, with both A and B functions of  $y_p$ .

Solving  $y_p$  from Eq. 3.30 and substituting  $y_p$  into Eq. 3.16 makes it possible to determine  $V_p$ . Since Eq. 3.30 is not linear in  $y_p$ , a Newton-Raphson iteration was again used for a solution.

### 3.6 Summary of Computational Procedures

In solving the equations of free-surface unsteady flow, Eqs. 1.1 and 1.2 and Eqs. 3.1 and 3.4, by the system of specified intervals, the steps of computing velocity  $V$  and depth  $y$  at various times and positions along the conduit are as follows.

(1) Values of  $V$  and  $y$  at various positions along the channel for the steady-state condition of constant base flow,  $Q_b$ , are determined from a computation of the backwater curve.

(2) The upstream boundary conditions are evaluated.

(3) The downstream boundary conditions are evaluated.

(4) Values of  $V$  and  $y$  at time  $t + \Delta t$  along the channel are computed from the known values of  $V$  and  $y$  at time  $t$ .

(5) Steps (2), (3), and (4) are repeated as long as desired or meaningful.

To benefit other investigators, the computational procedures and programs are reproduced in Appendix 3. Appendix 3 gives the computation details of the numerical integration method using the specified interval scheme of the method of characteristics. It includes (1) flow chart, (2) computer program, (3) definitions of variables and (4) sample input and output. Additional subroutines were developed to compute the boundary conditions for supercritical regime and for lateral inflow at specified locations.\*

### 3.7 Effect of Variations in Computational Parameters

The discrepancy between a computed value and the observed value from a physical experiment is attributable to numerous sources of errors. These errors are generally the result of systematic and random errors in the observational system and possible systematic errors in computational procedures. Random errors are a result of unavoidable accidental variations in the physical systems. The discussion that follows will be concerned with errors in the computational procedure.

Computational errors emanating from procedures in this study are the result of:

(1) The approximation of infinitesimal variations by finite values. This is a result of assuming in general, linear relations rather than the true curvilinear relations. This is a systematic error. However, the propagation of this error is not readily determined since it may be positive or negative during different stages of the computations.

(2) Truncation of numerical values. This is due to the limited precision of any discrete-element calculator.

(3) Round off in the printed output. The printed output of any computed value from a digital computer differs from the internally generated value by the amount the value is rounded off in conversion to numeric form. The computer used for these calculations rounds off in a manner similar to manual calculators.

The following discussion evaluates the significance of the controllable variables in the solution of the unsteady flow equations. These equations are considered under the computational parameters of incremental length and incremental time interval during which the integration process proceeds.

The effect of variations in the hydraulic parameters of roughness and the velocity distribution coefficients is discussed in Part I, Hydrology Paper No. 43.

Determination of computational parameter  $\Delta t$ . The grid sizes of  $\Delta x$  and  $\Delta t$  in the computational scheme, Fig. 3.2, is limited by the characteristic directions  $\xi_+$ ,  $\xi_-$ , encountered during the integration.

Referring to Fig. 3.3, in order for R to lie in the interval A-C for all conditions of flow, it is necessary that the ratio of  $\Delta t/\Delta x$  be less than the value of  $dt/dx$  assumed at the location R. This condition must exist throughout the integration solution.

In order to assure that this condition exists, it is necessary that  $\Delta t$  be computed from

$$\Delta t = \Delta x / [V + \sqrt{gA/B}]$$

\* Originals of all computer-program and punched-card decks are deposited with the Office of Research, Federal Highway Administration, U.S. Department of Transportation, Washington, D.C.

in which

- (1)  $V$  is the maximum anticipated velocity, and
- (2)  $A/B$  is a maximum for free surface flow.

Effect of computational parameter  $\Delta x$ . The method of characteristics using a specified intervals system gives the complete numerical solution of the free-surface unsteady flow. The accuracy of the results depends on the size of the rectangular grids  $\Delta x$  and  $\Delta t$  of Fig. 3.2. In this section only the effect of  $\Delta x$  is discussed;  $\Delta t$  will be discussed in the next section.

If  $n$  is the number of intervals along the conduit and  $x_L$  is the length of the conduit, then

$$\Delta x = \frac{x_L}{n} \quad (3.32)$$

Since  $x_L$  is assumed to be fixed,  $n$  is arbitrarily selected as any even number, thus  $\Delta x$  is determined. The smaller the  $\Delta x$ , presumably the more accurate are the results. But also, the smaller the  $\Delta x$ , the greater the required computing time. In compromising these two conditions to satisfy the objectives of this study, several values of  $n$  for the fixed  $x_L$  were tried.

Figure 3.7 shows the effect of the size of  $\Delta x$  on the depth hydrographs at three positions along the conduit. The upper graph is the depth hydrograph at a position 50.0 feet downstream from the inlet and for a  $\Delta x$  of 40.91, 20.45, 10.23, and 5.12 feet corresponding to  $n$  values of 20, 40, 80, and 160, respectively. The center and lower graphs are the depth hydrographs at 410.0 feet from the inlet, and 771.7 feet from the inlet, respectively. The initial condition for each computation is the steady-state water surface for a free outfall.

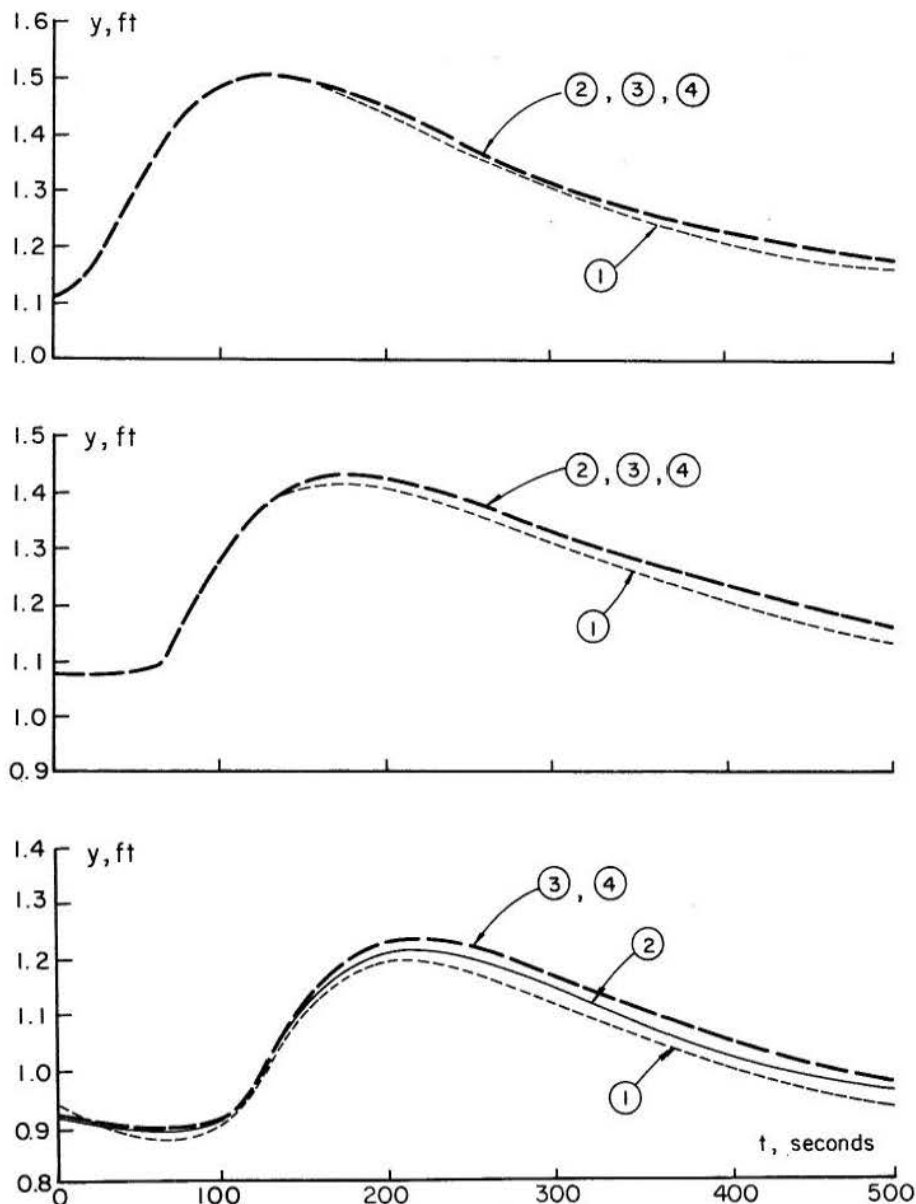


Fig. 3.7. Effect of  $\Delta x$  on hydrographs at various positions along the conduit; (1)  $\Delta x = 40.91$  ft, (2)  $\Delta x = 20.45$  ft, (3)  $\Delta x = 10.23$  ft, and (4)  $\Delta x = 5.12$  ft, at three locations of conduit  $x = 50.0$  ft (upper graph),  $x = 410.0$  ft (center graph) and  $x = 771.7$  ft (lower graph).

Comparing the depth hydrographs of Fig. 3.7 with the given inflow discharge hydrograph of Fig. 3.4, it was found that:

(1) The critical portion of the conduit for computing depth hydrographs is near the outlet where there is the greatest curvature of the water surface profile. The maximum differences between the computed depths, with  $\Delta x$  being 40.91 and 5.12 feet, are approximately 0.3, 0.6, and 1.0 percent of the conduit diameter at 50.0, 410.0, and 771.7 feet from the inlet, respectively.

(2) There is no significant increase in accuracy over 0.005 feet or 0.15 percent of the conduit diameter when  $\Delta x$  is less than 10.23 feet. Therefore, a  $\Delta x$  equal to 10.23 feet, or  $n$  equal to 80, was selected for computation in the other portions of this study.

The peak depth  $y_p$  and the time to peak depth  $T_p$  are two important parameters describing a depth hydrograph. These two parameters are defined and shown graphically in Fig. 3.8. The required accuracy of a computed hydrograph at various positions along the conduit can be measured by the peak depth,  $y_p$ , relative to the diameter,  $D$  of the conduit, for various lengths  $\Delta x$ . Also, the accuracy can be measured by the time to peak depth,  $T_p$ , relative to the time to peak discharge,  $t_p$ , of the inflow discharge hydrograph of Fig. 3.4, for various lengths  $\Delta x$  and the same positions,  $x$ .

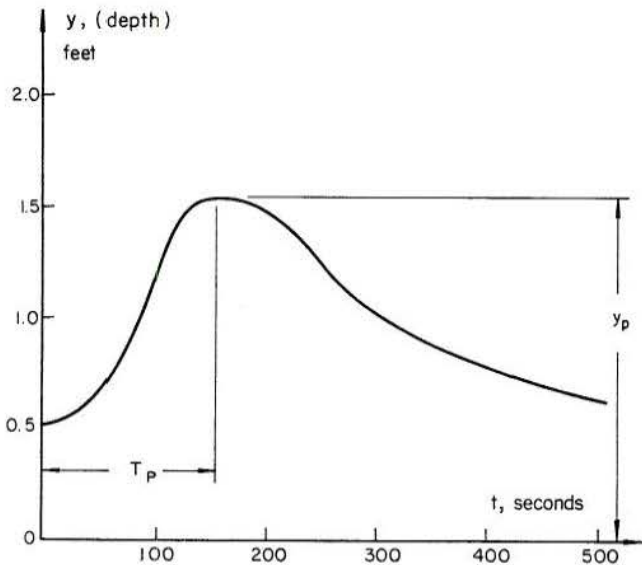


Fig. 3.8. Characteristics of the depth hydrograph with  $T_p$  the time at peak depth, and  $y_p$  the peak depth.

From the selected criteria for defining the accuracy of a computed hydrograph for a given  $\Delta x$ , it was found that the percentage differences of  $y_p$ ,

$$\frac{(y_p)_i - (y_p)_{\min}}{D} \times 100$$

in which the index "min" refers to the depth  $y_p$  of the smallest difference used,  $\Delta x = 5.12$  ft, and the index "i" refers to depths of any other  $\Delta x > 5.12$  ft, ranged from 0.0 percent to 2.1 percent for  $\Delta x$  ranging from 5.12 ft to 40.91 ft, and at various positions  $x$ , as shown in Table 3.1. At the upstream part of the conduit there was no significant difference between  $y_p/D$  measure for different values of  $\Delta x$ , as expected. At the approximate middle of the conduit there was a 0.2 percent difference. At the downstream end, the difference was 2.1 percent. No significant change in the percentage difference of  $y_p$  to  $D$  was found when  $\Delta x$  was reduced below 10.23 ft.

In using the other parameter,  $T_p$ , to define the accuracy of computed depth hydrographs with different values of  $\Delta x$  and various positions  $x$ , the measure of accuracy was

$$\frac{(T_p)_i - (T_p)_{\min}}{t_p} \times 100$$

in which the indices "min" and "i" refer to the  $\Delta x = 5.12$  ft and all others  $\Delta x$ , respectively. It was found that there were no significant percentage differences for values  $\Delta x > 5.12$  ft, and various positions  $x$ . The percentages were about 1.2 percent at the upstream, 2.0 percent at the middle, and 8.5 percent at the downstream part of the conduit. It was also found that there was no significant change of the percentages of  $T_p$  to  $t_p$  (which was about 1.9 percent) when  $\Delta x$  was reduced below 10.23 ft, as shown in Table 3.2.

Tables 3.1 and 3.2 show the percentage differences of  $y_p$  to the diameter  $D$  of the conduit, and  $T_p$  to  $t_p$ , respectively, with different values of  $\Delta x$  and various positions,  $x$ . These values at even distances (0, 50, 100, ...ft) were computed by linear interpolation from the values in the grid system of Fig. 3.2; therefore, some error may have been introduced. However, the change in shape of the depth hydrograph due to varying  $\Delta x$  was considered to be small. Larger  $\Delta x$  produced a lower and later peak depth.

As previously mentioned, the smaller the  $\Delta x$ , the longer the computing time required. For these particular values in the hydrograph and the specified grid system computer program, the relation between the time required for the CDC 6600 computer and the various  $\Delta x$  or  $n$  values is shown in Fig. 3.9. This relation is approximately a power function because the number of computational locations in the  $(x, t)$ -plane is proportional to the square of the  $x$ -positions for a constant time position.

Table 3.1. Difference in  $y_p$  computed from various sizes of  $\Delta x$   
(in percent of conduit diameter  $D$ )

$\Delta x$	DISTANCE, ft																
(ft)	0	50	100	150	200	250	300	350	400	450	500	550	600	650	700	750	800
40.91	0	-0.02	-0.16	-0.04	-0.06	-0.08	-0.11	-0.16	-0.24	-0.31	-0.41	-0.50	-0.59	-0.70	-0.94	-1.43	-2.07
20.45	0	-0.01	-0.02	-0.02	-0.03	-0.04	-0.04	-0.06	-0.10	-0.13	-0.18	-0.22	-0.27	-0.39	-0.42	-0.66	-0.99
10.23	0	0	-0.01	0	-0.01	-0.01	-0.01	-0.02	-0.03	-0.04	-0.06	-0.08	-0.09	-0.11	-0.14	-0.23	-0.39

Table 3.2. Difference in  $T_p$  computed from various sizes of  $\Delta x$   
(in percent of  $t_p$ )

$\Delta x$	DISTANCE, ft																
(ft)	0	50	100	150	200	250	300	350	400	450	500	550	600	650	700	750	800
40.91	1.23	-0.09	0.18	0.14	-1.21	-0.36	-1.62	-2.04	-2.02	-1.81	-1.09	1.21	-0.96	-1.43	-8.47	-7.32	-3.48
20.45	-0.40	-0.09	0	0.14	0.05	-0.06	0	-0.40	-0.40	-1.81	-2.73	-0.42	-0.40	0	-3.58	-4.07	-2.04
10.23	0.41	0	0	0.14	0.05	0	0	-0.22	-0.40	0	-1.90	-0.24	-0.42	0	-1.49	-1.62	-0.41

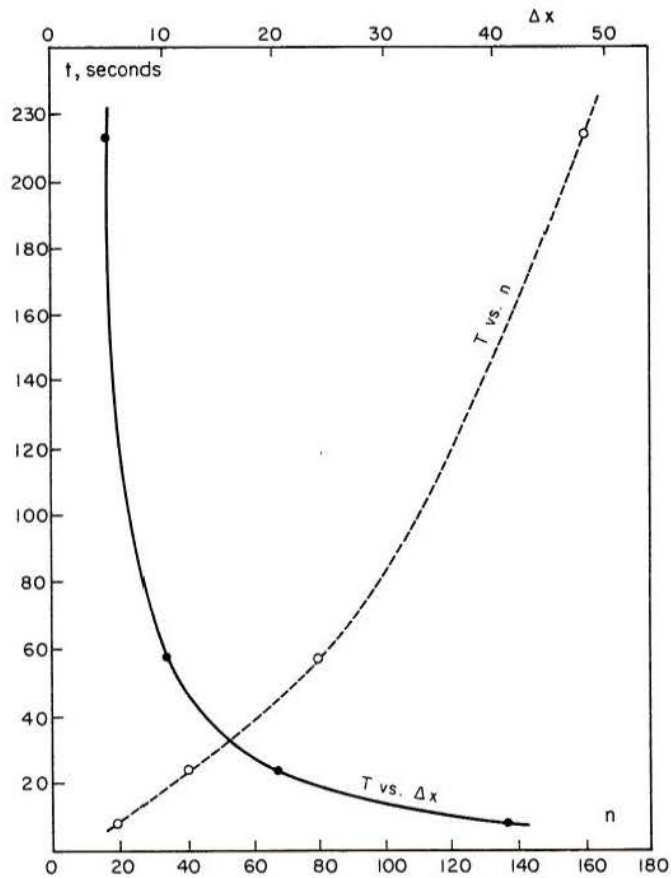


Fig. 3.9. Relations between  $n$  and  $\Delta x$  and the computer time,  $T$ , required for CDC 6600 computer.

COMPARISON OF THREE FINITE DIFFERENCE SCHEMES  
OF NUMERICAL INTEGRATION

4.1. Criteria for Comparison

The comparison of three finite-difference schemes for numerical integration and numerical computer solution and the eventual selection of the most desirable scheme for particular applications depend on simplicity, stability, accuracy, flexibility, and resulting computer time. The three schemes to be compared are: diffusing, Lax-Wendroff, and specified intervals scheme in the method of characteristics.

The simplicity of a particular scheme is related to both the algebraic description of its numerical algorithm and the computer programming involved. Generally, if the algebra is kept simple for understanding the computer programming is usually also simplified. Frequently, however, this may lead to numerous programming decisions to insure that conditions outside the range of the simplified assumptions are either included or deliberately excluded. Thus, simplified algebra does not necessarily infer simplicity in the computer algorithm.

The stability of a solution infers that the process will converge to a real solution. This criterion is satisfied in the case of solving the De Saint Venant equations if the mesh size  $\Delta t/\Delta x$  ratio is less than  $dt/dx$ , for any part of the  $(x,t)$ -plane used in the integration solutions. If this condition is not satisfied, the solution will fluctuate about the correct value with increasing amplitude. Eventually, the results may exceed the capacity of computer.

The accuracy of a solution method in this study infers that the algorithm will reproduce the initial conditions for the steady state boundary conditions. As a corollary, the algorithm should be able to compute the steady state conditions from any arbitrary initial conditions. If the algorithm satisfies this criterion, it may be inferred that there will be good agreement between the computed and the observed quantities. The difference between these two can then be attributed to the limitations of the underlying assumptions of the theoretical equations and the limitations of accurately estimating the geometric and hydraulic parameters.

The flexibility of a computer algorithm depends on the range of conditions the algorithm will accommodate. For the unsteady flow solutions, it is desirable that the algorithm provide for all conditions of depth, velocity, and discharge within the expected physical ranges. Generally, this must include both the subcritical and the supercritical conditions. Since numerical procedures at some stage require interpolations, a computer decision is required to determine the appropriate interpolation.

4.2 Properties of Diffusing Scheme

The diffusing scheme is the simplest of the three compared schemes to develop and represent in algebraic form. This can be seen from Table 2.1, wherein the partial derivatives are represented as ratios of finite differences. This simplicity, of algebraic form, however, limits accuracy and flexibility.

The stability of the diffusing scheme is assured provided the ratio of  $\Delta t/\Delta x$  does not exceed the

absolute maximum value of  $dt/dx$  at any point in the  $(x, t)$ -plane during the integration process.

The accuracy of the scheme may suffer during eventual periods of supercritical flow. This is because the characteristics intersect at a relatively great distance from the solution point. Figure 4.1 graphically presents this relationship. The accuracy of the diffusing scheme is further limited because the dependent variables are assumed to vary linearly within the interval of  $2\Delta x$ . Thus, if the actual value of a dependent variable at a given  $x$ -position is more than the interpolated value, the computed value at the same position for a later time will be less than it should be. This effect produces a dampening effect in time at a fixed location. Figure 4.1 demonstrates this effect for the depth at a location near the free-fall outlet. The greater the curvature of the free surface the more pronounced is this effect

To reduce this effect the physical size of  $\Delta x$  may be reduced but this results in an increase of the computer time needed. The computer time increases by the square of the number  $n$  of distance intervals,  $\Delta x$ . Subsequent comparison indicate that the diffusing scheme requires more computer time than the other two schemes.

4.3 Properties of Lax-Wendroff Scheme

The Lax-Wendroff scheme is an improvement over the diffusing scheme in that it accommodates the curvature in the variation of dependent variables. This, however, involves a more complicated numerical algorithm.

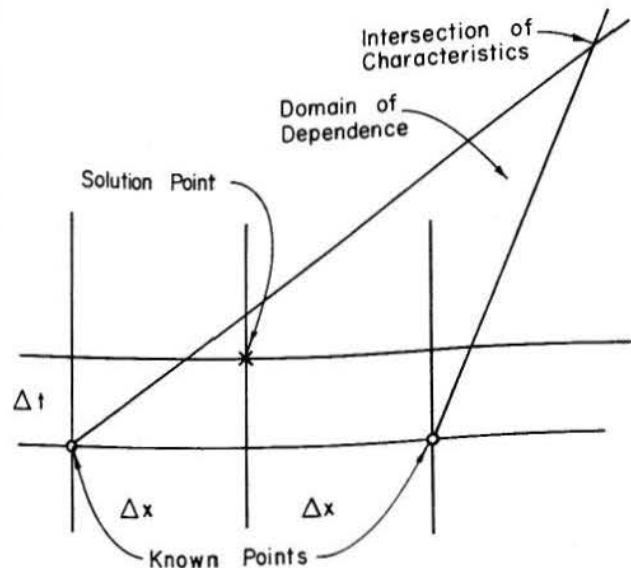


Fig. 4.1. Effect of characteristic slopes.

The Lax-Wendroff scheme results in a more accurate solution in comparison with the diffusing scheme for the same  $\Delta x$  and  $\Delta t$  intervals without a significant increase in computer time. An indication of this improved accuracy is demonstrated in Fig. 4.2. The Lax-Wendroff method consistently produces the same depth over a very large period of time, whereas, the diffusing produces a consistent change.

With regard to its flexibility in accommodating a wide range of flow conditions, the Lax-Wendroff scheme possesses the same inherent limitations as the diffusing scheme. Thus, by the Lax-Wendroff scheme the further the intersection of the two characteristic curves from the solution point, the less accurate the solution.

#### 4.4 Properties of Specified Intervals Scheme of the Method of Characteristics

The complications inherent in the specified intervals scheme of the method of characteristics are justified because of its inherent accuracies. The

basis for this is that the points of solutions are at the intersections of characteristic curves, rather than at any point within the domain of dependence.

The linear interpolation of this scheme is made without the need of a computer decision. All flow conditions can be accommodated by this scheme.

The accuracy of this scheme is demonstrated in Fig. 4.2, and is very good when compared to the diffusing and Lax-Wendroff schemes.

It is apparent that this finite-difference scheme of the method of characteristics produces a rapidly convergent and stable value. It is comparable to the same property of the Lax-Wendroff scheme.

The non-linear interpolation of the method of characteristics for dependent variables along distances for a given time is an improvement over the linear interpolation. However, linear interpolation is used in producing results (C) of Fig. 4.2 for this method of characteristics.

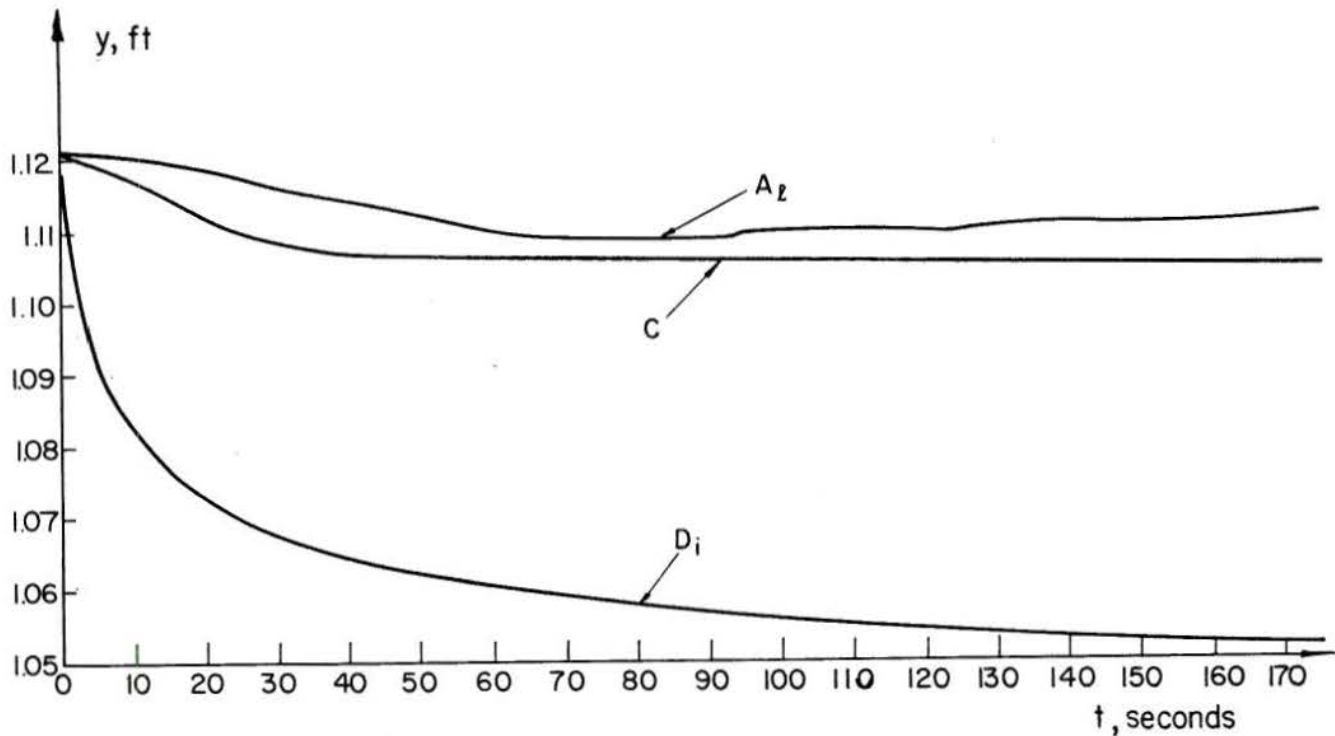


Fig. 4.2. Comparison of diffusing scheme ( $D_i$ ), Lax-Wendroff scheme ( $A_2$ ), and the specified intervals scheme of method of characteristics (C) in reproducing the steady initial conditions along the conduit, at the distance  $x = 796.7$  ft.

## Chapter 5

### CONCLUSIONS AND RECOMMENDATIONS

#### 5.1 Conclusions

1. Numerical integration solutions to the differential equations of gradually varied free-surface unsteady flow in prismatic channels and conduits have been reviewed, evaluated, and compared, both by the integration of the two partial differential equations and by their equivalent, four ordinary characteristic differential equations.

2. Numerical integration schemes, their solutions and their resulting computer programs are compared on the basis of their simplicity, stability, accuracy, flexibility, and the resulting computer time needed under given physical conditions.

3. Second-order or non-linear interpolations for dependent variables in the finite-difference schemes, for both the Lax-Wendroff scheme and the specified intervals scheme of the method of characteristics, were found to be necessary if maximum accuracy is to be obtained.

4. Solutions by the specified intervals scheme of the method of characteristics, with the second-order or non-linear interpolations for dependent variables, do not significantly require more computer time for a given accuracy comparable to the accuracy of solutions by any other scheme.

5. The Lax-Wendroff finite-difference scheme requires some particular programming considerations and adjustments in the case of supercritical flow.

6. The finite-difference specified intervals scheme of the method of characteristics with the

second-order of non-linear interpolations of dependent variables is sufficiently flexible to accommodate a large range of flow conditions.

7. Numerical integration by the specified intervals scheme of the method of characteristics with the second-order or non-linear interpolations of dependent variables in the writers' opinion should be used in general for studies of gradually varied free-surface unsteady flow.

#### 5.2 Recommendations

Four recommendations for further studies are present in the following:

1. Other numerical integration finite-difference schemes, periodically appearing in the literature or not studied in this paper, should be investigated and compared with the recommended finite-difference specified intervals scheme of the method of characteristics. This should be done to find whether improvements in overall applicability can be attained.

2. The finite-difference specified intervals scheme of the method of characteristics may be further improved by considering the curvilinear nature of the characteristic curves. Thus, a better method of interpolation may be designed.

3. For the integration of gradually varied free-surface unsteady flow equations the use of a hybrid computer should be particularly investigated.

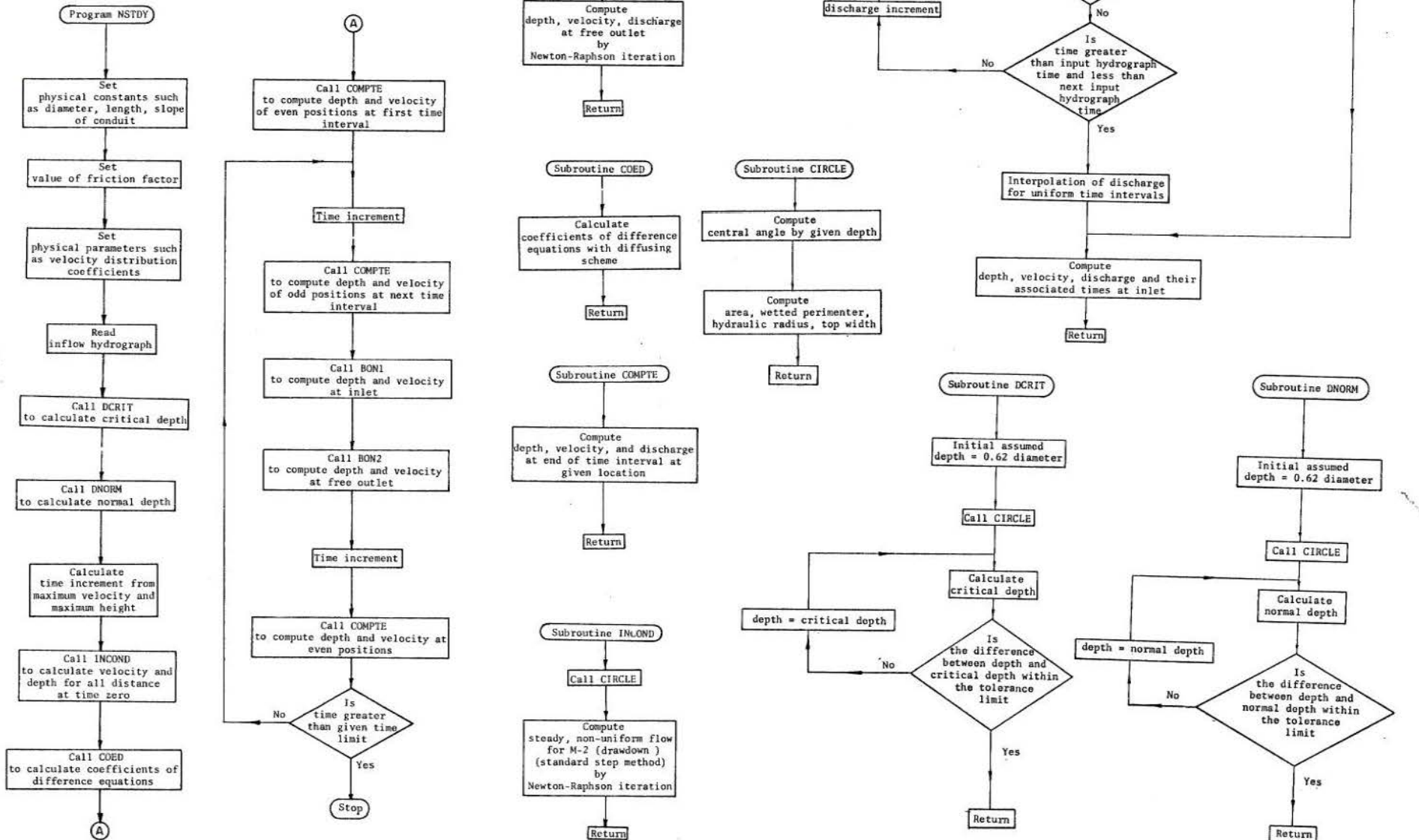
4. Computer times and computer costs should be systematically investigated for the most popular digital computers and for various finite-difference schemes.

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# APPENDIX I COMPUTATIONAL DETAILS OF DIFFUSING SCHEME

## A.I. FLOW CHART



Input hydrograph coordinates (Q,t) may be selected arbitrarily.





SUBROUTINE FOR COMPUTING UPSTREAM BOUNDARY CONDITION

```

SUBROUTINE BON1
DIMENSION TG(200), QI(200)
DIMENSION G(400), H(400), V(400)
DIMENSION G(330)
COMMON DN,H1,A,AP,DP,A1,C1,D1,E1,DZ,AZ,CZ,EZ,VP,DTA,DT,THA,DC
COMMON DX,XF,GR,ALPHA,BETA,SO,F,H,V,Q,DX,DT,T,TO,TF,N,FB,FC,B
COMMON M,MM,MM,L,I,PERD,DDT,VA,IQ,TQ,QI,NGCD
COMMON HA,HH,VM,HT,VT,NPT,HH,G,QII
COMMON THETA,WP,R,DEPTH,VC
COMMON QMAX(400),VMAX(400),HMAX(400)
COMMON TQMAX(400),TVMAX(400),THMAX(400)
1 IF (IG,GE,NGCD) 2,3
2 QT=QI(NGCD)
3 GO TO 6
4 IF (T,GE,TQ(IG).AND,T,LT,TQ(IQ+1)) 5,4
5 TG=IQ+1
6 GO TO 1
7 QT=QI(IQ)+QI(IQ+1)-QI(IQ)*(T-TQ(IQ))/(TQ(IQ+1)-TQ(IQ))
8 HN=H(1)
9 THETA=2.0*ATANF((SGRTF(D*H(2)-H(2)**2))/(D*0.5-H(2)))
10 IF (THETA) 7,8,8
11 THETA=6.28318+THETA
12 A=0.125*(THETA-SINF(THETA))*(D*D)
13 WP=D*0.5*THETA
14 R=A/WP
15 A2=V(2)*ALPHA/GR
16 SF=1.25*F*D2*V(2)*V(2)/K
17 E2=SF-SO
18 SW=SGRTF((D*HN-HN**H))
19 THETA=2.0*ATANF((SWRTF(SG))/(D*0.5-HN))
20 IF (THETA) 10,11,11
21 THETA=6.28318+THETA
22 AX=0.125*(THETA-SINF(THETA))*(D*D)
23 FH=HN-A2*(V(3)+VA-V(1))-DT/AX-D2*DX/DT*(V(3)+QT/AX-VA-V(1))-C2*(HA*U
24 I+H(3)-H(1))-4.0*DX*E2
25 DAX=0.25*D*D*(1.0-COSF(THETA))/SQ
26 FPH=1.0-(A2-B2*DX/DT)*(GT*DAX/(AX*AX))
27 HNU=HN-FH/FPH
28 IF (ABS(FHNU-HN)-0.0001) 13,12,12
29 HN=HNU
30 GO TO 9
31 H(1)=HNU
32 Q(1)=GT
33 V(1)=GT/AX
34 IF (H(1),LT,HMAX(1)) GO TO 14
35 HMAX(1)=H(1)
36 THMAX(1)=T
37 IF (V(1),LT,VMAX(1)) GO TO 15
38 VMAX(1)=V(1)
39 TVMAX(1)=T
40 IF (Q(1),LT,QMAX(1)) GO TO 16
41 QMAX(1)=Q(1)
42 TGMAX(1)=1
43 RETURN
44 END

```

SUBROUTINE FOR COMPUTING CRITICAL DEPTH

```

SUBROUTINE DCRIT
DIMENSION TG(200), QI(200)
DIMENSION G(400), H(400), V(400)
DIMENSION G(330)
COMMON DN,H1,A,AP,DP,A1,C1,D1,E1,DZ,AZ,CZ,EZ,VP,DTA,DT,THA,DC
COMMON DX,XF,GR,ALPHA,BETA,SO,F,H,V,Q,DX,DT,T,TO,TF,N,FB,FC,B
COMMON M,MM,MM,L,I,PERD,DDT,VA,IQ,TQ,QI,NGCD
COMMON HA,HH,VM,HT,VT,NPT,HH,G,QII
COMMON THETA,WP,R,DEPTH,VC
COMMON QMAX(400),VMAX(400),HMAX(400)
COMMON TQMAX(400),TVMAX(400),THMAX(400)
1 THETA=2.0*ATANF((SGRTF(D*DX-DX**2))/(D*0.5-DX))
2 IF (THETA) 2,3,3
3 THETA=6.28318+THETA
4 A=0.125*(THETA-SINF(THETA))*(D*D)
5 B=D*SINF(THETA*0.5)
6 DC=DX-(B*(A**3)-ALPHA*(D*Q(1)**2)/K)/(3.0*(D**3)**2)-(2.0*(A**3)
7 I-COSF(THETA*0.5))/(SINF(THETA*0.5)))
8 IF (ABS(F/C-DX)-0.0001) 5,4,4
9 DX=DC
10 GO TO 1
11 VC=Q(1)/A
12 RETURN
13 END

```

SUBROUTINE FOR COMPUTING DOWNSTREAM BOUNDARY CONDITION

```

SUBROUTINE BON2
DIMENSION TG(200), QI(200)
DIMENSION G(400), H(400), V(400)
DIMENSION G(330)
COMMON DN,H1,A,AP,DP,A1,C1,D1,E1,DZ,AZ,CZ,EZ,VP,DTA,DT,THA,DC
COMMON DX,XF,GR,ALPHA,BETA,SO,F,H,V,Q,DX,DT,T,TO,TF,N,FB,FC,B
COMMON M,MM,MM,L,I,PERD,DDT,VA,IQ,TQ,QI,NGCD
COMMON HA,HH,VM,HT,VT,NPT,HH,G,QII
COMMON THETA,WP,R,DEPTH,VC
COMMON QMAX(400),VMAX(400),HMAX(400)
COMMON TQMAX(400),TVMAX(400),THMAX(400)
HP=H(1)
VP=V(1)
THETA=2.0*ATANF((SGRTF(D*HP-HP*HP))/(D*0.5-HP))
1 IF (THETA) 1,1,2
2 THETA=THETA+0.20318
3 AP=0.125*(THETA-SINF(THETA))*(D*D)
4 BP=D*SINF(D*0.5*THETA)
5 THETA=2.0*ATANF((SGRTF(D*HX-HX*HX))/(D*0.5-HX))
6 IF (THETA) 4,5,5
7 THETA=6.28318+THETA
8 A=0.125*(THETA-SINF(THETA))*(D*D)
9 B=D*SINF(THETA*0.5)
10 VX=SGRTF(A*GR/B)
11 CT=COSE(THETA*0.5)/SINF(THETA*0.5)
12 FORG=AP/(BP*DX)*(VX+V(MM)-V(MMM))-VM)+VP/UX*(HX+H(MM)-H(MMM))-HM)+I
13 H(2)=H(1)+H(MM)-H(MMM)-HM)/DT
14 FPR1=AP/(D*DX*2.0*VX)*((GR-A*GR*2.0*CT)/(D*B))+VP/UA+1.0/DT
15 DC=HX-FORG/FPR1
16 IF (ABS(F/DC-HX)-0.0001) 7,6,6
17 HX=DC
18 GO TO 3
19 H(MM)=DC
20 THETA=2.0*ATANF((SGRTF(D*DC-DC*DC))/(D*0.5-DC))
21 IF (THETA) 8,9,9
22 THETA=6.28318+THETA
23 A=0.125*(THETA-SINF(THETA))*(D*D)
24 B=D*SINF(THETA*0.5)
25 V(MM)=SGRTF(A*GR/B)
26 Q(MM)=A*V(MM)
27 I=MM
28 IF (H(I),LT,HMAX(I)) GO TO 10
29 HMAX(I)=H(I)
30 THMAX(I)=T
31 IF (V(I),LT,VMAX(I)) GO TO 11
32 VMAX(I)=V(I)
33 TVMAX(I)=T
34 IF (Q(I),LT,QMAX(I)) GO TO 12
35 QMAX(I)=Q(I)
36 TQMAX(I)=T
37 RETURN
38 END

```

SUBROUTINE FOR COMPUTING GEOMETRIC PARAMETERS OF CIRCULAR SEGMENT

```

SUBROUTINE CIRCLL
DIMENSION TG(200), QI(200)
DIMENSION G(400), H(400), V(400), X(330)
COMMON DN,H1,A,AP,DP,A1,C1,D1,E1,DZ,AZ,CZ,EZ,VP,DTA,DT,THA,DC
COMMON DX,XF,GR,ALPHA,BETA,SO,F,H,V,Q,DX,DT,T,TO,TF,N,FB,FC,B
COMMON M,MM,MM,L,I,PERD,DDT,VA,IQ,TQ,QI,NGCD
COMMON HA,HH,VM,HT,VT,NPT,HH,G,QII
COMMON THETA,WP,R,DEPTH,VC
THETA=2.0*ATANF((SGRTF(DIA*DCPTH-DEPTH**2))/(DIA*0.5-DEPTH))
1 IF (THETA) 1,2,2
2 THETA=6.28318+THETA
3 A=0.125*(THETA-SINF(THETA))*(DIA*DIA)
4 WP=(DIA*0.5)*THETA
5 R=A/WP
6 B=DIA*SINF(THETA/2.0)
7 RETURN
8 END

```

SUBROUTINE FOR COMPUTING COEFFICIENTS IN DIFFERENCE EQUATIONS

```

SUBROUTINE COEV
DIMENSION TQ(200), QI(200)
DIMENSION Q(400), H(400), V(400)
DIMENSION G(330)
COMMON DN,H1,A,AP,OP,A1,C1,D1,E1,OP2,A2,C2,E2,VP,UTA,UT,HX,UC
COMMON D,XO,XF,GR,ALPHA,BETA,SO,F,H,V,Q,UX,DT,T,TU,TF,N,FB,FC,OB
COMMON N,MM,MMH,L,I,PERD,DDT,VA,IQ,TU,QI,NUCD
COMMON HA,HM,VM,HT,VT,NPT,HH,G,GII
COMMON THETA,W,P,R,DEPTH,VC
VT=(V(I+1)+V(I-1))*0.5
HT=(H(I+1)+H(I-1))*0.5
THETA=2.0*ATANF((SQRTF(D*HT-HT**2))/(D*0.5-HT))
IF (THETA) 1,2,2
ThETA=0.20310+THETA
A=0.125*(THETA-SINF(ThETA))*(D*D)
WP=D*0.5*THETA
R=A/WP
B=D*SINF(THETA*0.5)
A1=A/(VT*0)
D1=1.0/VT
A2=VT*ALPHA/GR
SF=0.125*F*0.2*VT*VT/R
E2=SF-SO
RETURN
END
    
```

SUBROUTINE FOR COMPUTING DEPTH & VELOCITY AT END OF TIME INTERVAL

```

SUBROUTINE COMPT
DIMENSION TQ(200), QI(200)
DIMENSION G(400), H(400), V(400)
DIMENSION G(330)
COMMON DN,H1,A,AP,OP,A1,C1,D1,E1,OP2,A2,C2,E2,VP,UTA,UT,HX,UC
COMMON D,XO,XF,GR,ALPHA,BETA,SO,F,H,V,Q,UX,DT,T,TU,TF,N,FB,FC,OB
COMMON N,MM,MMH,L,I,PERD,DDT,VA,IQ,TU,QI,NUCD
COMMON HA,HM,VM,HT,VT,NPT,HH,G,GII
COMMON THETA,W,P,R,DEPTH,VC
COMMON QMAX(400),VMAX(400),HMAX(400)
COMMON TVMAX(400),TVMAX(400),TVMAX(400)
H(I)=HT-(DT/2.0*UX*0.1)*(A1*(V(I+1)-V(I-1))+H(I+1)-H(I-1))
V(I)=VT-(DT/2.0*UX*0.1)*(A2*(V(I+1)-V(I-1))+H(I+1)-H(I-1))/(2.0*UX+E2)
VP=V(I)
Q(I)=V(I)*A
IF (H(I).LT.HMAX(I)) GO TO 1
HMAX(I)=H(I)
THMAX(I)=T
IF (V(I).LT.VMAX(I)) GO TO 2
VMAX(I)=V(I)
TVMAX(I)=T
IF (Q(I).LT.QMAX(I)) GO TO 3
QMAX(I)=Q(I)
TQMAX(I)=T
RETURN
END
    
```

SUBROUTINE FOR COMPUTING NORMAL DEPTH

```

SUBROUTINE DNORM
DIMENSION TQ(200), QI(200)
DIMENSION G(400), H(400), V(400)
DIMENSION G(330)
COMMON DN,H1,A,AP,OP,A1,C1,D1,E1,OP2,A2,C2,E2,VP,UTA,UT,HX,UC
COMMON D,XO,XF,GR,ALPHA,BETA,SO,F,H,V,Q,UX,DT,T,TU,TF,N,FB,FC,OB
COMMON N,MM,MMH,L,I,PERD,DDT,VA,IQ,TU,QI,NUCD
COMMON HA,HM,VM,HT,VT,NPT,HH,G,GII
COMMON THETA,W,P,R,DEPTH,VC
THETA=2.0*ATANF((SQRTF(D*H1-H1**2))/(D*0.5-H1))
IF (THETA) 2,3,3
ThETA=0.20310+THETA
A=0.125*(THETA-SINF(THETA))*(D*D)
WP=(0.5*D)*THETA
R=A/WP
B=D*SINF(THETA*0.5)
DN=H1-(WP*(F*G(I)*WII)/(D*0*GR*SU*K*R*A))/(3.0*0)/K-2.0/G/SINF(ThETA)
IF (ABSF(DN-H1)-0.0001) 7,4,4
IF (D-DN) 5,5,5
DN=DN*0.5
GO TO 4
H1=DN
GO TO 1
RETURN
END
    
```

SUBROUTINE FOR COMPUTING INITIAL CONDITION

```

SUBROUTINE INCOND
DIMENSION TQ(200), QI(200)
DIMENSION G(400), D(400), V(400), X(350)
COMMON DN,H1,A,AP,OP,A1,C1,D1,E1,OP2,A2,C2,E2,VP,UTA,UT,HX,UC
COMMON DIA,XO,XF,GR,ALPHA,BETA,SO,F,H,V,Q,UX,DT,T,TU,TF,N,FB,FC,OB
COMMON N,MM,MMH,L,I,PERD,DDT,VA,IQ,TU,QI,NUCD
COMMON HA,HM,VM,HT,VT,NPT,HH,X,QB
COMMON THETA,W,P,R,DEPTH,VC
DTOL=0.00001
IF (DN-DC) 1,1,2
K=1
GO TO 17
DIN=(DN+UC)*0.5
DEPTH=DC
CALL CIRCLE
VV=G0/A
VH=(VV*VV)/(2.0*GR)
S1=F*VH/(4.0*R)
EE1=DC+ALPHA*VH
D(2*N+1)=DC
Q(2*N+1)=QB
V(2*N+1)=VV
NCOUNT=0
DO 10 L=1,N
DEPTH=DIN
CALL CIRCLE
HFTH=0.5*THETA
DIHET=4.0/(DIA*SINF(HFTH))
DAREA=G.L25*QI*A*DIA*(1.0-COSF(THETA))*DT*NET
DN0=5*DIA*DIHET
DRA=(WP*DAREA+NDW)/(WP*WP)
DENG=1.0-(QB*UB)/(GR*(A**3))*DAREA
DSLO=-F*QB*G0*(2.0*R*A*DAREA+(A**2)*DRA)/(8.0*G0*(K*A**2)**2)
VV=G0/A
VH=(VV*VV)/(2.0*GR)
S2=F*VH/(4.0*R)
SF=(S1+S2)*0.5
EE2=DIN+ALPHA*VH
FRATIO=(EE2-EE1+2.0*UX*(SO-SF))/(DENU+(EE2-EE1)*USLO/(SO-SF))
DCOM=DIN-FRATIO
IF (DCOM) 5,4,6
WRITE (6,19)
GO TO 18
DCOM=ABSF(DCOM)
IF (ABSF(DCOM-DIN)-DTOL) 15,15,7
IF (0.82*DIA-DCOM) 8,14,14
DIN=DCOM*0.5
IF (0.82*DIA-DIN) 10,10,11
DIN=DIN*0.5
NCOUNT=NCOUNT+1
GO TO 9
IF (NCOUNT-20) 12,12,13
GO TO 3
WRITE (6,20)
GO TO 18
DIN=DCOM
GO TO 3
DIN=DCOM
S1=S2
EE1=EE2
I1=2*(N-L)+1
D(I1)=DIN
V(I1)=VV
Q(I1)=QB
CONTINUE
GO TO 18
WRITE (6,21) K
RETURN
C-----
FORMAT (* DCOM EQUALS ZERO *)
FORMAT (25H D2 MUCH GREATER THAN DIA)
FORMAT (* STOP *,I3)
END
    
```





TIME IS 1.1400844E+02 SEC.

Table with 4 columns: PNT, H, V, Q. Contains 21 rows of data points.

TIME IS 1.79201039E+02 SEC.

Table with 4 columns: PNT, H, V, Q. Contains 21 rows of data points.

TIME IS 1.30304342E+02 SEC.

Table with 4 columns: PNT, H, V, Q. Contains 21 rows of data points.

TIME IS 1.95504588E+02 SEC.

Table with 4 columns: PNT, H, V, Q. Contains 21 rows of data points.

TIME IS 1.46809941E+02 SEC.

Table with 4 columns: PNT, H, V, Q. Contains 21 rows of data points.

MAXIMUM VALUES AND TIMES AT EACH LOCATION

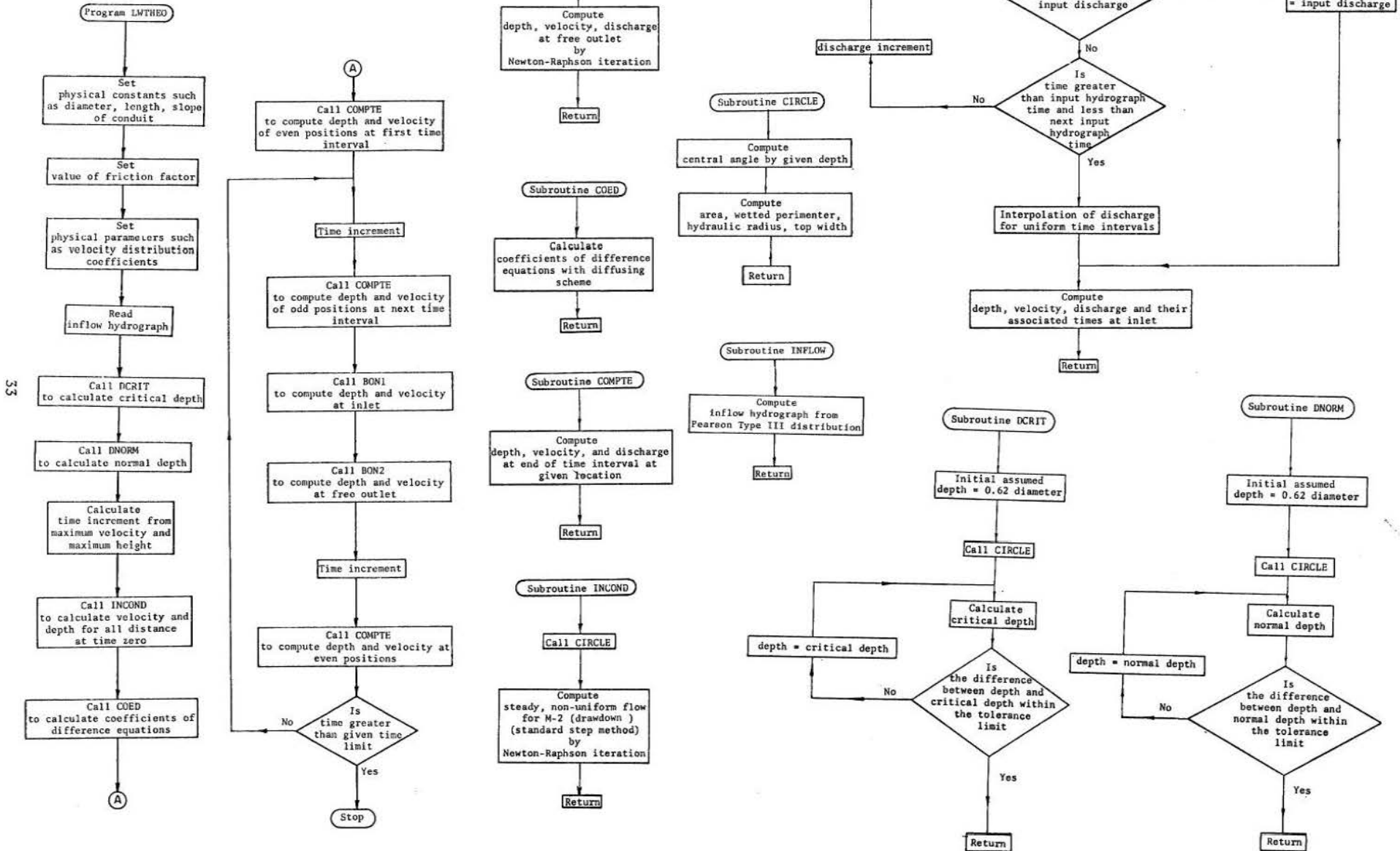
Table with 7 columns: DISTANCE, MAX DEPTH, TIME, MAX VEL, TIME, MAX Q, TIME. Contains 21 rows of maximum values and times.

TIME IS 1.62905490E+02 SEC.

Table with 4 columns: PNT, H, V, Q. Contains 21 rows of data points.

# APPENDIX 2 COMPUTATIONAL DETAILS OF LAX-WENDROFF SCHEME

## A.2.1. FLOW CHART



## A.2.2. FORTRAN IV COMPUTER PROGRAM

MAIN PROGRAM FOR UNSTEADY FLOW BY LAX-WENDROFF SCHEME

```

PNCGRKX LxTHEO(INPUT,OUTPUT,TAPE5=INPUT,TAPL6=OUTPUT)
DIMENSION TG(3300), QI(3300)
DIMENSION G(400), H(400), V(400)
DIMENSION G1(330)
COMMON DN,H1,A,AP,DP,A1,C1,D1,E1,b2,A2,Cc,E2,VP,DTA,QT,hd,DC
CGHHA D,XO,XF,GR,ALPHA,BETA,SO,F,H,V,w,DX,DT,T,TO,TF,N,FB,FC,B
CUPHHA M,MN,MM,L,I,PERD,DDT,VA,IG,TQ,QI,NGCD
COMMON HA,HH,VM,HT,VT,NPT,HH,G,QII
COMMON THETA,WP,R,DEPT,VC,J,HN,VN
COMMON LMAX(400),VMAX(400),HMAX(400)
COMMON TGMAX(400),TVMAX(400),THMAX(400)
INTEGER KUN
DO 1 I=1,400
HMAX(I)=G
VMAX(I)=C
QMAX(I)=G
1 CONTINUE
C-----INPUT WHICH MAY BE ALTERED
D=2.9262
SG=0.001
QB=10.
T=0.0
XO=0.0
XF=821.70
F=0.012
ALPHA=1.00
BETA=1.00
GR=32.175
H1=0.4*D
IQ=1
C-----END OF INPUT WHICH MAY BE ALTERED
C-----INITIAL TIME, FINAL TIME, INITIAL HEIGHT, NUMBER OF POINTS PER ROW
TO=0.
TF=200.
N=20
NI=N+1
IXOX=1
NPO=6
FB=0.109394
FC=-0.17944
NT=NPO
C-----CALCULATION OF CRITICAL AND NORMAL DEPTH AT BASE FLOW
QII=QB
CALL DNORM
DX=DN
DNQD=DN
CALL DCRIT
DCQB=DC
XL=XF-XO-4.5*DC
FD=SQRTF(GR*A/B)
FM=N
C-----CALCULATION OF DT FROM MAXIMUM VELOCITY AND MAXIMUM HEIGHT
DX=XL/FM
RA=1.0/(2.C*VC)
DI=RA*DX
DT=DI*0.5
DO 2 J=2,3300
CALL INFLOW
NGCD=J
2 CONTINUE
WRITE (6,11) DNQD,DCQB,DN,DC
WRITE (6,12) N,DX,DT,XO,XF,TO,TF,SO,D,F
WRITE (6,13) RA,H1,PERD,FB,FC
C-----CALCULATION OF INITIAL CONDITIONS
C-----HEIGHTS AT PARTICULAR DISTANCES FROM INLET END
CALL INCOND
DTA=DT
3 IF (NPO-NT) 4,4,6
4 WRITE (6,14) T
NT=0
WRITE (6,15)
DO 5 I=1,N1,IXOX
WRITE (6,16) I,H(I),V(I),Q(I)
5 CONTINUE

```

```

LWT 1
LWT 2
LWT 3
LWT 4
LWT 5
LWT 6
LWT 7
LWT 8
LWT 9
LWT 10
LWT 11
LWT 12
LWT 13
LWT 14
LWT 15
LWT 16
LWT 17
LWT 18
LWT 19
LWT 20
LWT 21
LWT 22
LWT 23
LWT 24
LWT 25
LWT 26
LWT 27
LWT 28
LWT 29
LWT 30
LWT 31
LWT 32
LWT 33
LWT 34
LWT 35
LWT 36
LWT 37
LWT 38
LWT 39
LWT 40
LWT 41
LWT 42
LWT 43
LWT 44
LWT 45
LWT 46
LWT 47
LWT 48
LWT 49
LWT 50
LWT 51
LWT 52
LWT 53
LWT 54
LWT 55
LWT 56
LWT 57
LWT 58
LWT 59
LWT 60
LWT 61
LWT 62
LWT 63
LWT 64
LWT 65
LWT 66
LWT 67
LWT 68
LWT 69
LWT 70
LWT 71
LWT 72
LWT 73
LWT 74

```

```

6 NT=NT+1
T=T+DTA
QA=Q(2)
HA=H(2)
VA=V(2)
QM=G(N)
HM=H(N)
VM=V(N)
QN=G(N+1)
HN=H(N+1)
VN=V(N+1)
DO 7 I=2,N
C-----CALCULATION OF COEFFICIENTS AND SOLUTION OF DIFFERENCE EQUATIONS
CALL COED
CALL COMPT
7 CONTINUE
C-----CALCULATION OF INLET BOUNDARY CONDITIONS
CALL BUN1
HD=HN
C-----CALCULATION OF OUTLET BOUNDARY CONDITIONS
CALL BUN2
IF (TF-T) 8,3,3
8 CONTINUE
NPG=N1/50+1
DO 9 III=1,NPG
II=50*III-49
IL=II+49
WRITE (6,17)
WRITE (6,18)
DO 9 I=II,IL
X=(I-1)*DX
WRITE (6,19) X,HMAX(I),THMAX(I),VMAX(I),TVMAX(I),QMAX(I),TWMAX(I)
IF (I.EQ.N1) GO TO 10
9 CONTINUE
10 CALL EXIT
C-----
11 FORMAT (*1DNQB = *E16.8// *DCQB = *E16.8//
1 *DNPF = *E16.8// *DCPF = *E16.8//
12 FORMAT (* N = *I5//
1 *DX = *E16.8//
2 *DT = *E16.8//
3 *XU = *E16.8//
4 *XF = *E16.8//
5 *TU = *E16.8//
6 *TF = *E16.8//
7 *SO = *E16.8//
8 *D = *E16.8//
9 *F = *E16.8//
13 FORMAT (* RA = *E16.8//
1 *H1 = *E16.8//
2 *PERD = *E16.8//
3 *FB = *E16.8//
4 *FC = *E16.8//
14 FORMAT (1H1,7HTIME IS,16.0,5M SEC.)
15 FORMAT (2X,3HPNT,10X,1HH,17X,1HV,17X,1HG)
16 FORMAT (1X,14,2X,E16.0,cX,E16.0,cX,E16.0)
17 FORMAT (// *1 MAXIMUM VALUES AND TIMES AT EACH LOCATION//)
18 FORMAT (* DISTANCE MAX DEPTH TIME MAX VEL TIME MAX G
1 TIME*)
19 FORMAT (F8.2,3(4X,F6.2,2X,F7.2))
END

```

### SUBROUTINE FOR COMPUTING GEOMETRIC PARAMETERS OF CIRCULAR SEGMENT

```

SUBROUTINE CIRCLE
DIMENSION TG(3300), QI(3300)
DIMENSION G(400), D(400), V(400), X(330)
COMMON DN,H1,A,AP,DP,A1,C1,D1,E1,b2,A2,Cc,E2,VP,DTA,QT,hd,DC
COMMON DIA,XO,XF,GR,ALPHA,BETA,SO,F,D,V,w,DX,DT,T,TO,TF,N,FB,FC,B
COMMON M,MN,MM,L,I,PERD,DDT,VA,IG,TQ,QI,NGCD
COMMON HA,HH,VM,HT,VT,NPT,HH,X,QB
COMMON THETA,WP,R,DEPT,VC,J,HN,VN
THETA=2.0*ATANF((SQRTF(DIA*DEPTH-DEPTH**2))/(DIA/2.0-DEPTH))
IF (THETA) 1,2,2
1 THETA=6.28318-THETA
2 A=0.125*(THETA-SINF(THETA))*(DIA*DIA)
WP=(DIA/2.0)*THETA
R=A/WP
B=DIA*SINF(THETA/2.C)
RETURN
END
CIR 1
CIR 2
CIR 3
CIR 4
CIR 5
CIR 6
CIR 7
CIR 8
CIR 9
CIR 10
CIR 11
CIR 12
CIR 13
CIR 14
CIR 15
CIR 16
CIR 17-

```



SUBROUTINE FOR COMPUTING INITIAL CONDITION

```

SUBROUTINE INCOND      INC 1
DIMENSION I(13300), G(13300) INC 2
DIMENSION Q(400), D(400), V(400), X(330) INC 3
COMMON DN,H1,A,AP,DP,A1,C1,D1,E1,p2,A2,C2,e2,VP,DTA,DT,HO,UC INC 4
COMMON DIA,XD,XF,GR,ALPHA,BETA,SO,FO,V,GO,DX,DT,T,TO,TF,N,FO,FC,BO INC 5
COMMON M,MM,MMH,L,I,PERD,DDT,VA,I,GT,GI,N,UCD INC 6
COMMON HA,HM,VM,HT,VT,NPT,HH,G,GI INC 7
COMMON THETA,RP,R,DEPTH,VC,J,HN,VN INC 8
DTOL=0.0001 INC 9
ELEV=100.0 INC 10
AA=4.5*DC INC 11
ELLV=ELEV+SO*XX INC 12
IF (DA-DC) 1,1,2 INC 13
K=1 INC 14
1  GU TO 17 INC 15
2  DIN=1.75*DC INC 16
DEPTH=DC INC 17
CALL CIRCLE INC 18
VV=GO/A INC 19
VH=(VV*V)/(2.0*GR) INC 20
S1=F*VH/(4.0*R) INC 21
EL1=UC+ALPHA*VH INC 22
X(N+1)=AF-XX INC 23
G(N+1)=DC INC 24
Q(N+1)=GB INC 25
V(N+1)=VV INC 26
NCOUNT=0 INC 27
DO 16 L=1,N INC 28
XX=XX+DX INC 29
DEPTH=DIN INC 30
CALL CIRCLE INC 31
HFTH=G.5*THETA INC 32
DTHET=4.0/(DIA*SINF(HFTH)) INC 33
DAKEA=0.125*(A*VJA*(1.0-COSF(THETA))*DTHET INC 34
DA=0.5*DIA*DTHET INC 35
WP=0.5*DIA*THETA INC 36
DRA=(A*DAKEA-A*W)/(I*P*WP) INC 37
DENG=1.0-(GO*UD/(GR*(A**3)))*DAREA INC 38
DSLO=-F*GD*UD*(2.0*GR*A*DAKEA+(A**2)*DRA)/(16.0*GN*(I*(K*A**2)**2)) INC 39
VV=QB/A INC 40
VH=(VV*V)/(2.0*GR) INC 41
S2=F*VH/(4.0*R) INC 42
SF=(S1+S2)/2.0 INC 43
EL2=DIN+ALPHA*VH INC 44
FRATIO=(EL2-EE1+LX*(SO-SF))/(DENG+LEE2-EE1)*DSL/(SO-SF) INC 45
DCOM=DIN-FRATIO INC 46
IF (DCOM) 5,4,0 INC 47
WRITE (6,19) INC 48
GO TO 16 INC 49
5  DCOM=ABSF(DCOM) INC 50
6  IF (ABSF(DCOM-DIN)-DTOL) 19,15,7 INC 51
7  IF (0.02*DIA-DCOM) 8,14,14 INC 52
8  DIN=DCOM/2.0 INC 53
9  IF (0.82*DIA-DIN) 10,10,11 INC 54
10 DIN=DIN/2.0 INC 55
NCOUNT=NCOUNT+1 INC 56
GO TO 9 INC 57
11 IF (NCOUNT-20) 12,12,12 INC 58
12 GO TO 3 INC 59
13 WRITE (6,20) INC 60
GO TO 16 INC 61
14 DIN=DCOM INC 62
GO TO 3 INC 63
15 DIN=DCOM INC 64
S1=S2 INC 65
EE1=EE2 INC 66
II=N-L+1 INC 67
X(II)=XF-XX INC 68
D(II)=DIN INC 69
V(II)=VV INC 70
Q(II)=GB INC 71
16 CONTINUE INC 72
GO TO 16 INC 73
17 WRITE (6,21) K INC 74
18 RETURN INC 75
C----- INC 76
19 FORMAT (* DCOM EQUALS ZERO *) INC 77
20 FORMAT (25H D2 MUCH GREATER THAN DIA) INC 78
21 FORMAT (* STOP *,I3) INC 79
END INC 80

```

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SUBROUTINE FOR COMPUTING INFLOW HYDROGRAPH

```

SUBROUTINE INFLCW      INF 1
DIMENSION TQ(3300), GI(3300) INF 2
DIMENSION G(400), H(400), V(400) INF 3
DIMENSION G(3300) INF 4
COMMON DN,H1,A,AP,DP,A1,C1,D1,E1,p2,A2,C2,e2,VP,DTA,DT,HO,UC INF 5
COMMON DX,XD,XF,GR,ALPHA,BETA,SO,FO,V,GO,DX,DT,T,TO,TF,N,FO,FC,BO INF 6
COMMON M,MM,MMH,L,I,PERD,DDT,VA,I,GT,GI,N,UCD INF 7
COMMON HA,HM,VM,HT,VT,NPT,HH,G,GI INF 8
COMMON THETA,RP,R,DEPTH,VC,J,HN,VN INF 9
OO=10.0 INF 10
TP=100.0 INF 11
TG=150.0 INF 12
UU=TP/(TG-TP) INF 13
AJ=J INF 14
TQ(J)=(AJ-1.0)*DT INF 15
QI(J)=QE+GO*(EXPF(-(TQ(J)-TP)/(TG-TP)))*(TG(J)/TP)**OO INF 16
RETURN INF 17
END INF 18

```

SUBROUTINE FOR COMPUTING DEPTH & VELOCITY AT END OF TIME INTERVAL

```

SUBROUTINE COMPIE      COM 1
DIMENSION G(3300) COM 2
DIMENSION TQ(3300), G(3300) COM 3
DIMENSION G(400), H(400), V(400) COM 4
COMMON DN,H1,A,AP,DP,A1,C1,D1,E1,p2,A2,C2,e2,VP,DTA,DT,HO,UC COM 5
COMMON DX,XD,XF,GR,ALPHA,BETA,SO,FO,V,GO,DX,DT,T,TO,TF,N,FO,FC,BO COM 6
COMMON M,MM,MMH,L,I,PERD,DDT,VA,I,GT,GI,N,UCD COM 7
COMMON HA,HM,VM,HT,VT,NPT,HH,G,GI COM 8
COMMON THETA,RP,R,DEPTH,VC,J,HN,VN COM 9
COMMON QMAX(400),VMAX(400),HMAX(400) COM 10
COMMON TQMAX(400),TVMAX(400),THMAX(400) COM 11
Z1=DT/DX COM 12
Z2=A/B COM 13
Z3=ALPHA/BETA COM 14
Z4=GR/BETA COM 15
H(I)=H(I)-0.5*Z1*(Z2*(V(I+1)-V(I-1))+V(I)*ln((I+1)-H(I-1)))+5*(Z1*COM 16
121)*((Z3+1.0)*V(I)*Z2*(V(I+1)-2.0*V(I)+V(I-1))+Z2*Z4*V(I)**2)*ln(COM 17
21+1)-2.0*H(I))+H(I-1)) COM 18
V(I)=V(I)-0.5*Z1*(Z3*V(I)*ln((I+1)-H(I-1))+Z4*ln((I+1)-H(I-1))+2.0*DCOM 19
1X*Z4*E2)+0.5*Z1*Z1*(Z3*Z3*V(I)**2+Z4*Z2)*(V(I+1)-2.0*V(I)+V(I-1))COM 20
2+(Z3+1.0)*Z4*V(I)*H(I+1)-2.0*H(I)+H(I-1)) COM 21
Q(I)=V(I)*A COM 22
IF (H(I).LT.HMAX(I)) GO TO 1 COM 23
HMAX(I)=H(I) COM 24
THMAX(I)=T COM 25
1 IF (V(I).LT.VMAX(I)) GO TO 2 COM 26
VMAX(I)=V(I) COM 27
2 IF (Q(I).LT.QMAX(I)) GO TO 3 COM 28
QMAX(I)=Q(I) COM 29
3 RETURN COM 30
END COM 31

```

SUBROUTINE FOR COMPUTING COEFFICIENTS IN DIFFERENCE EQUATIONS

```

SUBROUTINE COED      CUE 1
DIMENSION TQ(3300), G(3300) CUE 2
DIMENSION G(400), H(400), V(400) CUE 3
DIMENSION G(3300) CUE 4
COMMON DN,H1,A,AP,DP,A1,C1,D1,E1,p2,A2,C2,e2,VP,DTA,DT,HO,UC CUE 5
COMMON DX,XD,XF,GR,ALPHA,BETA,SO,FO,V,GO,DX,DT,T,TO,TF,N,FO,FC,BO CUE 6
COMMON M,MM,MMH,L,I,PERD,DDT,VA,I,GT,GI,N,UCD CUE 7
COMMON HA,HM,VM,HT,VT,NPT,HH,G,GI CUE 8
COMMON THETA,RP,R,DEPTH,VC,J,HN,VN CUE 9
THETA=2.0*ATANF(SURTF(D*(H(I)-H(I)**2))/(D/2.0-n(I))) CUE 10
IF (THETA) 1,2,2 CUE 11
1 THETA=0.20310+THETA CUE 12
2 A=0.125*(THETA-SINF(THETA))*(D*D) CUE 13
WP=(D/2.0)*THETA CUE 14
R=A/WP CUE 15
B=D*SINF(THETA/2.0) CUE 16
A1=A/(V(I)*B) CUE 17
C1=1.0 CUE 18
D1=1.0/V(I) CUE 19
E1=0.0 CUE 20
B2=BETA/GR CUE 21
A2=V(I)*ALPHA/GR CUE 22
C2=1.0 CUE 23
SF=.125*F*B2*V(I)*V(I)/R CUE 24
E2=SF*SO CUE 25
RETURN CUE 26
END CUE 27

```

## SUBROUTINE FOR COMPUTING UPSTREAM BOUNDARY CONDITION

```

SUBROUTINE BONZ
DIMENSION TQ(3300), QI(3300)
DIMENSION Q(400), H(400), V(400)
DIMENSION G(330)
COMMON DN,H1,A,AP,DP,A1,C1,D1,E1,DZ,AZ,CZ,EZ,VP,DTA,QT,HB,DC
COMMON DX,XF,GR,ALPHA,BETA,SOF,H,V,QU,DX,DT,T,TU,TF,N,FB,FC,B
COMMON M,MM,MMH,L,I,PERD,DDT,VA,IG,IU,UI,NUCD
COMMON HA,HH,VM,HT,VT,NPT,HH,G,QII
COMMON THETA,WP,R,DEPT,VC,J,HN,VN
COMMON QMAX(400),VMAX(400),HMAX(400)
COMMON TQMAX(400),TVMAX(400),THMAX(400)
1 IF (IG,UI,NUCD) Z=3
2 QT=UI*(NUC)
3 GU TO 6
4 IF (T,UB,TW(I),AND,T,LT,TW(I+1)) 5+4
5 IU=IU+1
6 GO TO 1
7 QT=Q(I)+(QI(IQ+1)-QI(IQ))*(T-TG(IU))/(TG(IQ+1)-TG(IU))
8 HL=H(I)
9 THETA=2.0*ATANF((SQRT(D*H(I)-H(I)**2))/(D/2.0-H(I)))
10 IF (THETA) 7+8+9
11 THETA=6.28318+THETA
12 A=0.125*(THETA-SINF(THETA))*(D*D)
13 WP=(D/2.0)*THETA
14 R=A/WP
15 AZ=V(I)*ALPHA/GK
16 SF=0.125*F*BZ*V(I)*V(I)/R
17 EZ=SF-SO
18 THETA=2.0*ATANF((SQRT(D*HL-HL**2))/(D/2.0-HL))
19 SO=SQRT(D*HL-HL**2)
20 IF (THETA) 1G,11+11
21 THETA=6.28318+THETA
22 AX=D.125*(THETA-SINF(THETA))*(U*U)
23 FH=HL-AZ*(V(3)+V(1)-QT/AX)-BZ*DX/DT*(V(3)+QT/AX-VA-V(1))-CZ*(HABU)
24 DAX=(D*U/D*0.0*(1.0-COSF(THETA)))*2.0/U*(D-Z.0*HL)**2/SU+4.0*SU)
25 FPH=1.0-(AZ-BZ*DX/DT)*(U/DAX)/(AX*AX)
26 HNU=HL-FH/FPH
27 IF (ABS(FHNU-HL)-0.00001) 13,12,12
28 HL=HNU
29 GO TO 9
30 H(I)=HNU
31 Q(I)=QT
32 V(I)=QT/AX
33 IF (H(I),LT,HMAX(I)) GO TO 14
34 HMAX(I)=H(I)
35 THMAX(I)=T
36 IF (V(I),LT,VMAX(I)) GO TO 15
37 VMAX(I)=V(I)
38 TVMAX(I)=T
39 IF (W(I),LT,QMAX(I)) GO TO 16
40 QMAX(I)=Q(I)
41 TQMAX(I)=T
42 RETURN
43 END

```

## SUBROUTINE FOR COMPUTING CRITICAL DEPTH

```

SUBROUTINE DCRIT
DIMENSION G(330)
DIMENSION Q(400), H(400), V(400)
DIMENSION TQ(3300), QI(3300)
COMMON DN,H1,A,AP,DP,A1,C1,D1,E1,DZ,AZ,CZ,EZ,VP,DTA,QT,HB,DC
COMMON DX,XF,GR,ALPHA,BETA,SOF,H,V,QU,DX,DT,T,TU,TF,N,FB,FC,B
COMMON M,MM,MMH,L,I,PERD,DDT,VA,IG,IU,UI,NUCD
COMMON HA,HH,VM,HT,VT,NPT,HH,G,QII
COMMON THETA,WP,R,DEPT,VC,J,HN,VN
1 THETA=2.0*ATANF((SQRT(D*DX-DX**2))/(D/2.0-DX))
2 IF (THETA) 2+3+3
3 THETA=6.28318+THETA
4 A=0.125*(THETA-SINF(THETA))*(D*D)
5 B=D*SINF(THETA/2.0)
6 DC=DX-(B*(A**3)-ALPHA*(D*QI)**2)/UR/(3.0*(U*A)**2-12.0*(A**3))
7 I=COSF(THETA/2.0)/(SINF(THETA/2.0))
8 IF (ABS(F(U-DX)-0.0001) 5+4+4
9 DX=DC
10 GO TO 1
11 VC=QII/A
12 RETURN
13 END

```

## SUBROUTINE FOR COMPUTING DOWNSTREAM BOUNDARY CONDITION

```

SUBROUTINE BONZ
DIMENSION TQ(3300), QI(3300)
DIMENSION Q(400), H(400), V(400)
DIMENSION G(330)
COMMON DN,H1,A,AP,DP,A1,C1,D1,E1,DZ,AZ,CZ,EZ,VP,DTA,QT,HB,DC
COMMON DX,XF,GR,ALPHA,BETA,SOF,H,V,QU,DX,DT,T,TU,TF,N,FB,FC,B
COMMON M,MM,MMH,L,I,PERD,DDT,VA,IG,IU,UI,NUCD
COMMON HA,HH,VM,HT,VT,NPT,HH,G,QII
COMMON THETA,WP,R,DEPT,VC,J,HN,VN
COMMON QMAX(400),VMAX(400),HMAX(400)
COMMON TQMAX(400),TVMAX(400),THMAX(400)
HP=(HN+H(N))/2.0
VP=(VN+V(N))/2.0
THETA=2.0*ATANF((SQRT(D*HP-H*HP))/(D/2.0-HP))
IF (THETA) 1+1+2
1 THETA=6.28318+THETA
2 AP=0.125*(THETA-SINF(THETA))*(D*D)
3 BP=D*SINF(THETA/2.0)
4 THETA=2.0*ATANF((SQRT(D*HD-HB*HB))/(D/2.0-HB))
5 IF (THETA) 4+5+5
6 THETA=6.28318+THETA
7 A=0.125*(THETA-SINF(THETA))*(D*D)
8 B=D*SINF(THETA/2.0)
9 VX=SQRT(A*GR/B)
10 CTN=COSF(THETA/2.0)/SINF(THETA/2.0)
11 FURG=(H(N)+HB-HM-HN)/DT+VP/DX*(HN+HB-HM-H(N))+AP/(B*P*DX)*(VN+VX-VMBU)
12 I=V(N)
13 FPR=AP/(B*P*DX*2.0*VX)*(GR-A*GR*2.0*CTN/(B*B))+VP/DX+1.0/DT
14 DC=HD-FURG/FPR
15 IF (ABS(F(DC-HB)-0.0001) 7+6+6
16 HB=DC
17 GO TO 3
18 HIN+1=DC
19 THETA=2.0*ATANF((SQRT(U*DC-DC*DC))/(D/2.0-DC))
20 IF (THETA) 8+9+9
21 THETA=6.28318+THETA
22 A=0.125*(THETA-SINF(THETA))*(D*D)
23 B=D*SINF(THETA/2.0)
24 VIN+1=SQRT(A*GR/B)
25 Q(N+1)=A*V(N+1)
26 I=N+1
27 IF (H(I),LT,HMAX(I)) GO TO 10
28 HMAX(I)=H(I)
29 THMAX(I)=T
30 IF (V(I),LT,VMAX(I)) GO TO 11
31 VMAX(I)=V(I)
32 TVMAX(I)=T
33 IF (Q(I),LT,QMAX(I)) GO TO 12
34 QMAX(I)=Q(I)
35 TQMAX(I)=T
36 RETURN
37 END

```

## SUBROUTINE FOR COMPUTING NORMAL DEPTH

```

SUBROUTINE DNORM
DIMENSION TQ(3300), QI(3300)
DIMENSION Q(400), H(400), V(400)
DIMENSION G(330)
COMMON DN,H1,A,AP,DP,A1,C1,D1,E1,DZ,AZ,CZ,EZ,VP,DTA,QT,HB,DC
COMMON DX,XF,GR,ALPHA,BETA,SOF,H,V,QU,DX,DT,T,TU,TF,N,FB,FC,B
COMMON M,MM,MMH,L,I,PERD,DDT,VA,IG,IU,UI,NUCD
COMMON HA,HH,VM,HT,VT,NPT,HH,G,QII
COMMON THETA,WP,R,DEPT,VC,J,HN,VN
1 THETA=2.0*ATANF((SQRT(D*H1-H1**2))/(D/2.0-H1))
2 IF (THETA) 2+3+3
3 THETA=6.28318+THETA
4 A=0.125*(THETA-SINF(THETA))*(D*D)
5 WP=(D/2.0)*THETA
6 R=A/WP
7 B=D*SINF(THETA/2.0)
8 DN=H1-(WP*(FQII*QI))/(6.0*GR*SO*R*H(A))/(3.0*B)/N-2.0/SINF(THETA)DN)
9 I=2.0)
10 IF (ABS(DN-H1)-0.0001) 5+4+4
11 H1=DN
12 GO TO 1
13 RETURN
14 END

```

### A.2.3. DEFINITION OF VARIABLES

NAME	DEFINITION	STATEMENT	NUMBER(S)				
A	AREA OF CIRCULAR SEGMENT	DNO 13	DCR 13	COE 13	801 23		
		R02 22	B02 37	CIR 12			
AX	INTERMEDIATE AREA	R01 33					
AJ	(?)	INF 14					
ALPHA	VELOCITY DIST FACTOR-ENERGY	LWT 26					
AP	INTERMEDIATE AREA	R02 17					
A1	(2)	COE 17					
A2	(2)	COE 22	R01 26				
R	FREE SURFACE WIDTH	DNO 16	DCR 14	COE 16	802 23		
		R02 38	CIR 15				
BETA	VELOCITY DIST FACTOR-MOMENTUM	LWT 27					
RP	INTERMEDIATE FREE SURFACE WIDTH	R02 18					
R2	(2)	COE 21					
CTN	COTANGENT OF ANGLE	R02 25					
C1	(2)	COE 18					
C2	(2)	COE 23					
D	DIAMETER OF PIPE	LWT 19	INC 24	INC 69			
DAREA	DERIVATIVE OF AREA WITH DEPTH	INC 34					
DAX	DERIVATIVE OF AX	R01 36					
DC	CRITICAL DEPTH	DCR 15	B02 29				
DCOM	COMPUTED DEPTH	INC 46	INC 50				
DCQB	CRITICAL DEPTH	LWT 48					
DENG	(?)	INC 38					
DEPTH	DEPTH OF FLOW	INC 17	INC 30				
DIN	INITIAL VALUE OF DEPTH	INC 16	INC 53	INC 55	INC 62		
		INC 64					
DN	NORMAL DEPTH	DNO 17					
DNQB	NORMAL DEPTH	LWT 46					
DQA	DERIVATIVE OF HYD RADIUS WITH DEPTH	INC 37					
DSLO	DERIVATIVE OF EN SLOPE WITH DEPTH	INC 39					
DT	INCREMENT OF TIME	LWT 55	LWT 56				
DTA	INCREMENT OF TIME	LWT 67					
DTHT	DERIVATIVE OF THETA WITH DEPTH	INC 33					
DTOL	MAX ERROR IN DEPTH CALCULATIONS	INC 9					
DW	DERIVATIVE OF WP WITH DEPTH	INC 35					
DX	INCREMENT OF DISTANCE	LWT 45	LWT 53	DCR 18			
D1	(2)	COE 19					
EE1	ENERGY AT KNOWN DEPTH	INC 22	INC 66				
EE2	ENERGY AT UNKNOWN DEPTH	INC 44					
ELEV	ELEVATION	INC 10	INC 12				
E1	(2)	COE 20					
E2	(2)	COE 25	B01 28				
F	DARCY-WEISBACH FRICTION FACTOR	LWT 25					
FB	FACTOR IN REY. NO.-DARCY FUNCTION	LWT 39					
FC	FACTOR IN REY. NO.-DARCY FUNCTION	LWT 40					
FD	CELERITY OF WAVE	LWT 50					
FH	(2)	R01 34					
FM	NUMNER OF DISTANCE INTERVALS	LWT 51					
FORG	(2)	R02 26					
FPH	(2)	R01 37					
FPRI	(2)	R02 28					
FRATIO	(?)	INC 45					
GR	ACCELERATION OF GRAVITY	LWT 28					
H	DEPTH OF FLOW	COM 16	B01 42	B02 33			
HA	DEPTH OF FLOW	LWT 78					
HB	DEPTH OF FLOW	LWT 93	B02 31				
HPTH	ONE-HALF THETA	INC 32					
HL	INTERMEDIATE DEPTH	R01 19	B01 40				
HM	INTERMEDIATE DEPTH	LWT 81					
HMAX	MAXIMUM DEPTH	LWT 14	COM 24	B01 46	B02 43		
HN	INTERMEDIATE DEPTH	LWT 84					
HNU	COMPUTED DEPTH	R01 38					
HP	INITIAL DEPTH	R02 12					
HI	INITIAL DEPTH	LWT 29	DNO 20				

I	(1)	R02 41					
IXOX	(1)	LWT 37					
II	(1)	LWT 100	INC 67				
IL	(1)	LWT 101					
IQ	(1)	LWT 30	B01 16				
K	(1)	INC 14					
N	NCOUNT ITERATION CONTROL COUNTER	INC 27	INC 56				
N	NUMBER OF DISTANCE INTERVALS	LWT 15					
NT	TIME INTERVAL, PRT OUT CTRL COUNTER	LWT 41	LWT 70	LWT 75			
NPG	PRINT OUT CONTROL	LWT 99					
NPD	TIME INTERVAL PRINTOUT LIMIT	LWT 38					
NQCONO.	OF INFLOW DISCHARGE VALUES	LWT 59					
NI	NO. OF X-POSITIONS	LWT 36					
Q	COMPUTED DISCHARGE	INC 25	INC 71	COM 22	B01 43		
		R02 40					
QA	INTERMEDIATE DISCHARGE	LWT 77					
QB	BASE FLOW	LWT 21					
QI	INTERPOLATED DISCHARGE	INF 16					
QII	INTERMEDIATE DISCHARGE	LWT 43					
QM	INTERMEDIATE DISCHARGE	LWT 80					
QMAX	MAX. DISCHARGE	LWT 16	COM 30	B01 52	B02 49		
QN	INTERMEDIATE DISCHARGE	LWT 83					
QO	EXCESS OF PEAK Q OVER QR	INF 10					
QT	INFLOW DISCHARGE	B01 13	B01 18				
R	HYDRAULIC RADIUS	DNO 15	COE 15	B01 25	CIR 14		
RA	RATIO OF DT TO DX	LWT 54					
SF	AVERAGE FRICTION SLOPE	INC 43	COE 24	B01 27			
S0	CHANNEL BED SLOPE	LWT 20					
S0	(2)	R01 30					
S1	FRICTION SLOPE	INC 21	INC 65				
S2	FRICTION SLOPE	INC 42					
T	TIME IN SECONDS	LWT 22	LWT 76				
TVMAX	TIME OF MAX. VFLOCITY	COM 28	B01 50	B02 47			
TF	TIME LIMIT FOR SOLUTION	LWT 34					
TG	TIME TO C. OF G. OF INFLOW HYDROGRAPH	INF 12					
THETA	CENTRAL ANGLE OF CIRCULAR SEGMENT	DNO 10	DNO 12	DCR 10	DCR 12		
		COE 10	COE 12	B01 20	B01 22		
		R01 29	B01 32	B02 14	B02 16		
		R02 19	B02 21	B02 34	B02 36		
		CIR 9	CIR 11				
THMAX	TIME OF MAX. DEPTH	COM 25	B01 47	B02 44			
TO	(?)	LWT 33					
TP	TIME TO PEAK INFLOW HYD. DISCHARGE	INF 11					
TO	(2)	INF 15					
TOMAX	TIME OF MAX. DISCHARGE	COM 31	B01 53	B02 50			
UU	(?)	INF 13					
V	VELOCITY	INC 26	INC 70	COM 19	B01 44		
		R02 39					
VV	VELOCITY	INC 19	INC 40				
VX	CELERITY OF WAVE	B02 24					
VA	INTERMEDIATE VELOCITY	LWT 79					
VC	CRITICAL VELOCITY	DCR 20					
VH	VELOCITY HEAD	INC 20	INC 41				
VM	INTERMEDIATE VELOCITY	LWT 82					
VMAX	MAXIMUM VELOCITY	LWT 15	COM 27	B01 49	B02 46		
VN	INTERMEDIATE VELOCITY	LWT 85					
VP	AVERAGE VELOCITY	R02 13					
WP	WETTED PERIMETER	INC 36	DNO 14	COE 14	B01 24		
		CIR 13					
X	POSITION ALONG CHANNEL	LWT 105	INC 23	INC 68			
XX	POSITION OF CRITICAL DEPTH	INC 11	INC 29				
XF	TOTAL LENGTH OF CHANNEL	LWT 24					
XL	WORKING LENGTH OF CHANNEL	LWT 49					
X0	INITIAL POSITION ALONG CHANNEL	LWT 23					
Z1	(2)	COM 12					
Z2	(2)	COM 13					
Z3	(?)	COM 14					
Z4	(2)	COM 15					

(1) DO-LOOP COUNTER OR VARIABLE SUBSCRIPT  
 (2) INTERMEDIATE VARIABLE

# A.2.4. SAMPLE OUTPUT

(No input required)

DNQB = 1.22058217E+00  
 DCQB = 1.00785443E+00  
 DNPB = 1.22058217E+00  
 DCPB = 1.00785443E+00  
 N = 20  
 DX = 4.09592327E+01  
 DT = 2.07629670E+00  
 XO = 0.  
 XF = 8.21700000E+02  
 TO = 0.

TF = 2.00000000E+02  
 SO = 1.00000000E+03  
 D = 2.42620000E+00  
 F = 1.20000000E+02  
 RA = 1.02613185E+01  
 H1 = 1.22859067E+00  
 PERD = -4.13844033-195  
 FB = 1.09394000E+01  
 FC = -1.74440000E+01

TIME IS 3-77333405E+01 SEC.

TIME PNT	U	V	W
1	1.36443243E+00	4.80214770E+00	1.00000000E+00
2	1.40577995E+00	4.2908447E+00	1.00000000E+00
3	1.35449109E+00	4.73627395E+00	1.00000000E+00
4	1.33795097E+00	4.24709718E+00	1.00000000E+00
5	1.31619048E+00	4.14411150E+00	1.00000000E+00
6	1.29709912E+00	4.0400444E+00	1.00000000E+00
7	1.27791452E+00	4.01064334E+00	1.00000000E+00
8	1.26504211E+00	3.95027718E+00	1.00000000E+00
9	1.25251210E+00	3.90741156E+00	1.00000000E+00
10	1.24215502E+00	3.86875754E+00	1.00000000E+00
11	1.23361813E+00	3.84344848E+00	1.00000000E+00
12	1.22642763E+00	3.82907932E+00	1.00000000E+00
13	1.22050524E+00	3.82424108E+00	1.00000000E+00
14	1.21393226E+00	3.82444906E+00	1.00000000E+00
15	1.20774829E+00	3.84134111E+00	1.00000000E+00
16	1.19998430E+00	3.86147713E+00	1.00000000E+00
17	1.19040050E+00	3.89401477E+00	1.00000000E+00
18	1.17848469E+00	3.95444488E+00	1.00000000E+00
19	1.15346622E+00	4.07346220E+00	1.00000000E+00
20	1.10637096E+00	4.2985443E+00	1.00000000E+00
21	1.00969596E+00	4.87774416E+00	1.00000000E+00

TIME IS 0. SEC.

TIME PNT	U	V	W
1	1.422751153E+00	3.73669328E+00	1.00000000E+01
2	1.422725474E+00	3.77772414E+00	1.00000000E+01
3	1.422693597E+00	3.73900045E+00	1.00000000E+01
4	1.422653979E+00	3.74060524E+00	1.00000000E+01
5	1.422604705E+00	3.74259158E+00	1.00000000E+01
6	1.422543390E+00	3.74505704E+00	1.00000000E+01
7	1.422467883E+00	3.74815459E+00	1.00000000E+01
8	1.422373955E+00	3.75194492E+00	1.00000000E+01
9	1.422256764E+00	3.75673459E+00	1.00000000E+01
10	1.422110173E+00	3.76271430E+00	1.00000000E+01
11	1.421924208E+00	3.77024748E+00	1.00000000E+01
12	1.421694368E+00	3.77977110E+00	1.00000000E+01
13	1.421400589E+00	3.79190400E+00	1.00000000E+01
14	1.421025613E+00	3.80750455E+00	1.00000000E+01
15	1.420543223E+00	3.82774429E+00	1.00000000E+01
16	1.419913617E+00	3.85453487E+00	1.00000000E+01
17	1.419075871E+00	3.89066717E+00	1.00000000E+01
18	1.417928145E+00	3.94121002E+00	1.00000000E+01
19	1.416278538E+00	4.01804754E+00	1.00000000E+01
20	1.413665538E+00	4.14024455E+00	1.00000000E+01
21	1.410763468E+00	4.87274427E+00	1.00000000E+01

TIME IS 5-03111207E+01 SEC.

TIME PNT	U	V	W
1	1.31801795E+00	4.77574477E+00	1.00000000E+00
2	1.4963482E+00	4.4500410E+00	1.00000000E+00
3	1.43921743E+00	4.70054956E+00	1.00000000E+00
4	1.4263577E+00	4.5059470E+00	1.00000000E+00
5	1.40027274E+00	4.3005971E+00	1.00000000E+00
6	1.3774943E+00	4.1581253E+00	1.00000000E+00
7	1.35470831E+00	4.3374747E+00	1.00000000E+00
8	1.33221433E+00	4.28135437E+00	1.00000000E+00
9	1.31206383E+00	4.1742417E+00	1.00000000E+00
10	1.29291286E+00	4.1034447E+00	1.00000000E+00
11	1.27531013E+00	4.04044127E+00	1.00000000E+00
12	1.25964440E+00	3.9800747E+00	1.00000000E+00
13	1.24563914E+00	3.9444446E+00	1.00000000E+00
14	1.23314072E+00	3.9237947E+00	1.00000000E+00
15	1.22143404E+00	3.91201484E+00	1.00000000E+00
16	1.20971111E+00	3.91454794E+00	1.00000000E+00
17	1.19868162E+00	3.93029537E+00	1.00000000E+00
18	1.17421516E+00	3.9415570E+00	1.00000000E+00
19	1.15262146E+00	4.1077737E+00	1.00000000E+00
20	1.10467303E+00	4.3274143E+00	1.00000000E+00
21	1.1000437E+00	4.8704746E+00	1.00000000E+00

TIME IS 1-25777802E+01 SEC.

TIME PNT	U	V	W
1	1.42434317E+00	4.00749144E+00	1.00000000E+01
2	1.42423340E+00	3.81004023E+00	1.02965757E+01
3	1.42309333E+00	3.78135739E+00	1.01632011E+01
4	1.42307753E+00	3.76282451E+00	1.00829123E+01
5	1.42281519E+00	3.75372146E+00	1.00493455E+01
6	1.42264343E+00	3.75043137E+00	1.00188642E+01
7	1.42251288E+00	3.75063471E+00	1.0004471E+01
8	1.42232490E+00	3.75307406E+00	1.00035452E+01
9	1.42226261E+00	3.75721198E+00	1.00012404E+01
10	1.42210974E+00	3.76290478E+00	1.00002205E+01
11	1.42192228E+00	3.77031407E+00	0.99968795E+01
12	1.42168781E+00	3.7779425E+00	0.99936699E+01
13	1.42139155E+00	3.7919731E+00	0.99912124E+01
14	1.42101364E+00	3.8075308E+00	0.99885979E+01
15	1.42052689E+00	3.82783171E+00	0.9985630E+01
16	1.41989104E+00	3.85467048E+00	0.99822386E+01
17	1.41904270E+00	3.89104031E+00	0.99792128E+01
18	1.41786893E+00	3.94273126E+00	0.99846978E+01
19	1.41671610E+00	4.02696457E+00	1.00008499E+01
20	1.41218401E+00	4.21721439E+00	1.0033458E+01
21	1.01175211E+00	4.88339499E+00	1.00747613E+01

TIME IS 6-28888900E+01 SEC.

TIME PNT	U	V	W
1	1.03393203E+00	4.7140042E+00	1.00000000E+00
2	1.26439294E+00	4.4714003E+00	1.00000000E+00
3	1.2079356E+00	4.2051728E+00	1.00000000E+00
4	1.21206340E+00	4.0032408E+00	1.00000000E+00
5	1.48481239E+00	4.7645040E+00	1.00000000E+00
6	1.46292400E+00	4.700111E+00	1.00000000E+00
7	1.43451100E+00	4.6340113E+00	1.00000000E+00
8	1.41522990E+00	4.6638133E+00	1.00000000E+00
9	1.39295492E+00	4.4908674E+00	1.00000000E+00
10	1.36589417E+00	4.6130977E+00	1.00000000E+00
11	1.34658655E+00	4.3421414E+00	1.00000000E+00
12	1.32410398E+00	4.2702444E+00	1.00000000E+00
13	1.3023707E+00	4.2011241E+00	1.00000000E+00
14	1.28132962E+00	4.1404019E+00	1.00000000E+00
15	1.26153722E+00	4.1010777E+00	1.00000000E+00
16	1.24183718E+00	4.0722438E+00	1.00000000E+00
17	1.22147844E+00	4.0679408E+00	1.00000000E+00
18	1.19678711E+00	4.1065173E+00	1.00000000E+00
19	1.16451661E+00	4.1979424E+00	1.00000000E+00
20	1.11370489E+00	4.341842E+00	1.00000000E+00
21	1.01589295E+00	4.4973495E+00	1.00000000E+00

TIME IS 2-51555604E+01 SEC.

TIME PNT	U	V	W
1	1.26369703E+00	4.61171407E+00	1.26272130E+01
2	1.30843092E+00	4.0204347E+00	1.15337030E+01
3	1.47991237E+00	4.0115876E+00	1.12272947E+01
4	1.46804725E+00	3.93593454E+00	1.08722862E+01
5	1.45350423E+00	3.88089171E+00	1.06001942E+01
6	1.44380276E+00	3.83384050E+00	1.03955126E+01
7	1.43647535E+00	3.80499917E+00	1.02506293E+01
8	1.43103002E+00	3.7898873E+00	1.01530666E+01
9	1.42640571E+00	3.77944730E+00	1.00902771E+01
10	1.42358082E+00	3.77594459E+00	1.00514414E+01
11	1.4206071E+00	3.7778091E+00	1.00262383E+01
12	1.41761135E+00	3.78400859E+00	1.00147614E+01
13	1.41425813E+00	3.7942937E+00	1.00070489E+01
14	1.4102555E+00	3.80892741E+00	1.0002733E+01
15	1.40525563E+00	3.82877419E+00	1.00002145E+01
16	1.39877501E+00	3.8555938E+00	0.99865818E+01
17	1.3913001E+00	3.89270788E+00	0.9988014E+01
18	1.3775914E+00	3.9473051E+00	0.99873389E+01
19	1.35726467E+00	4.0421575E+00	1.00260889E+01
20	1.3122325E+00	4.26691408E+00	1.00254992E+01
21	1.01042234E+00	4.87974435E+00	1.00492244E+01

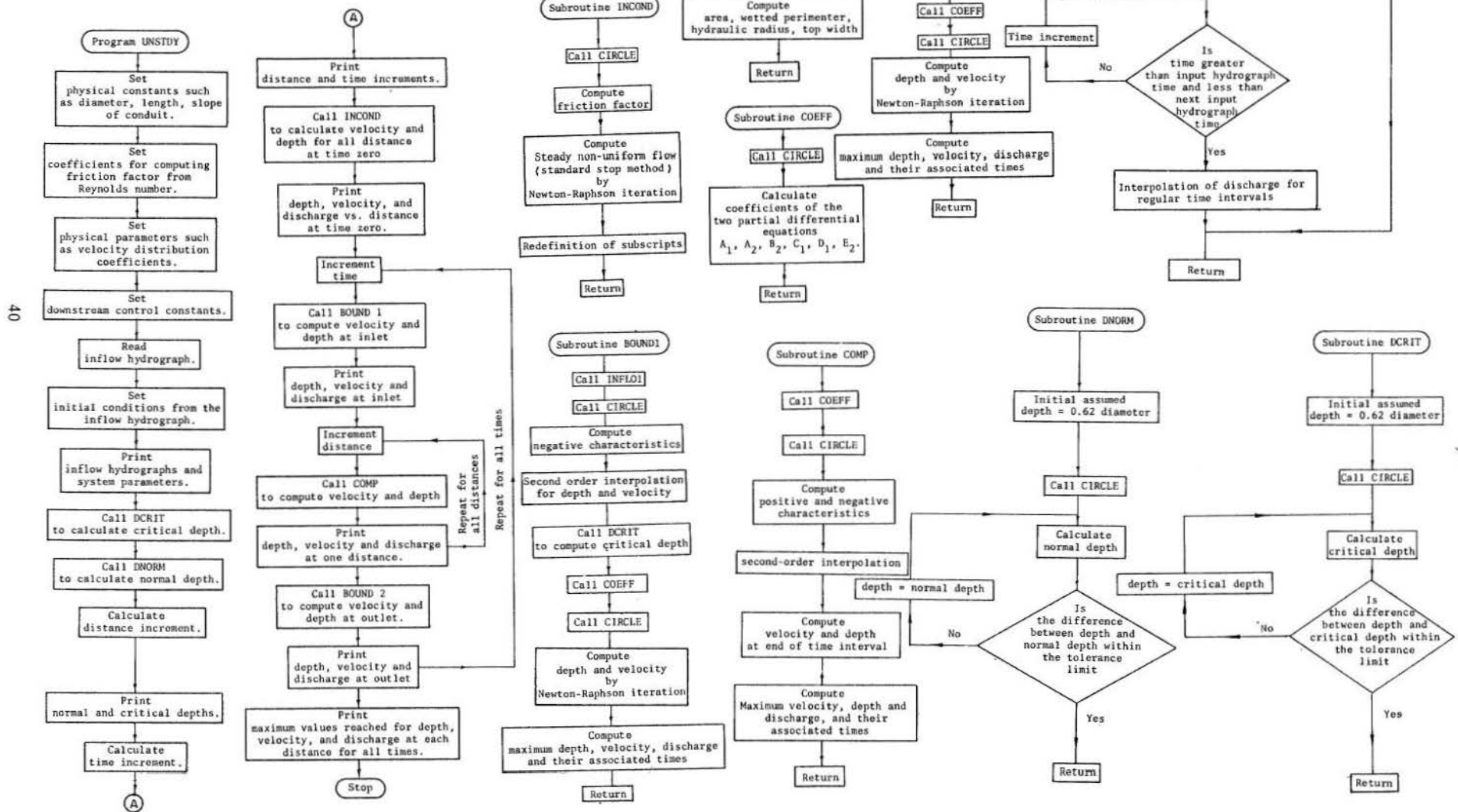
TIME IS 7-54666811E+01 SEC.

TIME PNT	U	V	W
1	1.08086001E+00	4.4244132E+00	1.00000000E+00
2	1.01044890E+00	3.0161391E+00	1.00000000E+00
3	1.58701484E+00	5.0257118E+00	1.00000000E+00
4	1.58056699E+00	4.0274039E+00	1.00000000E+00
5	1.55492461E+00	4.9274474E+00	1.00000000E+00
6	1.53670802E+00	4.4741230E+00	1.00000000E+00
7	1.51547683E+00	4.4171247E+00	1.00000000E+00
8	1.4942517E+00	4.7877147E+00	1.00000000E+00
9	1.47232865E+00	4.7371360E+00	1.00000000E+00
10	1.44986239E+00	4.6817407E+00	1.00000000E+00
11	1.42682953E+00	4.6216242E+00	1.00000000E+00
12	1.40323775E+00	4.601401E+00	1.00000000E+00
13	1.37909414E+00	4.6016937E+00	1.00000000E+00
14	1.35439035E+00	4.4413042E+00	1.00000000E+00
15	1.32904766E+00	4.3851142E+00	1.00000000E+00
16	1.30266633E+00	4.3314427E+00	1.00000000E+00
17	1.27421620E+00	4.1264355E+00	1.00000000E+00
18	1.24138424E+00	4.1104025E+00	1.00000000E+00
19	1.20103247E+00	4.0875031E+00	1.00000000E+00
20	1.14571302E+00	4.3430373E+00	1.00000000E+00
21	1.04321624E+00	4.2004411E+00	1.00000000E+00



APPENDIX 3  
 COMPUTATIONAL DETAILS OF FINITE-DIFFERENCE  
 SPECIFIED INTERVALS SCHEME OF THE METHOD  
 OF CHARACTERISTICS

A.3.1. FLOW CHART



### A.3.2. FORTRAN IV COMPUTER PROGRAM

MAIN PROGRAM FOR UNSTEADY FLOW BY METHOD OF CHARACTERISTICS

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PROGRAM UNSTDY                                UNS 2
1 (INPUT,OUTPUT,TAPE5=INPUT,TAPE6=OUTPUT)    UNS 4
C-----ATTENUATION ANALYSIS - CIRCULAR CROSS SECTION    UNS 6
C-----M1 INITIAL CONDITIONS                            UNS 8
C-----DETERMINATION OF HYDROGRAPH AT THE SPECIFIC POINT WITH TWO CONTROL    UNS 10
C----- ( AT UPSTREAM AND DOWNSHEAM ) BY THE METHOD OF CHARACTERISTICS    UNS 12
C----- FRICTION COEFFICIENT F VARIES WITH REYNOLDS NUMBER    UNS 14
DIMENSION D(500), DDT(500), DMAX(200), Q(500), QDT(500)    UNS 16
DIMENSION Q1(200), QMAX(200)                            UNS 18
DIMENSION TDMAX(200), TQ(200), TQMAX(200), TVMAX(200)    UNS 20
DIMENSION V(500), VDT(500), VMAX(200), X(500)           UNS 22
COMMON A,AR,AC,AD,AE,ALPHA,ALPHA,RC,RD,BETA,CD,CO,DC,DDT    UNS 24
COMMON DEPTH,DD,DIA,DIN,DM,DMAX,DMN,DDDT,DI,DTOL,DX,ED    UNS 26
COMMON F,FB,FC,FD,FNU,GR,I,ITD,ITOC,IXD,IXOC,J,MC,N,NQCD    UNS 28
COMMON NI,N1,Q,OB,QDT,QI,QIN,QMAX,QP,QR,REY,SO,T,TDMAX    UNS 30
COMMON TF,THEIA,TIO,TP,TD,TQMAX,TVMAX,V,VDT,VMAX,VV,WP    UNS 32
COMMON X,XE,XF,XX                                        UNS 34
C-----PHYSICAL CONSTANTS OF THE SYSTEM                UNS 36
DIA=2.9262                                             UNS 38
XF=R21.70                                             UNS 40
SO=0.001                                             UNS 42
C-----COEFFICIENTS FOR COMPUTING F FROM THE REYNOLDS NUMBER    UNS 44
FNU=0.0000141                                         UNS 46
FR=0.109394                                           UNS 48
FC=0.17944                                            UNS 50
C-----PHYSICAL PARAMETERS                              UNS 52
GP=32.175                                             UNS 54
ALPHA=1.00                                            UNS 56
BETA=1.00                                             UNS 58
C-----DOWNSHEAM CONTROL CONSTANTS                    UNS 60
CD=0.0                                                UNS 62
ED=1.35                                              UNS 64
C-----COMPUTATIONAL PARAMETERS                      UNS 66
N=20                                                 UNS 68
IXD=2                                                UNS 70
TF=200.                                             UNS 72
TIO=20.                                              UNS 74
DTOL=0.00001                                         UNS 76
C-----INFLOW HYDROGRAPH                              UNS 78
READ (5,200) NQCD                                     UNS 80
READ (5,210) (TQ(I),QI(I),I=1,NQCD)                 UNS 82
QR=QI(1)                                             UNS 84
QP=QI(2)                                             UNS 86
TP=TQ(2)                                             UNS 88
VOL=(QP-QR)*TP                                       UNS 90
QRA=QR/QP                                            UNS 92
N1=N+1                                              UNS 94
DO 10 I=1,N1                                         UNS 96
DMAX(I)=0.0                                         UNS 98
VMAX(I)=0.0                                         UNS 100
QMAX(I)=0.0                                         UNS 102
10 WRITE (6,220)                                     UNS 104
WRITE (6,270)                                       UNS 106
WRITE (6,230) QR                                     UNS 108
WRITE (6,250) QP                                     UNS 110
WRITE (6,260) TP                                     UNS 112
WRITE (6,320) QRA                                    UNS 114
WRITE (6,240) VOL                                    UNS 116
WRITE (6,270)                                       UNS 118
WRITE (6,280)                                       UNS 120
WRITE (6,290) SO                                     UNS 122
WRITE (6,300) ALPHA                                 UNS 124
WRITE (6,310) BETA                                 UNS 126

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WRITE (6,270)                                       UNS 128
WRITE (6,330) N,IXD,IF,TIO                          UNS 130
C-----COMPUTATION OF NORMAL DEPTH AND CRITICAL DEPTH    UNS 132
QQ=QB                                               UNS 134
CALL DCRTIT                                        UNS 136
CALL DNDRM                                        UNS 138
IF (DN-DC) 20,20,30                                UNS 140
20 WRITE (6,340)                                     UNS 142
GO TO 190                                          UNS 144
30 IF (CD) 40,50,40                                 UNS 146
40 MC=1                                             UNS 148
XX=0.0                                             UNS 150
DDUT=(DR/CD)**(1.0/FD)                             UNS 152
GO TO 60                                           UNS 154
50 MC=2                                             UNS 156
XX=4.5*DC                                          UNS 158
DDUT=DC                                             UNS 160
60 XF=XF-XX                                        UNS 162
AN=N                                               UNS 164
DX=XE/AN                                           UNS 166
WRITE (6,350) DN,DC                                UNS 168
C-----COMPUTATION OF DT ( TIME INCREMENT)             UNS 170
QQ=QP                                              UNS 172
CALL DCRTIT                                        UNS 174
DFPTH=DC                                           UNS 176
CALL CIRCLE                                        UNS 178
VC=QP/A                                             UNS 180
DFPTH=0.87*PIA                                     UNS 182
CALL CIRCLE                                        UNS 184
DTMAX=(DX**2.0*BETA)/(VC*(ALPHA+BETA)+SQRT(((ALPHA-BETA)**2)*VC*VC))    UNS 186
I=(4.0*BETA*GR*DM))                               UNS 188
DT=DTMAX*.5                                        UNS 190
CO=-DT/DX                                          UNS 192
NT=TF/DT                                           UNS 194
ITD=ITD/DT                                         UNS 196
WRITE (6,360) DX,DT                                 UNS 198
C-----COMPUTATION OF VELOCITY AND DEPTH FOR ALL DISTANCES X AT TIME 0.0    UNS 200
DIN=(DDUT+DN)/2.0                                  UNS 202
DFPTH=DDUT                                         UNS 204
CALL INCOND                                        UNS 206
QIN=QR                                             UNS 208
T=TQ(1)                                             UNS 210
WRITE (6,370) T                                     UNS 212
WRITE (6,380)                                       UNS 214
WRITE (6,390)                                       UNS 216
DO 70 I=1,N1,IXD                                   UNS 218
WRITE (6,400) X(I),DDI(I),V(I),Q(I)               UNS 220
CONTINUE                                           UNS 222
70 ITDC=I                                           UNS 224
DO 170 J=2,NT                                       UNS 226
T=T+DT                                             UNS 228
C-----COMPUTATION OF VELOCITY AND DEPTH FOR THE INLET AT TIME T    UNS 230
CALL ROUND1                                        UNS 232
IF (ITDC-ITD) 90,90,90                             UNS 234
80 WRITE (6,370) T                                  UNS 236
WRITE (6,380)                                       UNS 238
WRITE (6,390)                                       UNS 240
WRITE (6,400) X(I),DDI(I),VDT(I),QDT(I)          UNS 242
90 IXOC=I                                           UNS 244
DO 120 I=2,N                                       UNS 246
C-----COMPUTATION OF VELOCITY AND DEPTH AT TIME T    UNS 248
CALL COMP                                          UNS 250
IF (ITOC,EO,ITD,AND,IXOC,FQ,IXD) 100,110          UNS 252
100 IXOC=I                                          UNS 254
WRITE (6,400) X(I),DDI(I),VDT(I),QDT(I)          UNS 256
GO TO 120                                          UNS 258
110 IXOC=IXOC+1                                     UNS 260
120 CONTINUE                                        UNS 262
C-----COMPUTATION OF VELOCITY AND DEPTH FOR THE OUTLET AT TIME T    UNS 264
I=N1                                              UNS 266
CALL ROUND2                                        UNS 268
IF (ITOC-ITD) 140,130,140                          UNS 270
130 ITOC=I                                          UNS 272
WRITE (6,400) X(I),DDI(I),VDT(I),QDT(I)          UNS 274
GO TO 150                                          UNS 276
140 ITOC=ITOC+1                                     UNS 278
150 DO 160 I=1,N1                                   UNS 280
Q(I)=QDT(I)                                       UNS 282

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D(I)=DDT(I)          UNS 284
160 V(I)=VDT(I)      UNS 286
170 CONTINUE          UNS 288
    NPG=NI/50+1      UNS 290
    DO 180 I=1,NPG   UNS 292
    IT=50*I-49       UNS 294
    IL=I+49          UNS 296
    WRITE (6,410)    UNS 298
    WRITE (6,420)    UNS 300
    DO 180 I=I,IL    UNS 302
    WRITE (6,430) X(I),DMAX(I),TOMAX(I),VHMAX(I),TVMAX(I),QMAX(I),TOMAX
1(I)                UNS 306
    IF (I.EQ.NI) GO TO 190
180 CONTINUE          UNS 308
190 CALL EXIT        UNS 310
C-----            UNS 312
200 FORMAT (I3)      UNS 314
210 FORMAT (8F10.4)  UNS 316
220 FORMAT (1H)2X,*INFLOW HYDROGRAPH PARAMETERS*/      UNS 318
230 FORMAT (2X,*QB= *F10.5,*CFS*)                       UNS 320
240 FORMAT (2X,*WAVE VOLUME ABOVE BASE FLOW= *F8.2,*CU FT*) UNS 322
250 FORMAT (2X,*QP= *F10.5,*CFS*)                       UNS 324
260 FORMAT (2X,*TP= *F10.5,*SEC*)                       UNS 326
270 FORMAT (/)      UNS 328
280 FORMAT (2X,*SYSTEM PARAMETERS*/                     UNS 330
290 FORMAT (* SD =*F10.5)                                UNS 332
300 FORMAT (* ALPHA=*F10.5)                              UNS 334
310 FORMAT (* BETA=*F10.5)                              UNS 336
320 FORMAT (* QB/QP=*F10.5)                              UNS 338
330 FORMAT (* N =*15/* IX0 =*15/* IF =*F6.0/* T10 =*F10.5) UNS 340
340 FORMAT (* FLOW IS SUPERCRITICAL*)                   UNS 342
350 FORMAT (2X,*NORMAL DEPTH= *F6.4,*FT*4X,*CRITICAL DEPTH= *F6.4,*FNS
IT*/)            UNS 344
360 FORMAT (2X,*DX= *F8.5,*FT*4X,*DT= *F8.5,*SEC*/)    UNS 346
370 FORMAT (1H)5X,*CONDITIONS AT *9.3*SECONDS*/        UNS 348
380 FORMAT (2X,*DISTANCE*9X,*DEPTH*8X,*VELOCITY*
1 7X,*DISCHARGE*)  UNS 350
390 FORMAT (4X,*FT)*11X,*(FT)*10X,*(FPS)*11X,*(CFS)*1/  UNS 352
400 FORMAT (4(F10.4,5X))                                UNS 354
410 FORMAT (//)*1 MAXIMUM VALUES AND TIMES AT EACH LOCATION*/ UNS 356
420 FORMAT (* DISTANCE MAX DEPTH TIME MAX VEL TIME MAX Q
1 TIME*)         UNS 358
430 FORMAT (F8.2,3(4X,F6.2,2X,F7.2))                   UNS 360
END                                                       UNS 370

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SUBROUTINE FOR COMPUTING INITIAL CONDITION

```

SUBROUTINE INCOND          INC 2
DIMENSION D(500), DDT(500), DMAX(200), Q(500), QDT(500)  INC 4
DIMENSION QI(200), QMAX(200)                                INC 6
DIMENSION TOMAX(200), TQ(200), TOMAX(200), TVMAX(200)    INC 8
DIMENSION V(500), VDI(500), VMAX(200), X(500)            INC 10
COMMON A,AR,AC,AD,AE,ALPHA,BC,BC,HD,BETA,CO,CO,DU,DC,DDT  INC 12
COMMON DEPTH,DD,DIA,DIN,DM,DMAX,DN,DOU,DT,DTOL,DX,ED      INC 14
COMMON F,FR,FC,FD,FNU,GR,I,ITO,ITOC,IX0,IXOC,J,MC,N,NQCD  INC 16
COMMON NT,N1,Q,QR,ODT,DI,QIN,QMAX,QP,QQ,H,PEY,SD,T,TDMAX  INC 18
COMMON TF,THETA,T10,TP,TQ,TOMAX,TVMAX,V,VDI,VMAX,VV,WP    INC 20
COMMON X,XF,XF,XX     INC 22
CALL CIRCLF          INC 24
C CONDITION AT INITIAL POSITION                               INC 26
VV=QB/A              INC 28
VH=VV*VV/(2.0*GR)  INC 30
C COMPUTE REYNOLDS NUMBER                                  INC 32
RFY=VV*R/FNU        INC 34
C COMPUTE FRICTION FACTOR                                 INC 36
F=FB*REY**FC        INC 38
S1=F*VH/(4.0*R)    INC 40
EE1=DEPTH+ALPHA*VH INC 42
X(N+1)=XF-XX       INC 44
D(N+1)=DDIT        INC 46
V(N+1)=VV           INC 48
Q(N+1)=QB           INC 50
NCOUNT=0          INC 52

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C INTEGRATION OF STEADY FLOW                               INC 54
DO 150 L=1,N        INC 56
XX=XX+DX            INC 58
DFPTH=DN           INC 60
CALL CIRCLF        INC 62
VV=QB/A            INC 64
RFY=VV*R/FNU      INC 66
F=FB*REY**FC      INC 68
HFTH=0.5*THETA    INC 70
DTHET=4.0/(DIA*STNF(HFTH)) INC 72
DAREA=0.125*DIA*DTA*(1.0-COSF(THETA))*DTHET INC 74
DW=0.5*DIA*DTHET  INC 76
WP=0.5*DIA*THETA  INC 78
DRA=(WP*DAREA-ADQ)/(WP*HP) INC 80
DFNG=1.0-(QR*QB/(GR*(A**3)))*DAREA INC 82
DSL0=-F*QR*QB*(2.0*R*A*DAREA*(A**2)*DVA)/(R,GR*(R*A**2)**2) INC 84
VV=QB/A            INC 86
VH=VV*VV/(2.0*GR) INC 88
RFY=VV*R/FNU      INC 90
F=FB*REY**FC      INC 92
S2=F*VH/(4.0*R)   INC 94
SF=(S1+S2)/2.0    INC 96
EF2=DIN+ALPHA*VH  INC 98
FDATIO=(EF2-EF1+DX*(S0-SF))/(DFNG*(EF2-EF1)*DSL0/(S0-SF)) INC 100
C NEWTON-RAPHSON ITERATION                               INC 102
DCOM=DIN-FRATIO    INC 104
IF (DCOM) 30,20,40 INC 106
WRITE (6,200)      INC 108
GO TO 190          INC 110
30 DCOM=ABSF(DCOM) INC 112
IF (ABSF(DCOM-DIN)-DTOL) 130,130,50 INC 114
IF (0.2*DIA-DCOM) 60,120,120 INC 116
60 DIN=DCOM*0.5    INC 118
70 IF (0.2*DIA-DIN) 80,80,90 INC 120
80 DIN=DIN*0.5     INC 122
NCOUNT=NCOUNT+1   INC 124
GO TO 70           INC 126
90 IF (NCOUNT-20) 100,100,110 INC 128
100 GO TO 10       INC 130
110 WRITE (6,210)  INC 132
GO TO 190          INC 134
120 DIN=DCOM       INC 136
GO TO 10           INC 138
130 IF (ABSF(DCOM-DN),L.F,DTOL) 160,140 INC 140
C END OF NEWTON-RAPHSON                                  INC 142
DIN=DCOM           INC 144
S1=S2              INC 146
EF1=EF2           INC 148
C REDEFINITION OF SUBSCRIPTS                             INC 150
II=N-L+1          INC 152
X(II)=XF-XX       INC 154
D(II)=DIN         INC 156
V(II)=VV          INC 158
Q(II)=QB          INC 160
150 CONTINUE      INC 162
GO TO 180         INC 164
160 DFPTH=DN      INC 166
CALL CIRCLF       INC 168
VV=QB/A           INC 170
C CONSTANT CONDITIONS                                    INC 172
DO 170 J=L,N      INC 174
II=N-J+1          INC 176
X(II)=XF-XX       INC 178
D(II)=DN          INC 180
V(II)=VV          INC 182
Q(II)=QB          INC 184
XX=XX+DX          INC 186
170 CONTINUE      INC 188
180 RETURN        INC 190
190 CALL EXIT     INC 192
C-----            INC 194
200 FORMAT (* DCOM EQUALS ZERO *)                         INC 196
210 FORMAT (* INCOND DOES NOT CONVERGE*)                 INC 198
END                                                       INC 200

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SUBROUTINE FOR COMPUTING DEPTH & VELOCITY AT END OF TIME INTERVAL

```

SUBROUTINE COMP
C-----COMPUTATION OF VELOCITY AND DEPTH AT THE TIME T=DT BY KNOWING THE
C-----VELOCITY AND THE DEPTH AT THE TIME T
DIMENSION D(500), DDT(500), DMAX(200), Q(500), QDT(500)
DIMENSION QI(200), QMAX(200)
DIMENSION TDMAX(200), TQ(200), TQMAX(200), TVMAX(200)
DIMENSION V(500), VDT(500), VMAX(200), X(500)
COMMON A,AR,AC,AD,AE,ALPHA,H,HC,HD,HEIA,CI,CO,DC,DDT
COMMON DEPTH,DD,DIA,DIN,DM,DMAX,DN,DDUT,DT,DTOL,DX,FD
COMMON F,FB,FC,FD,FNU,GR,I,ITO,ITOC,IXO,IXOC,J,MC,N,NQCD
COMMON NT,N1,Q,QR,QDT,OI,GIN,QMAX,OP,OD,PREY,SO,T,TDMAX
COMMON TF,THEIA,TIO,TP,TQ,TQMAX,TVMAX,V,VDT,VMAX,VV,WP
COMMON X,XE,XF,XX
DD=D(I)
VV=V(I)
CALL COEFF
DEPTH=D(I)
CALL CIRCLE
C POSITIVE CHARACTERISTIC
CP=(2.0*BETA)/(V(I)*(ALPHA+HEIA)+SQRT(((ALPHA-BETA)**2)*V(I)*V(I)
+ (4.0*A*BFIA*GR/R)))
C NEGATIVE CHARACTERISTIC
CN=(2.0*BETA)/(V(I)*(ALPHA+HEIA)-SQRT(((ALPHA-BETA)**2)*V(I)*V(I)
+ (4.0*A*BFIA*GR/R)))
UP=CO/CP
UN=CO/CN
C 2ND ORDER INTERPOLATION
DR=D(I-1)*0.5*UP*(UP-1.)+D(I)*(1.-UP**2)+D(I+1)*0.5*UN*(UP+1.)
VR=V(I-1)*0.5*UP*(UP-1.)+V(I)*(1.-UP**2)+V(I+1)*0.5*UN*(UP+1.)
DS=D(I-1)*0.5*UN*(UN-1.)+D(I)*(1.-UN**2)+D(I+1)*0.5*UN*(UN+1.)
VS=V(I-1)*0.5*UN*(UN-1.)+V(I)*(1.-UN**2)+V(I+1)*0.5*UN*(UN+1.)
FCP=AC*CP-HC
FCN=AC*CN-HC
GCP=AR
GCM=AR
SCP=AF*CP
SCN=AE*CN
TCP=FCP*DR+GCP*VR-SCP*DI/CP
TCN=FCN*DS+GCM*VS-SCN*DI/CN
C VELOCITY AND DEPTH AT END OF TIME INTERVAL
VP=(FCP*TCP-FCN*TCN)/(FCP*GCM-FCN*GCP)
DP=(TCP*GCM-TCN*GCP)/(FCP*GCM-FCN*GCP)
IF (DP-0.02*DI) 20,10,10
WRITE (6,90) X(I),T
GO TO 80
20 DEPTH=DP
CALL CIRCLE
QDT(I)=A*VP
DDT(I)=DP
VDT(I)=VP
C MAX DEPTH, VELOCITY, AND DISCHARGE, AND THEIR ASSOCIATED TIMES
IF (DDT(I)-DMAX(I)) 40,40,30
30 DMAX(I)=DDT(I)
TDMAX(I)=T
40 IF (VDT(I)-VMAX(I)) 60,60,50
50 VMAX(I)=VDT(I)
TVMAX(I)=T
60 IF (QDT(I)-QMAX(I)) 80,80,70
70 QMAX(I)=QDT(I)
TQMAX(I)=T
80 RETURN
C-----
90 FORMAT (* FLOW IS FULL AT X = *F./2.* T = *T/2.)
END

```

SUBROUTINE FOR COMPUTING COEFFICIENTS IN ORDINARY DIFFERENTIAL EQUATIONS

```

SUBROUTINE COEFF
C-----COMPUTATION OF ALL COEFFICIENTS OF THE TWO PARTIAL DIFFERENTIAL
C-----EQUATIONS
DIMENSION D(500), DDT(500), DMAX(200), Q(500), QDT(500)
DIMENSION QI(200), QMAX(200)
DIMENSION TDMAX(200), TQ(200), TQMAX(200), TVMAX(200)
DIMENSION V(500), VDT(500), VMAX(200), X(500)
COMMON A,AR,AC,AD,AE,ALPHA,H,HC,HD,HEIA,CI,CO,DC,DDT
COMMON DEPTH,DD,DIA,DIN,DM,DMAX,DN,DDUT,DT,DTOL,DX,ED
COMMON F,FB,FC,FD,FNU,GR,I,ITO,ITOC,IXO,IXOC,J,MC,N,NQCD
COMMON NT,N1,Q,QR,QDT,OI,GIN,QMAX,OP,OD,PREY,SO,T,TDMAX
COMMON TF,THEIA,TIO,TP,TQ,TQMAX,TVMAX,V,VDT,VMAX,VV,WP
COMMON X,XE,XF,XX
DEPTH=DD
CALL CIRCLE
A)=A/(VV*R)
D)=1.0/VV
A2=ALPHA*VV/GR
B2=BETA/GR
C REYNOLDS NUMBER
REY=VV*R/FNU
C FRICTION FACTOR
F=FB*REY**FC
C ENERGY SLOPE
SF=.125*F*VV*VV/(R*GR)
E2=SF-SO
AR=A1*R2
AC=A1-A2
AD=A2*D1
AE=A1*E2
BC=B2
BD=-B2*D1
RETURN
END

```

SUBROUTINE FOR COMPUTING NORMAL DEPTH

```

SUBROUTINE DNDRM
DIMENSION D(500), DDT(500), DMAX(200), Q(500), QDT(500)
DIMENSION QI(200), QMAX(200)
DIMENSION TDMAX(200), TQ(200), TQMAX(200), TVMAX(200)
DIMENSION V(500), VDT(500), VMAX(200), X(500)
COMMON A,AR,AC,AD,AE,ALPHA,H,HC,HD,HEIA,CI,CO,DC,DDT
COMMON DEPTH,DD,DIA,DIN,DM,DMAX,DN,DDUT,DT,DTOL,DX,FD
COMMON F,FB,FC,FD,FNU,GR,I,ITO,ITOC,IXO,IXOC,J,MC,N,NQCD
COMMON NT,N1,Q,QR,QDT,OI,GIN,QMAX,OP,OD,PREY,SO,T,TDMAX
COMMON TF,THEIA,TIO,TP,TQ,TQMAX,TVMAX,V,VDT,VMAX,VV,WP
COMMON X,XE,XF,XX
DEPTH=D,62*DI
CALL CIRCLE
VV=Q/A
C REYNOLDS NUMBER
REY=VV*R/FNU
C FRICTION FACTOR
F=FB*REY**FC
C NEWTON-RAPHSON
DN=DEPTH-(WP-(F*(Q**2)/(H.*G**50*(R**2)*A))/(13.0*H)/H-2.0/SINF
+THEIA/2.0)
IF (ABS(DN-DEPTH)-DTOL) 30,20,20
20 DEPTH=DN
GO TO 10
30 RETURN
END

```

## SUBROUTINE FOR COMPUTING UPSTREAM BOUNDARY CONDITION.

```

SUBROUTINE BOUND1
C-----COMPUTATION OF VELOCITY AND DEPTH FOR X=0.0 AT THE TIME T
DIMENSION D(500), QDT(500), DMAX(200), Q(500), QDT(500)
DIMENSION DI(200), DMAX(200)
DIMENSION TDMAX(200), TO(200), TQMAX(200), TVMAX(200)
DIMENSION V(500), VDT(500), VMAX(200), X(500)
COMMON A,AB,AC,AD,AE,ALPHA,H,RC,HD,RE(A,CO),CO,D,DC,DDT
COMMON DEPTH,DD,DIA,DIN,DM,DMAX,DN,DOUJ,DT,DTOL,DX,ED
COMMON F,FB,FC,FD,FNU,GR,I,ITO,IIOC,IXD,IXOC,J,MC,N,NQCD
COMMON NT,N1,Q,QR,QDT,QI,QIN,DMAX,OP,OU,R,REY,SU,T,TDMAX
COMMON TF,THETA,TIO,TP,TQ,TQMAX,TVMAX,V,VDT,VMAX,VV,WP
COMMON X,XE,XF,XX
CALL INFL01
DEPTH=D(1)
CALL CIRCLE
C NEGATIVE CHARACTERISTIC
CM=(2.0*HETA)/(V(1)*(ALPHA+HETA))-SQRTF(((ALPHA-HETA)**2)*V(1)*V(1)
+ (4.0*A*BETA*GR/R))
IF (CM) 10,20,30
10 UN=CO/CM
C 2ND ORDER INTERPOLATION FOR DEPTH AND VELOCITY
DS=D(1)*0.5*UN*(UN-1.)+D(2)*(1.-UN**2)+D(3)*0.5*UN*(UN+1.)
VS=V(1)*0.5*UN*(UN-1.)+V(2)*(1.-UN**2)+V(3)*0.5*UN*(UN+1.)
GO TO 40
20 XS=X(1)
DS=D(1)
VS=V(1)
GO TO 40
30 QQ=QIN
CALL DCRIT
DP=DC
GO TO 80
40 DP=D(1)
VV=V(1)
CALL COEFF
FCM=AC*CM-RC
GCM=AB
SCM=AE*CM
ASMALL=DS-(SCM*CM*DT-GCM*VS)/FCM
BSMALL=-QIN*GCM/FCM
DP1=D(1)
50 RD=2.0*DP1/DIA-1.0
DEPTH=DP1
CALL CIRCLE
FDP1=DP1-ASMALL-(BSMALL/A)
FDP1P=1.0*(BSMALL/A**2)*((DIA*(1.0-COSF(THETA))/2.0)**(1.0/SQRTF(1.0-RD**2)))
C NEWTON-RAPHSON ITERATION
DP2=DP1-FDP1/FDP1P
IF (ABS(FDP2-DP1)-DTOL) 70,70,60
60 DP1=DP2
GO TO 50
C END OF NEWTON-RAPHSON
70 DP=DP1
80 IF (DP-0.82*DIA) 100,90,90
90 WRITE (6,170) X(1),T
GO TO 160
100 DEPTH=DP
CALL CIRCLE
VP=QIN/A
DNT(1)=DP
VNT(1)=VP
QNT(1)=QIN
C MAX DEPTH, VELOCITY, AND DISCHARGE, AND THEIR ASSOCIATED TIMES
IF (DDT(1)-DMAX(1)) 120,120,110
110 DMAX(1)=DDT(1)
TDMAX(1)=T
120 IF (VDT(1)-VMAX(1)) 140,140,130
130 VMAX(1)=VDT(1)
TVMAX(1)=T
140 IF (QDT(1)-QMAX(1)) 160,160,150
150 QMAX(1)=QDT(1)
TQMAX(1)=T
160 RETURN
C-----
170 FORMAT (' FLOW IS FULL AT X = ',F7.2,' T = ',F6.2)
END

```

## SUBROUTINE FOR COMPUTING DOWNSTREAM BOUNDARY CONDITION.

```

SUBROUTINE BOUND2
DIMENSION D(500), QDT(500), DMAX(200), Q(500), QDT(500)
DIMENSION QI(200), DMAX(200)
DIMENSION TDMAX(200), TQ(200), TQMAX(200), TVMAX(200)
DIMENSION V(500), VDT(500), VMAX(200), X(500)
COMMON A,AB,AC,AD,AF,ALPHA,H,HC,HD,HETA,CO,DC,DD,DC,DDT
COMMON DEPTH,DD,DIA,DIN,DM,DMAX,DN,DOUJ,DT,DTOL,DX,ED
COMMON F,FB,FC,FD,FNU,GR,I,ITO,IIOC,IXD,IXOC,J,MC,N,NQCD
COMMON NT,N1,Q,QR,QDT,QI,QIN,DMAX,OP,OU,R,REY,SU,T,TDMAX
COMMON TF,THETA,TIO,TP,TQ,TQMAX,TVMAX,V,VDT,VMAX,VV,WP
COMMON X,XE,XF,XX
DEPTH=D(N1)
CALL CIRCLE
C POSITIVE CHARACTERISTIC
CP=(2.0*BETA)/(V(N1)*(ALPHA+BETA))+SQRTF(((ALPHA-BETA)**2)*V(N1)*V(N1)
+ (4.0*A*BETA*GR/R))
UP=CO/CP
C 2ND ORDER INTERPOLATION FOR DEPTH AND VELOCITY
DR=D(N-1)*0.5*UP*(UP-1.)+D(N)*(1.-UP**2)+D(N1)*0.5*UP*(UP+1.)
VR=V(N-1)*0.5*UP*(UP-1.)+V(N)*(1.-UP**2)+V(N1)*0.5*UP*(UP+1.)
DN=D(N1)
VV=V(N1)
CALL COEFF
FCP=AC*CP-BC
GCP=AB
SCP=AE*CP
CSMALL=DR-(SCP*CP*DT-GCP*VV)/FCP
DSMALL=-GCP/FCP
DP1=D(N1)
10 RD=DP1*2.0/DIA-1.0
DEPTH=DP1
CALL CIRCLE
GO TO (20,30), MC
20 FN=CD*DP1**ED
FD1=CD*ED*DP1** (ED-1.0)
U=FD/A
FDP1=DP1-CSMALL-DSMALL*U
THETA2=THETA/2.0
DADD=(DIA/2.0)*(1.0-COSF(THETA))*(1.0/SQRTF(1.0-RD**2))
DIND=((A*FD1)-(FD*DADD))/(A*A)
GO TO 40
30 U=SQRTF(GR*A/R)
FDP1=DP1-CSMALL-DSMALL*U
THETA2=THETA/2.0
DIND=(2./DIA)*(1.0/SQRTF(1.0-RD**2))*(1.0/U)*((U/A)**2*(1.0-COSF(THETA2)/
(A*0*B)-(A*(DIA/2.0)*COSF(THETA2)/H**2)))
FDP1P=1.0-DSMALL*DIND
C NEWTON-RAPHSON ITERATION
DP2=DP1-FDP1/FDP1P
IF (ABS(FDP2-DP1)-DTOL) 60,60,50
50 DP1=DP2
GO TO 10
C END OF NEWTON-RAPHSON
60 DEPTH=DP2
IF (DEPTH-0.82*DIA) 80,70,70
70 WRITE (6,180) X(1),T
GO TO 170
80 CALL CIRCLE
DDT(N1)=DEPTH
GO TO (90,100), MC
90 QDT(N1)=CP*DEPTH**EN
VNT(N1)=QNT(N1)/A
GO TO 110
100 VNT(N1)=SQRTF(GR*A/R)
QNT(N1)=VNT(N1)*A
C MAX DEPTH, VELOCITY, AND DISCHARGE, AND THEIR ASSOCIATED TIMES
IF (DDT(N1)-DMAX(N1)) 130,130,120
120 DMAX(N1)=DDT(N1)
TDMAX(N1)=T
130 IF (VDT(N1)-VMAX(N1)) 150,150,140
140 VMAX(N1)=VDT(N1)
TVMAX(N1)=T
150 IF (QDT(N1)-QMAX(N1)) 170,170,160
160 QMAX(N1)=QDT(N1)
TQMAX(N1)=T
170 RETURN
C-----
180 FORMAT (' FLOW IS FULL AT X = ',F7.2,' T = ',F6.2)
END

```

SUBROUTINE FOR COMPUTING GEOMETRIC PARAMETERS OF CIRCULAR SEGMENT

```

SUBROUTINE CIRCLE
DIMENSION D(500), DDT(500), DMAX(200), Q(500), QDT(500)
DIMENSION QI(200), QMAX(200)
DIMENSION TDMAX(200), TQ(200), TQMAX(200), TVMAX(200)
DIMENSION V(500), VDT(500), VMAX(200), X(500)
COMMON A,AR,AC,AD,AE,AF,APHA,PH,PC,PD,PF,TA,CD,CO,DO,DC,DDT
COMMON DEPTH,DD,DIA,DIN,DM,DMAX,DN,DDDT,DI,DTOL,DX,ED
COMMON F,FR,FC,FD,FNU,GR,I,ITD,ITOC,IXU,IXUC,J,MC,N,NQCD
COMMON NT,NI,Q,QB,QDT,QI,QIN,QMAX,OP,OO,R,PEY,SQ,T,TDMAX
COMMON TF,THETA,TID,TP,TD,TQMAX,TVMAX,V,VDT,VMAX,VV,WP
COMMON X,XE,XF,XX
C
TEST TO INSURE DEPTH LESS THAN 0.82 DIA.
IF (DEPTH) 10,20,20
10 WRITE (6,100)
CALL EXIT
20 IF (DEPTH=0.82*DIA) 40,40,30
30 WRITE (6,110)
CALL EXIT
40 IF (DIA/2.0-DEPTH) 60,40,70
50 THETA=7.14159
GO TO 90
C
SURTENDED ANGLE
60 THETA=6.28318-2.0*ATANF((SQRT(DIA*DEPTH-DEPTH*DEPTH))/(DEPTH-DIA/
12.0))
GO TO 90
70 THETA=2.0*ATANF((SQRT(DIA*DEPTH-DEPTH*DEPTH))/(DIA/2.0-DEPTH))
IF (THETA) 80,90,90
80 THETA=6.28318+THETA
C
ARFA
90 A=0.125*(THETA-SINF(THETA))*DIA**2)
C
WETTED PERIMETER
WP=(DIA/2.0)*THETA
C
HYDRAULIC RADIUS
R=A/WP
C
SURFACE WIDTH
B=DIA*SINF(THETA/2.0)
C
HYDRAULIC DEPTH
DM=A/R
RETURN
C-----
100 FORMAT (* DEPTH IS NEGATIVE*)
110 FORMAT (* FLOW IS FULL*)
END

```

SUBROUTINE FOR COMPUTING CRITICAL DEPTH

```

SUBROUTINE DCHIT
DIMENSION D(500), DDT(500), DMAX(200), Q(500), QDT(500)
DIMENSION QI(200), QMAX(200)
DIMENSION TDMAX(200), TQ(200), TQMAX(200), TVMAX(200)
DIMENSION V(500), VDT(500), VMAX(200), X(500)
COMMON A,AR,AC,AD,AE,AF,APHA,PH,PC,PD,PF,TA,CD,CO,DO,DC,DDT
COMMON DEPTH,DD,DIA,DIN,DM,DMAX,DN,DDDT,DI,DTOL,DX,ED
COMMON F,FR,FC,FD,FNU,GR,I,ITD,ITOC,IXU,IXUC,J,MC,N,NQCD
COMMON NT,NI,Q,QB,QDT,QI,QIN,QMAX,OP,OO,R,PEY,SQ,T,TDMAX
COMMON TF,THETA,TID,TP,TD,TQMAX,TVMAX,V,VDT,VMAX,VV,WP
COMMON X,XE,XF,XX
DEPTH=0.62*DIA
10 CALL CIRCLE
C
NEWTON-RAPHSON
DC=DEPTH-(R*(A**3)-ALPHA*(Q**2)/Q)/(1.0*(18*A)**2)-(2.0*(A**DCR
13)*COSF(THETA/2.0)/(SINF(THETA/2.0)))
IF (ABSF(DC-DEPTH)-DTOL) 30,20,20
20 DEPTH=DC
GO TO 10
30 RETURN
END

```

SUBROUTINE FOR COMPUTING INFLOW HYDROGRAPH

```

SUBROUTINE INFLO1
C-----COMPUTATION OF THE INFLOW HYDROGRAPH
DISCHARGES AT IRREGULAR TIME INTERVALS
DIMENSION D(500), DDT(500), DMAX(200), Q(500), QDT(500)
DIMENSION QI(200), QMAX(200)
DIMENSION TDMAX(200), TQ(200), TQMAX(200), TVMAX(200)
DIMENSION V(500), VDT(500), VMAX(200), X(500)
COMMON A,AR,AC,AD,AE,AF,APHA,PH,PC,PD,PF,TA,CD,CO,DO,DC,DDT
COMMON DEPTH,DD,DIA,DIN,DM,DMAX,DN,DDDT,DI,DTOL,DX,ED
COMMON F,FR,FC,FD,FNU,GR,I,ITD,ITOC,IXU,IXUC,J,MC,N,NQCD
COMMON NT,NI,Q,QB,QDT,QI,QIN,QMAX,OP,OO,R,PEY,SQ,T,TDMAX
COMMON TF,THETA,TID,TP,TD,TQMAX,TVMAX,V,VDT,VMAX,VV,WP
COMMON X,XE,XF,XX
AJ=J
T=(AJ-1.0)*DT
ID=1
C
INTERPOLATION FOR REGULAR TIME INTERVALS
IF (T.GE.TQ(NQCD)) 10,20
10 QIN=QI(NQCD)
GO TO 50
20 IF (T.GE.TQ(IQ),AND,T,IT,TQ(IQ+1)) 40,30
30 ID=IQ+1
GO TO 20
40 QIN=QI(IQ)+(QI(IQ+1)-QI(IQ))*(T-TQ(IQ))/(TQ(IQ+1)-TQ(IQ))
GO TO 50
50 RETURN
END

```

### A.3.3. DEFINITION OF VARIABLES

NAME	DEFINITION	STATEMENT NUMBER(S)
A	AREA OF CIRCULAR SEGMENT	C1P 60
AR	(2)	COF 54
AC	(2)	COE 54
AD	(2)	COE 58
AE	(2)	COF 60
AJ	(1)	INF 28
ALPHA	VFL DISTRIBUTION FACTOR-ENERGY	UNS 56
AN	NUMBER OF DISTANCE INTERVALS	UNS 164
ASMLL	(2)	HO1 78
A1	(2)	COE 32
A2	(2)	COE 36
R	FREE SURFACE WIDTH	C1P 72
BC	(2)	COF 62
RD	(2)	COF 64
BETA	VFL DISTRIBUTION FACTOR-MOMENTUM	UNS 58
BSMLL	(2)	HO1 80
B2	(2)	COE 38
CD	OUTLET DISCHARGE COEFFICIENT	UNS 62
CM	NEGATIVE CHARACTERISTIC DIRECTION	HO1 34 COM 46
CO	MINUS DT/DX	UNS 192
CP	POSITIVE CHARACTERISTIC DIRECTION	COM 40 HO2 30
CSMLL	(2)	HO2 54
D	DEPTH OF FLOW AT TIME T	INC 46 INC 156 INC 180 UNS 284
DADD	DERIVATIVE OF AREA WITH DEPTH	HO2 78
DAREA	DERIVATIVE OF AREA WITH DEPTH	INC 74
DC	CRITICAL DEPTH	DCR 30
DCOM	COMPUTED DEPTH	INC 104 INC 112
DD	DEPTH	HO1 66 HO2 42 COM 28
DDT	DEPTH OF FLOW AT TIME T+DT	HO1 122 HO2 118 COM 98
DENG	(2)	INC 82
DEPTH	DEPTH OF FLOW	HO1 28 HO1 86 HO1 116 COE 28 HO2 24 HO2 62 HO2 108 COM 34 UNS 176 UNS 192 UNS 204 COM 92 INC 60 INC 166 DNO 24 DNO 46 DCR 24 DCR 36
DIA	DIAMETER OF PIPE	UNS 38
DIN	INITIAL VALUE OF DEPTH	INC 114 INC 122 INC 136 INC 144 UNS 202
DM	HYDRAULIC DEPTH	C1P 76
DMAX	MAXIMUM DEPTH	UNS 98 HO1 192 HO2 136 COM 106
DN	NORMAL DEPTH	DNO 40
DNOUT	DEPTH AT OUTLET	UNS 152 UNS 160
DP	DEPTH	HO1 62 HO1 108 COM 84
DP1	INITIAL VALUE OF DEPTH	HO1 62 HO1 102 HO2 58 HO2 102
DP2	COMPUTED VALUE OF DEPTH	HO1 98 HO2 98
DR	INTERPOLATED VALUE OF DEPTH	HO2 38 COM 56
DRA	(2)	INC 80
DS	INTERPOLATED VALUE OF DEPTH	HO1 44 HO1 52 COM 60
DSLO	DERIVATIVE OF ENERGY SLOPE WITH DEPTH	INC 84
DSMLL	(2)	HO2 56
DT	INCREMENT OF TIME	UNS 190
DTHET	DERIVATIVE OF THETA WITH DEPTH	INC 72
DTMAX	TIME OF MAXIMUM DEPTH	UNS 186
DTOL	MAXIMUM ERROR IN DEPTH CALCULATION	UNS 76
DUDD	(2)	HO2 40 HO2 90
DW	(2)	INC 76
DX	INCREMENT OF DISTANCE	UNS 166
D1	(2)	COE 34
ED	OUTLET DISCHARGE EXPONENT	UNS 64
EE1	ENERGY AT KNOWN DEPTH	INC 42 INC 148
EE2	ENERGY AT UNKNOWN DEPTH	INC 44
E2	(2)	COE 32
F	DISCHARGE RESISTANCE COEFF	INC 22 INC 60 INC 92 DNO 36 COE 44
F1	FACTOR IN BENT-MFL = BENT-FRAC	INC 44
F2	FACTOR IN BENT-MFL = BENT-FRAC	INC 44
F3	(2)	HO2 22 COM 44
F4	(2)	HO2 48 COM 96
F5	(2)	HO2 48
F6	(2)	HO2 48
F7	(2)	HO2 48
F8	(2)	HO2 48
F9	(2)	HO2 48
F10	(2)	HO2 48
F11	(2)	HO2 48
F12	(2)	HO2 48
F13	(2)	HO2 48
F14	(2)	HO2 48
F15	(2)	HO2 48
F16	(2)	HO2 48
F17	(2)	HO2 48
F18	(2)	HO2 48
F19	(2)	HO2 48
F20	(2)	HO2 48
F21	(2)	HO2 48
F22	(2)	HO2 48
F23	(2)	HO2 48
F24	(2)	HO2 48
F25	(2)	HO2 48
F26	(2)	HO2 48
F27	(2)	HO2 48
F28	(2)	HO2 48
F29	(2)	HO2 48
F30	(2)	HO2 48
F31	(2)	HO2 48
F32	(2)	HO2 48
F33	(2)	HO2 48
F34	(2)	HO2 48
F35	(2)	HO2 48
F36	(2)	HO2 48
F37	(2)	HO2 48
F38	(2)	HO2 48
F39	(2)	HO2 48
F40	(2)	HO2 48
F41	(2)	HO2 48
F42	(2)	HO2 48
F43	(2)	HO2 48
F44	(2)	HO2 48
F45	(2)	HO2 48
F46	(2)	HO2 48
F47	(2)	HO2 48
F48	(2)	HO2 48
F49	(2)	HO2 48
F50	(2)	HO2 48
F51	(2)	HO2 48
F52	(2)	HO2 48
F53	(2)	HO2 48
F54	(2)	HO2 48
F55	(2)	HO2 48
F56	(2)	HO2 48
F57	(2)	HO2 48
F58	(2)	HO2 48
F59	(2)	HO2 48
F60	(2)	HO2 48
F61	(2)	HO2 48
F62	(2)	HO2 48
F63	(2)	HO2 48
F64	(2)	HO2 48
F65	(2)	HO2 48
F66	(2)	HO2 48
F67	(2)	HO2 48
F68	(2)	HO2 48
F69	(2)	HO2 48
F70	(2)	HO2 48
F71	(2)	HO2 48
F72	(2)	HO2 48
F73	(2)	HO2 48
F74	(2)	HO2 48
F75	(2)	HO2 48
F76	(2)	HO2 48
F77	(2)	HO2 48
F78	(2)	HO2 48
F79	(2)	HO2 48
F80	(2)	HO2 48
F81	(2)	HO2 48
F82	(2)	HO2 48
F83	(2)	HO2 48
F84	(2)	HO2 48
F85	(2)	HO2 48
F86	(2)	HO2 48
F87	(2)	HO2 48
F88	(2)	HO2 48
F89	(2)	HO2 48
F90	(2)	HO2 48
F91	(2)	HO2 48
F92	(2)	HO2 48
F93	(2)	HO2 48
F94	(2)	HO2 48
F95	(2)	HO2 48
F96	(2)	HO2 48
F97	(2)	HO2 48
F98	(2)	HO2 48
F99	(2)	HO2 48
F100	(2)	HO2 48

FNU	KINEMATIC VISCOSITY	UNS 46
FRATIO	(2)	INC 100
GCM	(2)	HO1 74 COM 70
GCP	(2)	HO2 50 COM 48
GR	ACCELERATION OF GRAVITY	UNS 54
HFTM	(2)	INC 70
I	(1)	UNS 286
II	(1)	UNS 294 INC 192 INC 176
IL	LINE PRINTING INDEX	UNS 296
IQ	(1)	INF 42 INF 44
ITO	TIME INTERVAL BETWEEN OUTPUTS	UNS 194
ITOC	TIME, OUTPUT INTERVAL COUNTER	UNS 224 UNS 272 UNS 278
IXO	DISTANCE INTERVAL BETWEEN OUTPUTS	UNS 70
IXOC	DISTANCE, OUTPUT INTERVAL COUNTER	UNS 244 UNS 254 UNS 260
J	(1)	INC 158 UNS 226
L	(1)	INC 48
MC	BACKWATER PROFILE CODE	UNS 148 UNS 156
N	NUMBER OF X-INTERVALS	UNS 68
NCONT	ITERATION COUNTER	INC 52 INC 124
NP6	OUTPUT PAGE CONTROL	UNS 290
NQCD	NUMBER OF INPUT HYDROGRAPH POINTS	UNS 80
NT	NUMBER OF TIME INTERVALS	UNS 194
N1	NUMBER OF X-INTEGRATION LOCATIONS	UNS 94
Q	DISCHARGE AT TIME T	UNS 282 INC 50 INC 160 INC 184
QR	BASE DISCHARGE	UNS 84
QDT	DISCHARGE AT TIME T+DT	HO1 126 HO2 122 HO2 130 COM 96
QI	INPUT HYDROGRAPH DISCHARGE	UNS 82
QIN	INITIAL DISCHARGE	UNS 204 INF 48 INF 48
QMAX	MAXIMUM DISCHARGE AT TIME T	UNS 102 HO1 164 HO2 148 COM 118
QP	PEAK HYDROGRAPH DISCHARGE	UNS 86
QQ	DISCHARGE	UNS 134 UNS 172 HO1 58
QRA	PEAK TO BASE DISCHARGE RATIO	UNS 92
R	HYDRAULIC RADIUS	C1P 68
RD	DEPTH TO RADIUS RATIO	HO1 84 HO2 60
REY	REYNOLDS NUMBER	INC 34 INC 66 INC 90 DNO 32
REY	(2)	COE 42
SCM	(2)	HO1 76 COM 74
SCP	(2)	HO2 52 COM 72
SF	FRICTION SLOPE (AVERAGE)	INC 96 COE 50
S0	INVERT SLOPE	UNS 42
S1	FRICTION SLOPE AT KNOWN DEPTH	INC 40 INC 146
S2	FRICTION SLOPE AT UNKNOWN DEPTH	INC 94
T	TIME	UNS 210 UNS 228 INF 30
TCM	(2)	COM 78
TCP	(2)	COM 76
TDMAX	TIME TO MAXIMUM DEPTH	HO1 134 HO2 138 COM 108
TF	FINAL TIME FOR CALCULATION	UNS 72
THETA	CFNT ANG SUBTENDED BY FREE SURFACE	HO1 46 HO2 52 HO2 56
THETA2ONE-HALF THETA		HO2 76 HO2 88
TIO	TIME INTERVAL BETWEEN OUTPUTS	UNS 74
TP	TIME TO PEAK INPUT DISCHARGE	UNS 88
TQ	INFLOW HYDROGRAPH TIME	UNS 82
TOMAX	TIME TO MAXIMUM DISCHARGE	HO1 146 HO2 150 COM 120
TVMAX	TIME TO MAXIMUM VELOCITY	HO1 140 HO2 144 COM 114
U	VELOCITY	HO2 72 HO2 84
UN	(2)	HO1 40 COM 52
UP	(2)	HO2 34 COM 50
V	VELOCITY AT TIME T	UNS 286 INC 48 INC 158 INC 182
VC	CRITICAL VELOCITY	UNS 180
VDT	VELOCITY AT TIME T+DT	HO1 124 HO2 124 HO2 128 COM 100
VH	VELOCITY HEAD	INC 30 INC 88
VMAX	MAXIMUM VELOCITY	UNS 100 HO1 138 HO2 142 COM 112
VOL	VOLUME OF HYDROGRAPH WAVE	UNS 98
VP	VELOCITY	HO1 120 COM 42
VR	INTERPOLATED VALUE OF VELOCITY	HO2 40 COM 58
VS	INTERPOLATED VALUE OF VELOCITY	HO1 40 HO1 54 COM 62
VV	VELOCITY	HO1 50 HO2 44 COM 56 DNO 28
W	WETTED PERIMETER	INC 44 INC 78 INC 124
X	POSITION ALONG CHANNEL	INC 44 INC 124
Y	WATER SURFACE ELEVATION	INC 44
Z	INITIAL LENGTH OF CHANNEL	UNS 60
ZL	STARTING POSITION	HO2 50
ZR	ENDING POSITION	HO2 50

(1) = 100 - 1000 (2) = 1000 - 10000 (3) = 10000 - 100000 (4) = 100000 - 1000000 (5) = 1000000 - 10000000 (6) = 10000000 - 100000000 (7) = 100000000 - 1000000000 (8) = 1000000000 - 10000000000 (9) = 10000000000 - 100000000000 (10) = 100000000000 - 1000000000000 (11) = 1000000000000 - 10000000000000 (12) = 10000000000000 - 100000000000000 (13) = 100000000000000 - 1000000000000000 (14) = 1000000000000000 - 10000000000000000 (15) = 10000000000000000 - 100000000000000000 (16) = 100000000000000000 - 1000000000000000000 (17) = 1000000000000000000 - 10000000000000000000 (18) = 10000000000000000000 - 100000000000000000000 (19) = 100000000000000000000 - 1000000000000000000000 (20) = 1000000000000000000000 - 10000000000000000000000 (21) = 10000000000000000000000 - 100000000000000000000000 (22) = 100000000000000000000000 - 1000000000000000000000000 (23) = 1000000000000000000000000 - 10000000000000000000000000 (24) = 10000000000000000000000000 - 100000000000000000000000000 (25) = 100000000000000000000000000 - 1000000000000000000000000000 (26) = 1000000000000000000000000000 - 10000000000000000000000000000 (27) = 10000000000000000000000000000 - 100000000000000000000000000000 (28) = 100000000000000000000000000000 - 1000000000000000000000000000000 (29) = 1000000000000000000000000000000 - 10000000000000000000000000000000 (30) = 10000000000000000000000000000000 - 100000000000000000000000000000000 (31) = 100000000000000000000000000000000 - 1000000000000000000000000000000000 (32) = 1000000000000000000000000000000000 - 10000000000000000000000000000000000 (33) = 10000000000000000000000000000000000 - 100000000000000000000000000000000000 (34) = 100000000000000000000000000000000000 - 1000000000000000000000000000000000000 (35) = 1000000000000000000000000000000000000 - 10000000000000000000000000000000000000 (36) = 10000000000000000000000000000000000000 - 100000000000000000000000000000000000000 (37) = 100000000000000000000000000000000000000 - 1000000000000000000000000000000000000000 (38) = 1000000000000000000000000000000000000000 - 100 (39) = 100 - 1000 (40) = 1000 - 100 (41) = 100 - 1000 (42) = 1000 - 100 (43) = 100 - 1000 (44) = 1000 - 100 (45) = 100 - 1000 (46) = 1000 - 100 (47) = 100 - 1000 (48) = 1000 - 100 (49) = 100 - 1000 (50) = 1000 - 100 (51) = 100 - 1000 (52) = 1000 - 100 (53) = 100 - 1000 (54) = 1000 - 10000

### A.3.4. SAMPLE INPUT AND OUTPUT

**SAMPLE INPUT**  
 Same format as in A.1.4.

S 0.0 4.0 30.0 10.0 50.0 10. 80.0 4.0  
 200.0 4.0

**SAMPLE OUTPUT**

**INFLOW HYDROGRAPH PARAMETERS**  
 QB = 4.00000CFS  
 QP = 10.00000CFS  
 TP = 30.00000SEC  
 QB/JP = 4.0000  
 WAVE VOLUME ABOVE BASE FLOW = 180.00CU FT

**SYSTEM PARAMETERS**  
 SO = 0.0100  
 ALPHA = 1.00000  
 BETA = 1.00000  
 N = 20  
 IX0 = 2  
 TF = 20V  
 TIO = 20.00000  
 NORMAL DEPTH = 7.659FT CRITICAL DEPTH = 6.290FT

DX = 40.94348FT DT = 1.45574SEC

CONDITIONS AT 0.000SECONDS			
DISTANCE (FT)	DEPTH (FT)	VELOCITY (FPS)	DISCHARGE (CFS)
.0000	.7659	2.8522	4.0000
81.8870	.7658	2.8524	4.0000
163.7739	.7658	2.8528	4.0000
245.6609	.7636	2.8535	4.0000
327.5478	.7654	2.8542	4.0000
409.4348	.7648	2.8578	4.0000
491.3217	.7637	2.8639	4.0000
573.2087	.7612	2.8768	4.0000
655.0956	.7559	2.9052	4.0000
736.9826	.7433	2.9749	4.0000
818.8695	.6290	3.7688	4.0000

CONDITIONS AT 18.925SECONDS			
DISTANCE (FT)	DEPTH (FT)	VELOCITY (FPS)	DISCHARGE (CFS)
.0000	.9539	4.0904	7.7849
81.8870	.8362	3.3542	5.3194
163.7739	.7734	2.9100	4.1378
245.6609	.7659	2.8556	4.0049
327.5478	.7654	2.8550	4.0001
409.4348	.7648	2.8572	4.0000
491.3217	.7636	2.8643	4.0000
573.2087	.7612	2.8771	4.0000
655.0956	.7559	2.9071	4.0024
736.9826	.7427	3.0034	4.0338
818.8695	.6269	3.7940	4.0992

CONDITIONS AT 37.849SECONDS			
DISTANCE (FT)	DEPTH (FT)	VELOCITY (FPS)	DISCHARGE (CFS)
.0000	1.0700	4.4918	10.0000
81.8870	1.0308	4.4354	9.3859
163.7739	.9176	3.8867	7.0221
245.6609	.8053	3.1387	4.7230
327.5478	.7704	2.8920	4.0842
409.4348	.7651	2.8604	4.0002
491.3217	.7636	2.8643	4.0004
573.2087	.7612	2.8782	4.0020
655.0956	.7561	2.9125	4.0117
736.9826	.7390	3.0322	4.0441
818.8695	.6259	3.7904	4.0858

CONDITIONS AT 56.774SECONDS			
DISTANCE (FT)	DEPTH (FT)	VELOCITY (FPS)	DISCHARGE (CFS)
.0000	1.0404	4.0341	8.6452
81.8870	1.0060	4.3000	9.5438
163.7739	1.0442	4.3710	9.4136
245.6609	1.0048	4.3342	8.8619
327.5478	.8731	3.5932	6.0508
409.4348	.7883	3.0274	4.4214
491.3217	.7668	2.8881	4.0573
573.2087	.7616	2.8815	4.0091
655.0956	.7556	2.9140	4.0125
736.9826	.7453	3.0543	4.0500
818.8695	.6347	3.7870	4.0714

CONDITIONS AT 75.699SECONDS			
DISTANCE (FT)	DEPTH (FT)	VELOCITY (FPS)	DISCHARGE (CFS)
.0000	.3836	2.8384	4.8603
81.8870	.3967	2.8554	4.7178
163.7739	1.0199	3.9327	8.2021
245.6609	1.0257	4.1370	8.6958
327.5478	1.0230	4.3124	9.0317
409.4348	.7942	4.1384	7.9931
491.3217	.8344	3.3584	5.3101
573.2087	.7753	2.9823	4.2551
655.0956	.7568	2.9366	4.0500
736.9826	.7325	3.0785	4.0551
818.8695	.6240	3.7846	4.0621

CONDITIONS AT 94.623SECONDS			
DISTANCE (FT)	DEPTH (FT)	VELOCITY (FPS)	DISCHARGE (CFS)
.0000	.8295	2.5505	4.0000
81.8870	.8422	2.6384	4.2239
163.7739	.9002	3.0444	5.3496
245.6609	.9579	3.5015	6.7027
327.5478	.9923	3.8510	7.7304
409.4348	.9932	4.0117	8.0694
491.3217	1.0019	4.2687	9.6895
573.2087	.9129	3.8842	6.9592
655.0956	.7989	3.2461	4.8296
736.9826	.7390	3.1523	4.2044
818.8695	.6374	3.7954	4.1048

CONDITIONS AT 113.548SECONDS			
DISTANCE (FT)	DEPTH (FT)	VELOCITY (FPS)	DISCHARGE (CFS)
.0000	.8186	2.5984	4.0000
81.8870	.8251	2.6335	4.1063
163.7739	.8313	2.6343	4.1438
245.6609	.8615	2.8325	4.6823
327.5478	.9054	3.1650	5.6051
409.4348	.9484	3.5343	6.6769
491.3217	.9623	3.7555	7.2345
573.2087	.9696	3.9720	7.7310
655.0956	.9642	4.2182	8.1409
736.9826	.9394	3.8314	6.1098
818.8695	.7013	3.9941	4.9473

CONDITIONS AT 132.472SECONDS			
DISTANCE (FT)	DEPTH (FT)	VELOCITY (FPS)	DISCHARGE (CFS)
.0000	.8100	2.6348	4.0000
81.8870	.8177	2.6036	4.0942
163.7739	.8259	2.7020	4.2186
245.6609	.8290	2.6875	4.2113
327.5478	.8424	2.7624	4.4272
409.4348	.8704	2.9678	4.9761
491.3217	.9055	3.2611	5.7785
573.2087	.9322	3.5523	6.5494
655.0956	.9294	3.7297	6.8521
736.9826	.9372	4.1744	7.7537
818.8695	.8042	4.4732	7.4271

CONDITIONS AT 151.397SECONDS			
DISTANCE (FT)	DEPTH (FT)	VELOCITY (FPS)	DISCHARGE (CFS)
.0000	.8030	2.6600	4.0000
81.8870	.8100	2.6875	4.0770
163.7739	.8168	2.7072	4.1558
245.6609	.8265	2.7501	4.2411
327.5478	.8437	2.7525	4.3182
409.4348	.8342	2.7684	4.3757
491.3217	.8495	2.8817	4.6718
573.2087	.8720	3.0865	5.1840
655.0956	.8956	3.3874	5.9100
736.9826	.8662	3.5284	6.2386
818.8695	.8074	4.3000	6.5087

CONDITIONS AT 170.321SECONDS			
DISTANCE (FT)	DEPTH (FT)	VELOCITY (FPS)	DISCHARGE (CFS)
.0000	.7973	2.6952	4.0000
81.8870	.8036	2.7086	4.0632
163.7739	.8101	2.7291	4.1407
245.6609	.8161	2.7427	4.2043
327.5478	.8246	2.7814	4.3201
409.4348	.8301	2.8077	4.4074
491.3217	.8310	2.8107	4.4187
573.2087	.8360	2.8713	4.5524
655.0956	.8431	3.0100	4.8410
736.9826	.8439	3.3117	5.3636
818.8695	.7584	4.1660	5.7630

CONDITIONS AT 189.246SECONDS			
DISTANCE (FT)	DEPTH (FT)	VELOCITY (FPS)	DISCHARGE (CFS)
.0000	.7925	2.7188	4.0000
81.8870	.7991	2.7281	4.0537
163.7739	.8037	2.7414	4.1137
245.6609	.8102	2.7507	4.1494
327.5478	.8155	2.7742	4.2485
409.4348	.8220	2.8041	4.3119
491.3217	.8280	2.8471	4.4539
573.2087	.8272	2.8701	4.4839
655.0956	.8229	2.9204	4.5433
736.9826	.8097	3.0944	4.6919
818.8695	.7053	4.0044	5.0030

**MAXIMUM VALUES AND TIMES AT EACH LOCATION**

DISTANCE	MAX DEPTH	TIME	MAX VEL	TIME	MAX Q	TIME
.00	1.09	49.50	4.24	30.57	10.00	30.57
40.94	1.08	52.41	4.43	36.39	9.86	40.76
81.89	1.07	55.32	4.47	40.76	9.71	45.13
122.83	1.06	59.69	4.45	46.58	9.59	48.04
163.77	1.05	64.05	4.43	50.95	9.48	52.41
204.72	1.04	61.14	4.42	55.32	9.39	58.23
245.66	1.03	65.51	4.40	61.14	9.31	62.60
286.60	1.03	69.88	4.39	65.51	9.22	66.96
327.55	1.02	74.24	4.37	71.33	9.13	72.79
368.49	1.02	80.07	4.35	77.15	9.02	78.61
409.43	1.01	84.43	4.33	81.52	8.92	82.98
450.38	1.01	90.26	4.30	87.34	8.81	89.30
491.32	1.00	94.62	4.28	93.17	8.69	94.62
532.27	1.00	100.45	4.26	98.99	8.57	98.99
573.21	.99	105.27	4.24	103.36	8.46	104.81
614.15	.98	110.64	4.22	109.18	8.31	110.54
655.10	.97	116.46	4.22	115.00	8.21	115.00
696.04	.95	122.28	4.22	119.37	7.99	120.83
736.98	.95	128.10	4.24	125.10	8.13	124.65
777.93	.89	135.44	4.28	132.47	7.37	133.93
818.87	.87	143.93	4.45	133.93	7.44	133.93

Key Words: Finite-Difference Schemes, Unsteady Flow Equations, Method of Characteristics, Numerical Solutions of Differential Equations.

Abstract: This fourth part of a four-part series of hydrology papers on flood routing through storm drains presents the computer-oriented numerical methods on solving the quasi-linear hyperbolic partial differential equations known as De Saint-Venant equations of gradually varied free-surface unsteady flow. Formulation of various numerical finite-difference schemes either explicit schemes based on the two partial differential equations, unstable, diffusing, upstream differencing, leap frog, and Lax-Wendroff or the specified intervals scheme based on the method of characteristics is analyzed. A comparison between the specified intervals scheme of the method of characteristics, the Lax-Wendroff scheme and the diffusing scheme is discussed. Flow charts and computer programs for these various numerical methods are given in the appendices.

Reference: Yevjevich, Vujica and Albert H. Barnes, Colorado State University, Hydrology Paper No. 46 (November 1970) "Flood Routing Through Storm Drains, Part IV, Numerical Computer Methods of Solution".

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